

Coexistence of different structures for hadrons -- the case of the low-lying Ω_c baryons --

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1. What (structures) make hadrons
2. Ω_c baryons
-- Coexistence of MB and qqq structures --

Exotic hadrons from QCD

Coexistence of molecular and compact quark structures
near threshold regions

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1. Introduction — Coexistence
2. Pentaquarks, P_c and P_{cs}
3. Baryon octet of $sss = \Omega(2012)$
4. Summary

1. What (structures) make hadrons

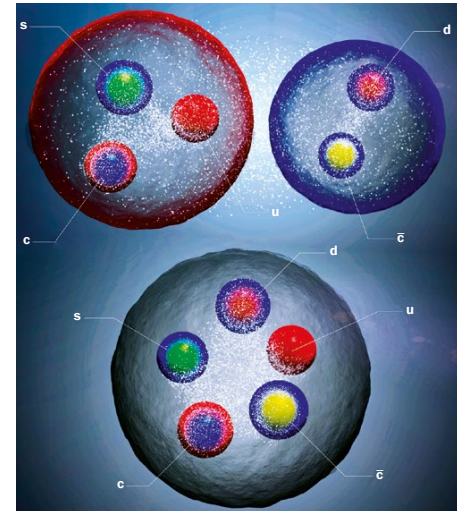
-- Coexistence of **MB** and *qqq* structures --



CERN COURIER , NOVEMBER/DECEMBER 2024

INSIDE PENTAQUARKS AND TETRAQUARKS

Marek Karliner and Jonathan Rosner ask what makes tetraquarks and pentaquarks tick, revealing them to be at times exotic **compact** states, at times hadronic **molecules** and at times **both** — with much still to be discovered.



But this is not a new story

Long history in hadron and nuclear physics

For $\Delta(1232)$

Doctor thesis

粒子と共鳴準位の混合効果について

益川 英
名古屋大学 物理教室



T. Maskawa

Progress of Theoretical Physics, Vol. 38, No. 1, July 1967

Mixing Effect between Particles and Resonances

Published paper

Toshihide MASKAWA, Hiroki KONDO
and Ziro MAKI*

Department of Physics, Nagoya University, Nagoya

**Research Institute for Fundamental Physics
Kyoto University, Kyoto*

(Received February 23, 1967)

p190-201, Only 3 citations

Molecular like states via clusterization

Not only in hadrons but also in nuclei

Alpha cluster (molecular) structure of nuclei

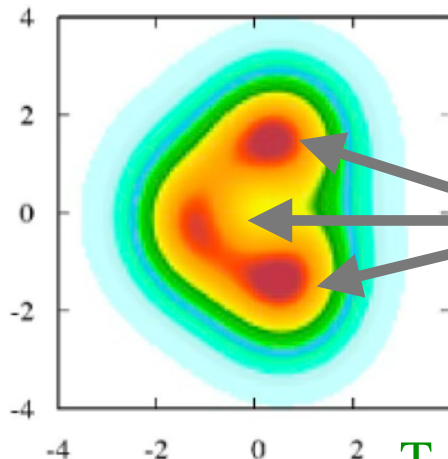


Prog. Theor. Phys. 40, 277 (1968)
Ikeda diagram

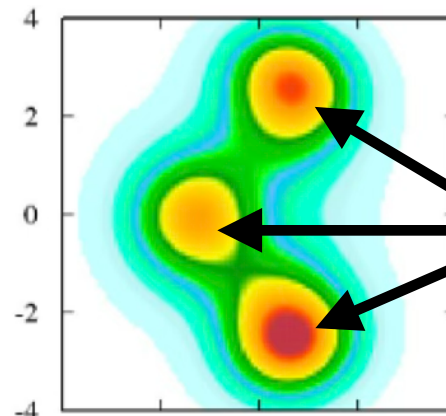
^8Be	^{12}C	^{16}O	^{20}Ne	^{24}Mg	^{28}Si	^{32}S
∞	∞ (7.27)	∞ (14.44)	∞ (19.17)	∞ (28.48)	∞ (38.46)	∞ (45.41)
	C	$\text{C} \circ$ (7.16)	$\text{C} \infty$ (11.89)	$\text{C} \infty$ (21.21)	$\text{C} \infty$ (31.19)	$\text{C} \infty$ (38.14)
		O	$\text{O} \circ$ (4.73)	$\text{O} \infty$ (14.05)	$\text{O} \infty$ (24.03)	$\text{O} \infty$ (30.96)
			$\text{C} \text{C}$ (13.93)	$\text{C} \text{C}$ (23.91)	$\text{C} \text{C}$ (30.86)	
			Ne	$\text{Ne} \circ$ (9.32)	$\text{Ne} \infty$ (18.29)	$\text{Ne} \infty$ (26.25)
				$\text{O} \text{C}$ (16.75)	$\text{O} \text{C}$ (23.70)	$\text{O} \text{C}$ (30.86)
			Mg	$\text{Mg} \circ$ (9.98)	$\text{Mg} \infty$ (19.96)	$\text{Mg} \infty$ (26.93)
				Si	$\text{Si} \circ$ (6.95)	$\text{Si} \infty$ (16.54)

Density distributions

^{12}C : Ground state



Disolved clusters

 ^{12}C : Excited (Hoyle) state

Well developed clusters

T. Otsuka, Y. Utsuno et al., Nat. Comm. 13, 2234 (2022)

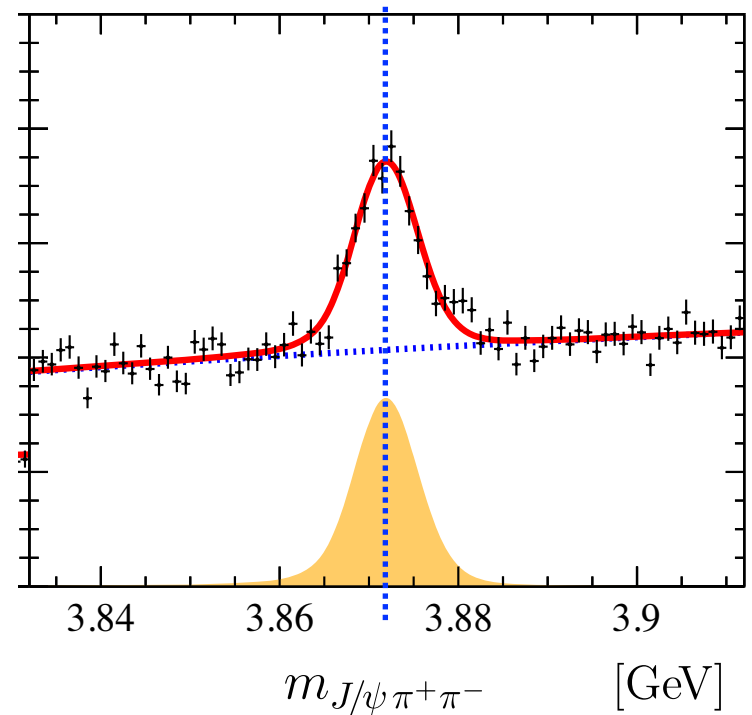
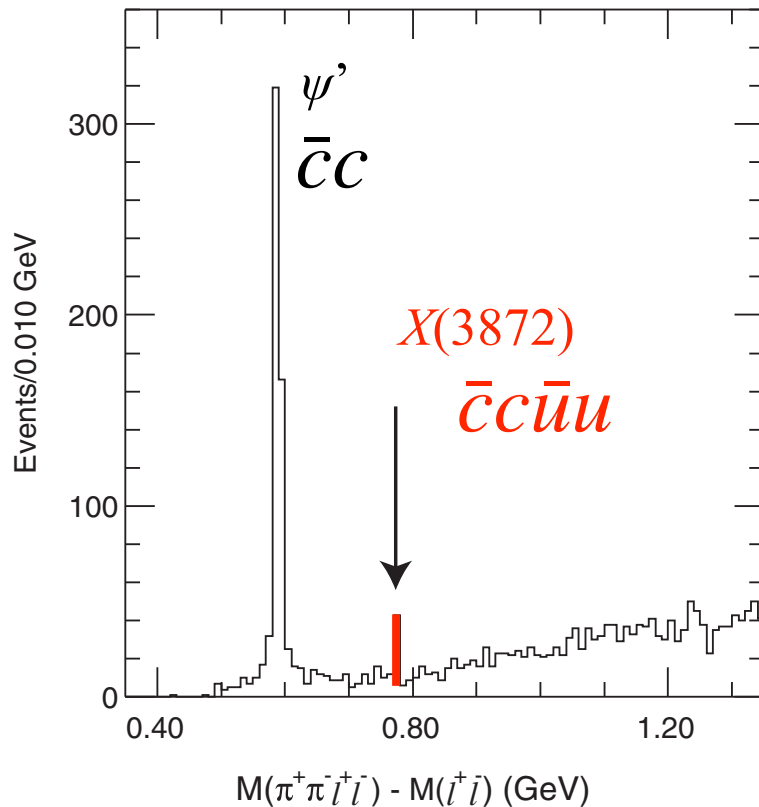
X(3872)

Belle@KEK, PRL91 262001 (2003)

LHCb, PRD 102, 092005 (2020)

$$B^\pm \rightarrow (J/\psi \pi^+ \pi^-) K^\pm$$

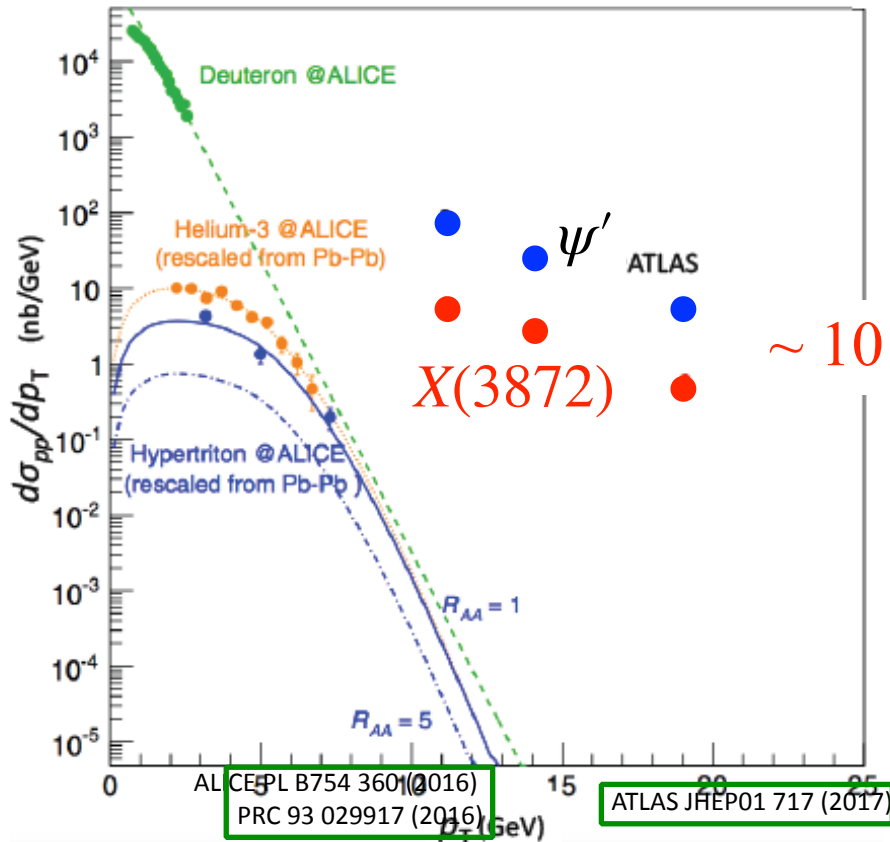
pp collisions at center-of-mass
energies of 7 and 8 TeV



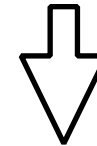
The peaks are almost at the threshold of $D\bar{D}^* \rightarrow$ **Hadronic molecule**

BUT

See Esposito et al., PRD 92 034028 (2015)



Large production rates
at large E (\sim TeV)
and p_T ($>$ GeV)



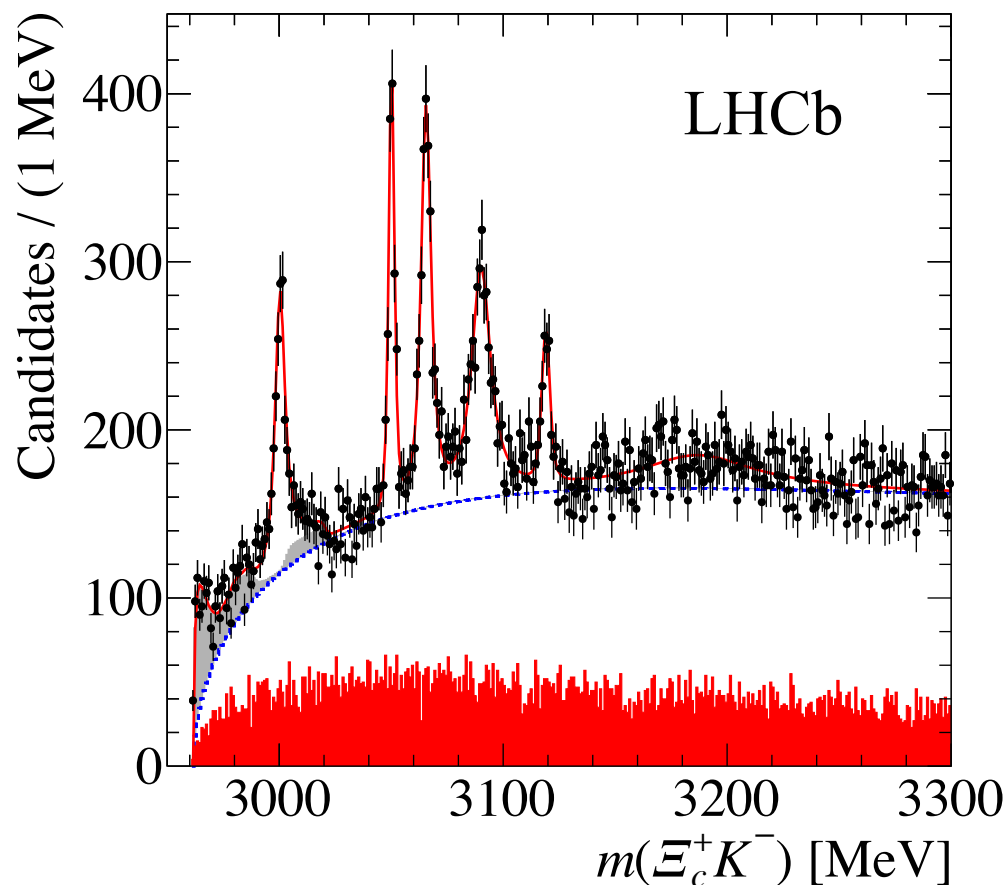
Implies admixture of
 $D^0 \bar{D}^{*0}$ and
compact quark core $\sim c\bar{c}$

We have performed analysis

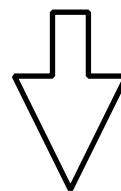
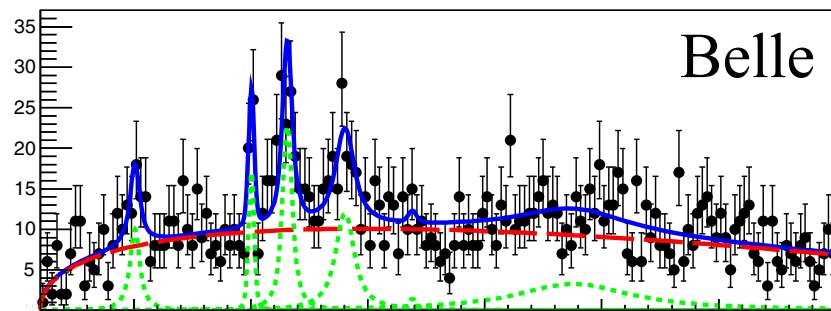
2. Ω_c baryons

Invariant mass analysis of $\Xi_c^+ K^- \rightarrow SSC$

Phys.Rev.Lett. 118 (2017) no.18, 182001

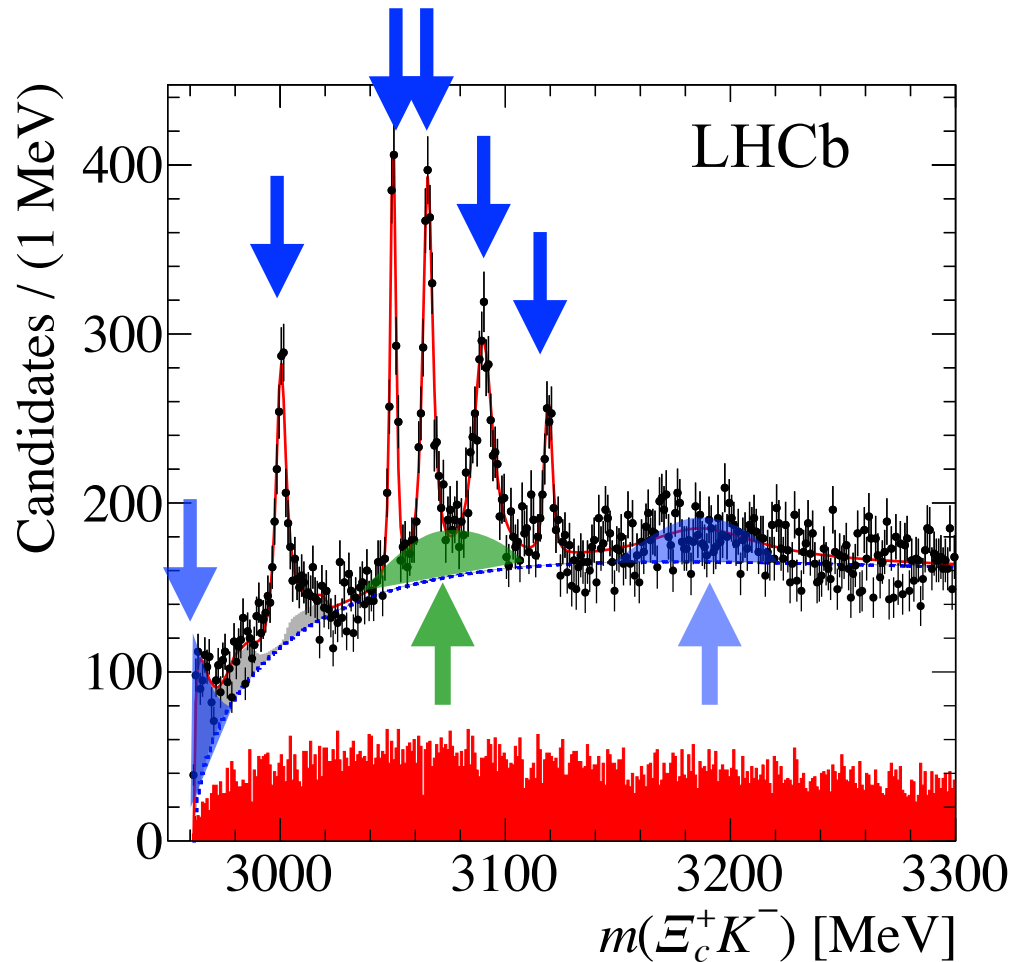


Phys.Rev.D97, 051102(R) (2018)

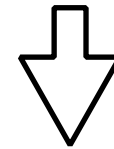


5 or 4 clear peaks

If we look more carefully . . .



- 5 sharp peaks
- Threshold enhancement
- Broad bump ~ 3200 MeV
- Broad bump ~ 3100 MeV



6 (or 8) signals

More from LHCb

Phys.Rev. D 104, L091102 (2021): Angular distribution \rightarrow **spin assignment**

Phys. Rev. Lett.131, 131902 (2023): More states $\Omega_c(3185), \Omega_c(3327)$

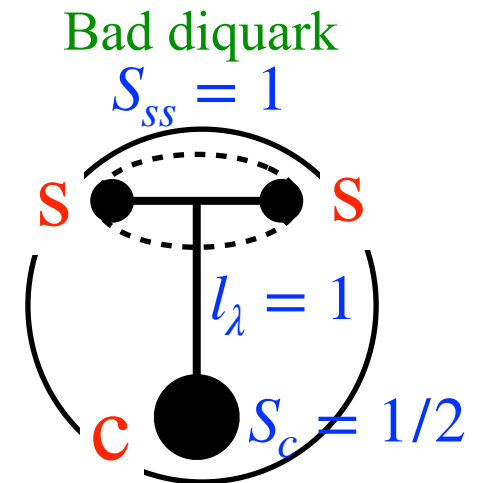
Option 1:

5 peaks implies **Quark core** structure

-- **λ mode** P-wave excitations --

Total spin J :

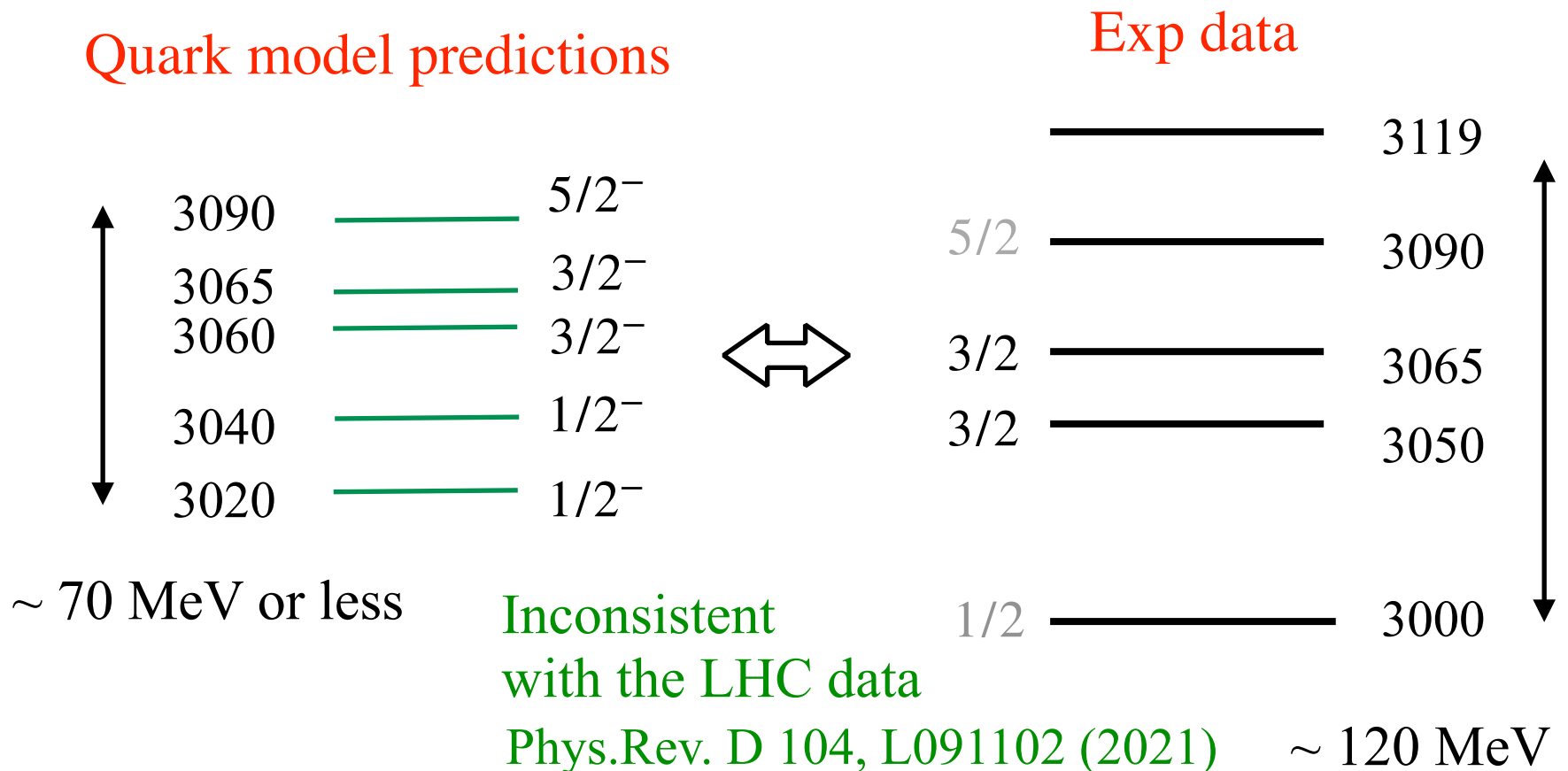
$$J = (S_{ss} + l_\lambda) + s_c$$
$$= \begin{cases} 0 & + & 1/2 & = & 1/2 \\ 1 & + & 1/2 & = & 1/2, 3/2 \\ 2 & + & 1/2 & = & 3/2, 5/2 \end{cases}$$



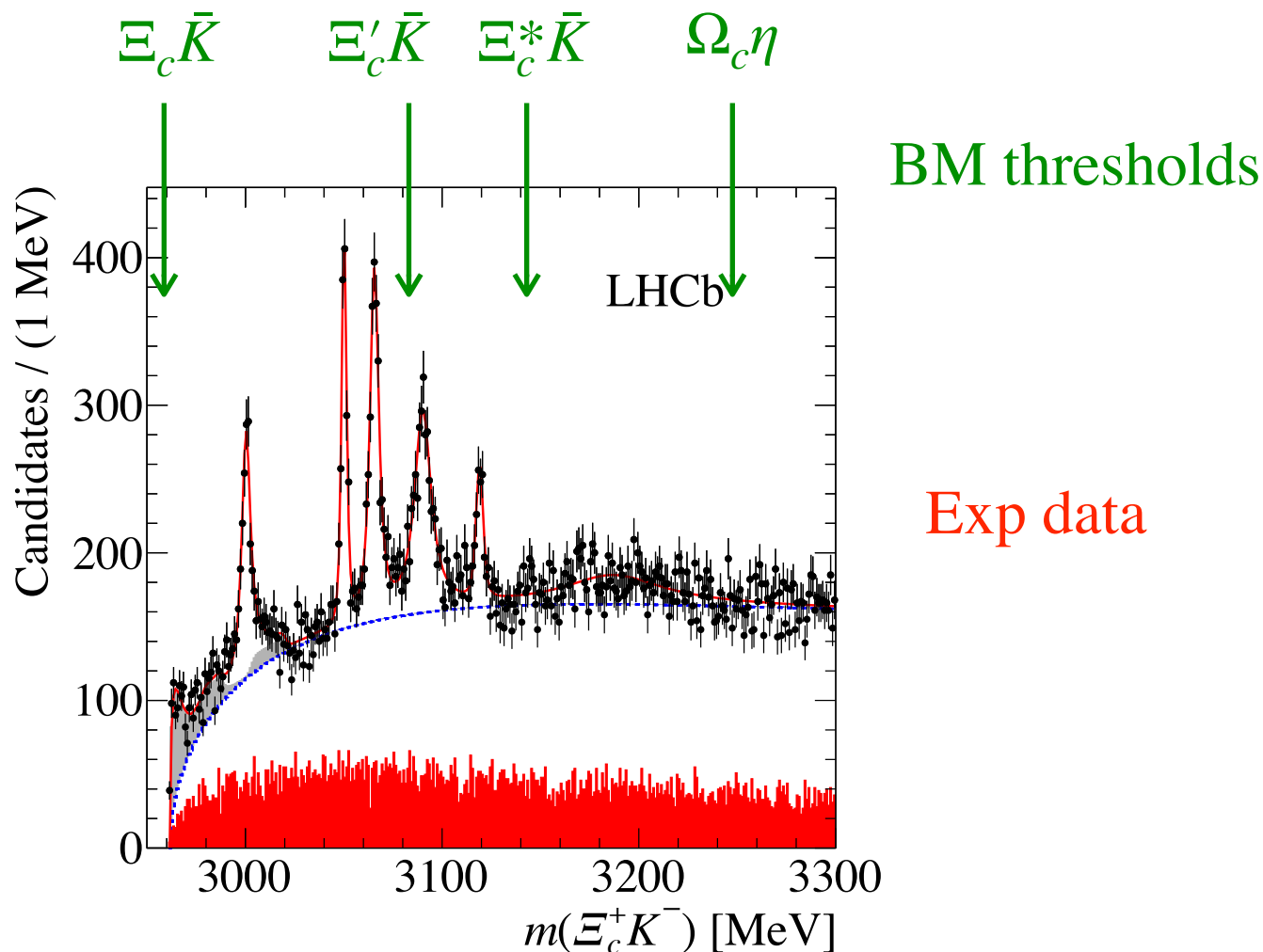
There are **5 states(!)** from lower λ modes
 $J^P = 1/2^-, 1/2^-, 3/2^-, 3/2^-, 5/2^-$

BUT difficulties

- (1) They distribute in a narrow energy range
Spin assignments do not follow data



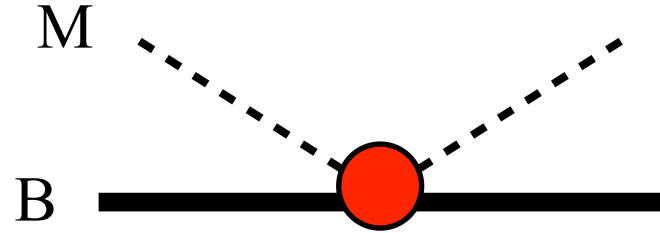
(2) There are MB channels ($\bar{K}\Xi_c$, $\bar{K}\Xi'_c$, $\bar{K}\Xi_c^*$) nearby.



Option 2:

Near threshold implies **MB molecule**

Weinberg-Tomozawa int.



$$L_{KKqq} = -F \frac{i}{8f_\pi^2} \text{tr}[\bar{B} \gamma^\mu V_\mu, B] - D \frac{i}{8f_\pi^2} \text{tr}\{\bar{B} \gamma^\mu V_\mu, B\}$$

Attractive *MB* channels form bound/resonant/virtual states

Option 3, Our choice:

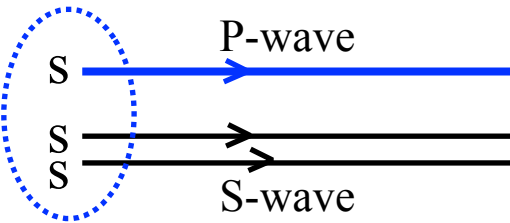
Coexistence of quark core and MB molecule

Understanding the low-lying Ω_c structures from a coupled-channel perspective

Zhang, Song, **Lu**, **Nagahiro** and Hosaka, Phys. Rev. D112, 0340035 (2025)

Quark core

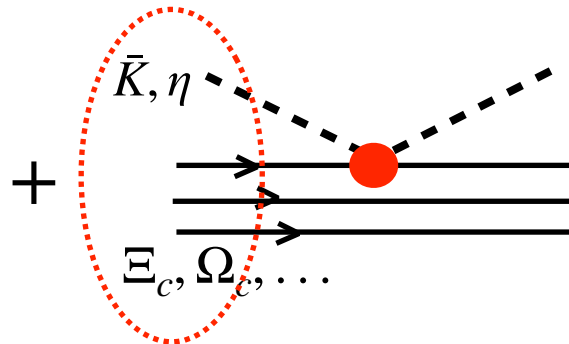
$P_{1/2,1/2, 3/2,3/2, 5/2}$



Quark model

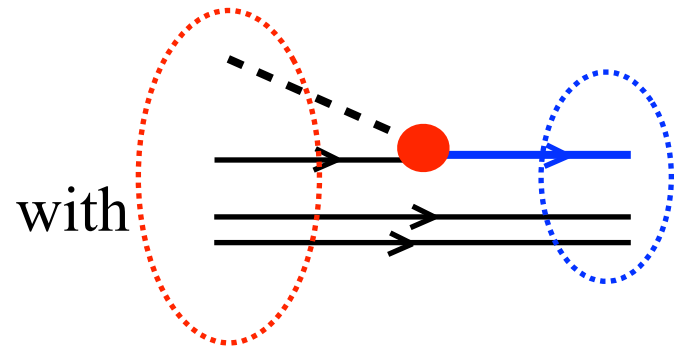
BM molecule

$\Xi_c \bar{K}, \dots$



S-wave interaction
Weinberg-Tomozawa

Coupling




Yukawa interaction

$$L_{KKqq} = -\frac{i}{8f_\pi^2} \bar{q} \gamma^\mu (\partial_\mu K^\dagger K - K^\dagger \partial_\mu K) q$$

$$L_{Kqq} = \frac{g_A}{2f_\pi} \bar{q} \gamma^\mu \gamma_5 \partial_\mu K q$$

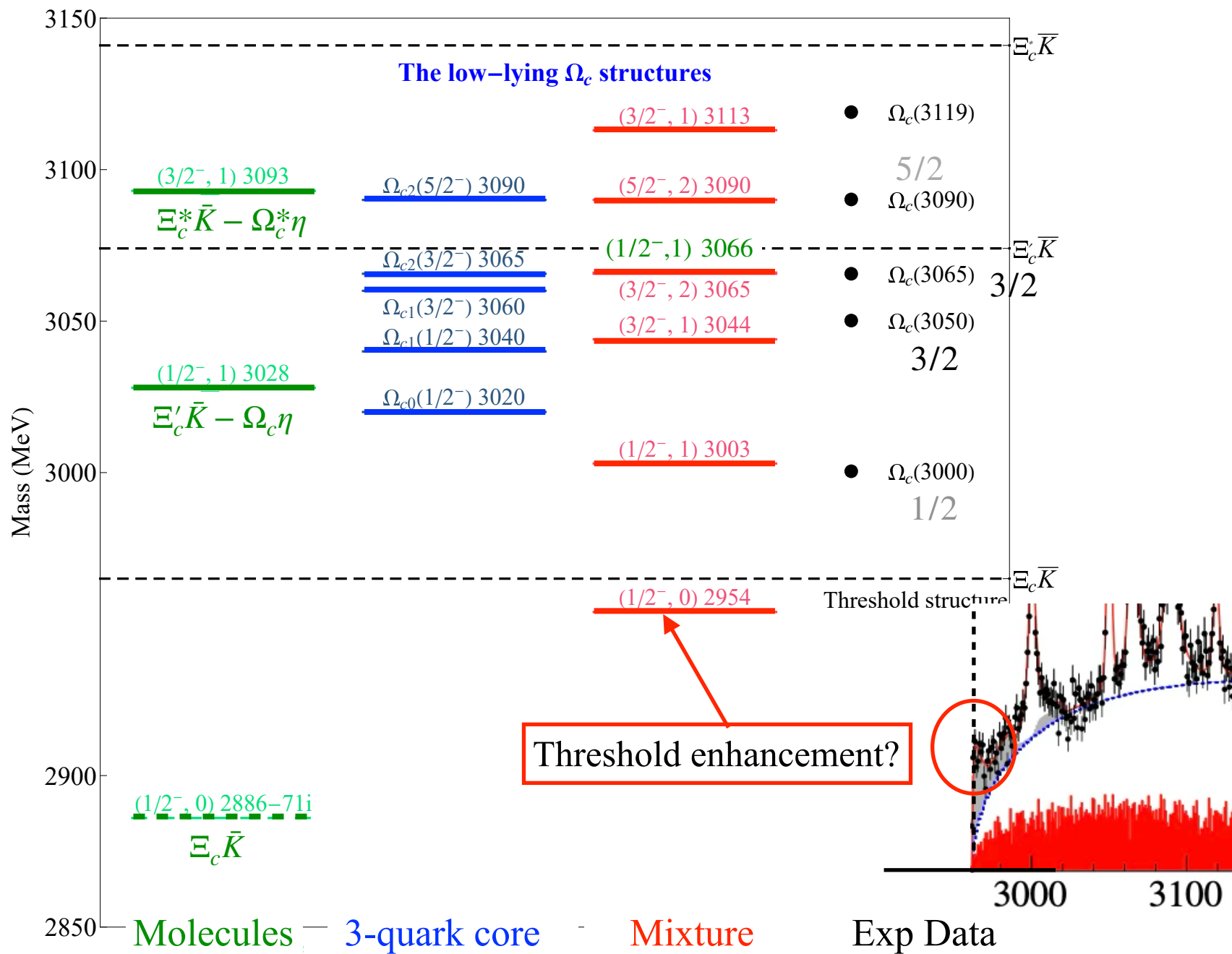
Coupling schemes for [Molecule]-[3quark]

$$\Omega_c \sim [{}^3P_{j,1/2_c}]^J$$


Light spin

Heavy quark symmetry
constrains the MB combination

Light spin	J^P	3-quarks	Molecule
$j = 2$	$5/2^-$	$\Omega_c({}^3P_{2,5/2})$	—
	$3/2^-$	$\Omega_c({}^3P_{2,3/2})$	—
$j = 1$	$3/2^-$	$\Omega_c({}^3P_{1,3/2})$	$\Xi_c^* \bar{K}, \Omega_c^* \eta$
	$1/2^-$	$\Omega_c({}^3P_{1,1/2})$	$\Xi_c' \bar{K}, \Omega_c \eta$
$j = 0$	$1/2^-$	$\Omega_c({}^3P_{0,1/2})$	$\Xi_c \bar{K}$



Summary

- Coexistence of different structures is relevant for various hadrons.
- Ω_c baryons has been studied.
- Computation of decay widths is in progress
- All Ω_c 's could be explained by coexistence with a prediction below the $\bar{K}\Xi_c$ threshold.