

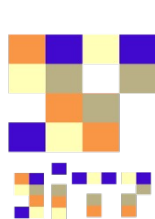
Lattice Study of Critical Point and Roberge-Weiss Phase Transition in Heavy-Quark Region of QCD

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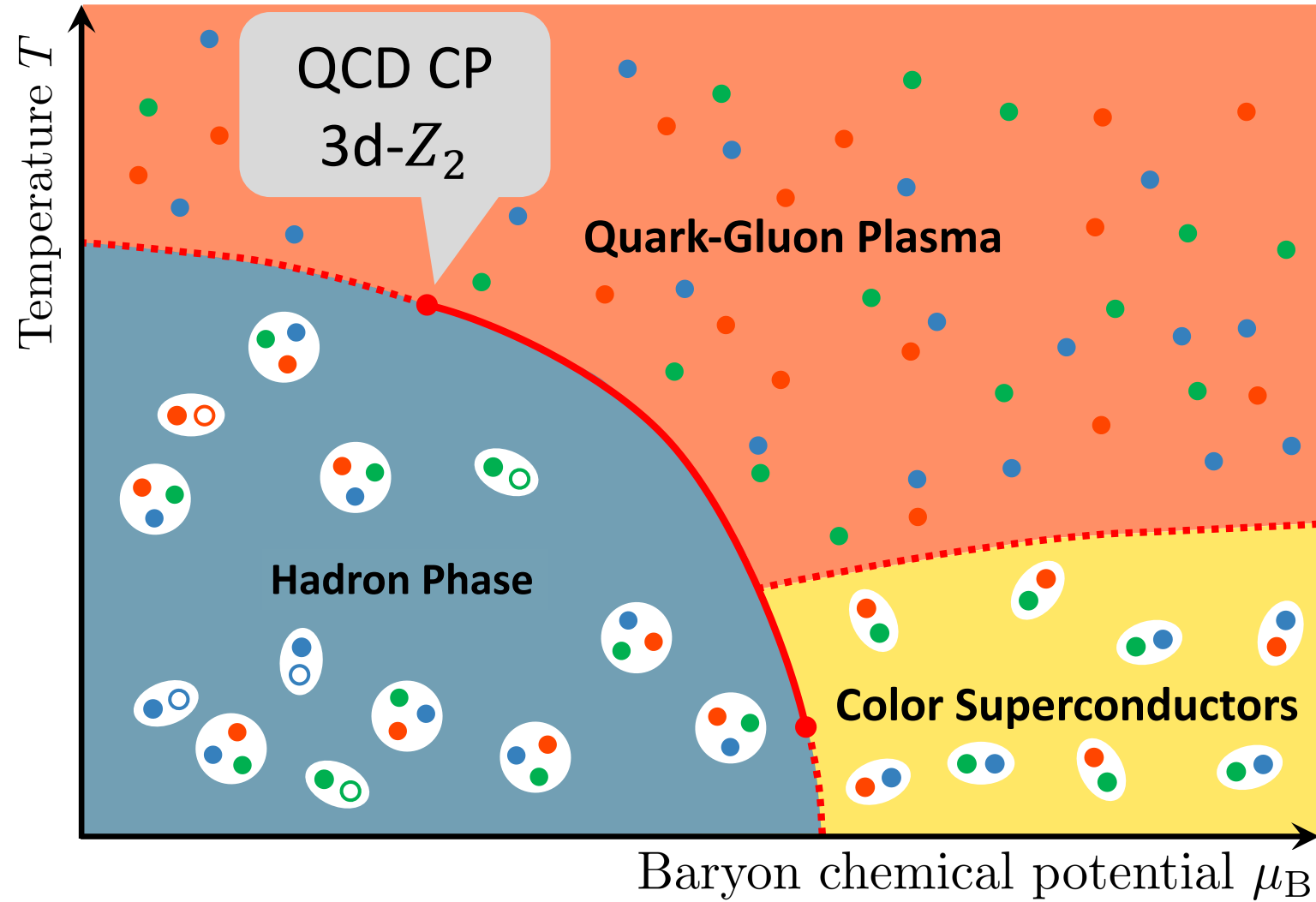
This talk is based on PRL 134, 162302, arXiv: 2508.20422[JPSJ, in press]

Institute of Physics, Academia Sinica



QCD Critical Point

QCD Phase diagram



QCD Critical Point(CP)

◆ Constraint to behavior of physical quantities

◆ Classification by universality class

QCD is asymptotic-free theory
Non-perturbative in low energy.
Analytic approach is too difficult!

◆ Effective Model Approach

➤ NJL model etc.

◆ FRG Approach

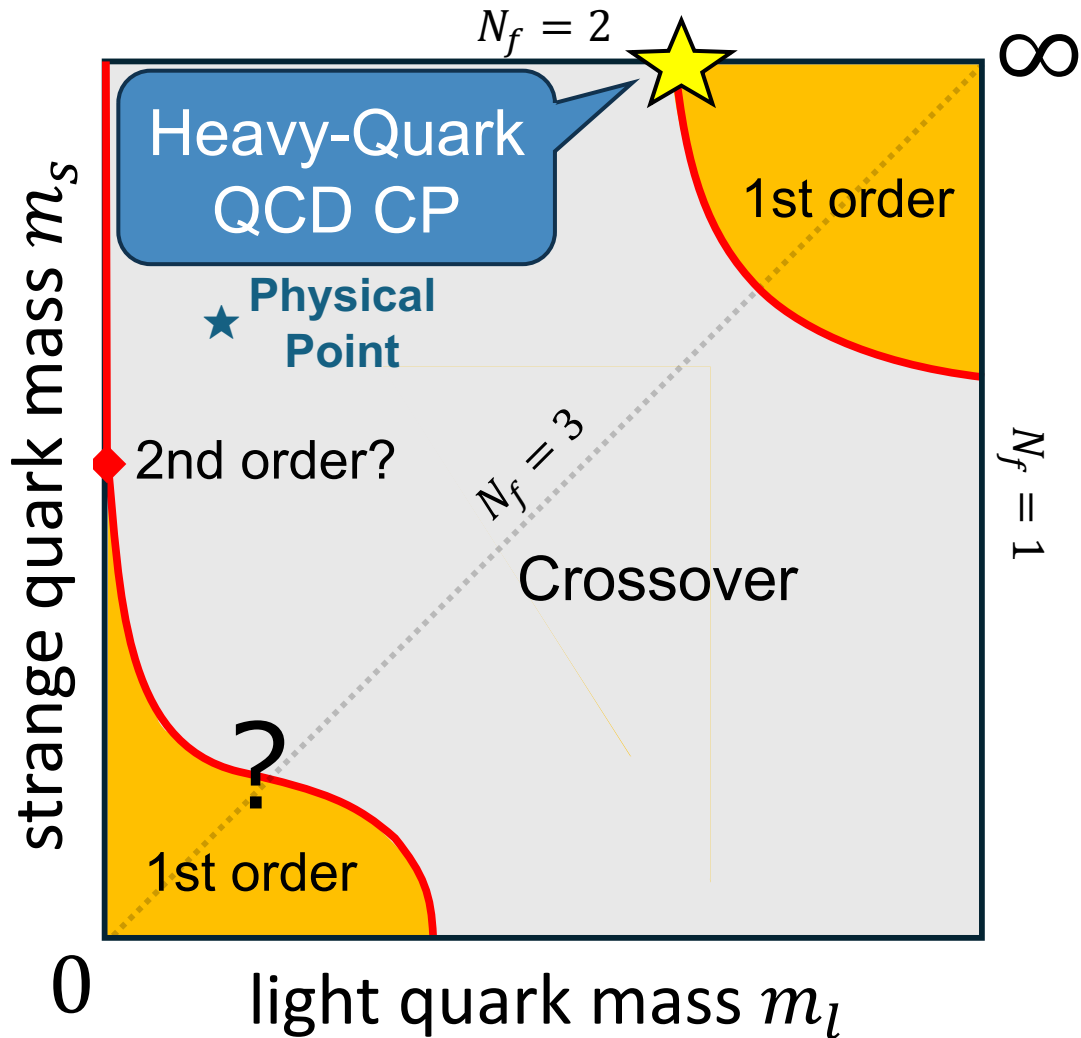
◆ Numerical Approach

➤ Monte Carlo simulation

➤ Tensor Network

➤ Quantum Computation

Columbia Plot



Columbia Plot :

The order of phase transition at $\mu = 0$

- ◆ $m \rightarrow m_{\text{phys}} : N_f = 2 + 1 \text{ QCD}$
→ Crossover
- ◆ $m \sim 0$: Light quark sector
→ 1st or Crossover is controversial.
(by Tsukuba group (~ 2020))
- ◆ $m \rightarrow \infty$: Pure Yang-Mills theory
→ 1st order

Method

Lee-Yang Zero Approach

Recent progress: Lee-Yang zeros

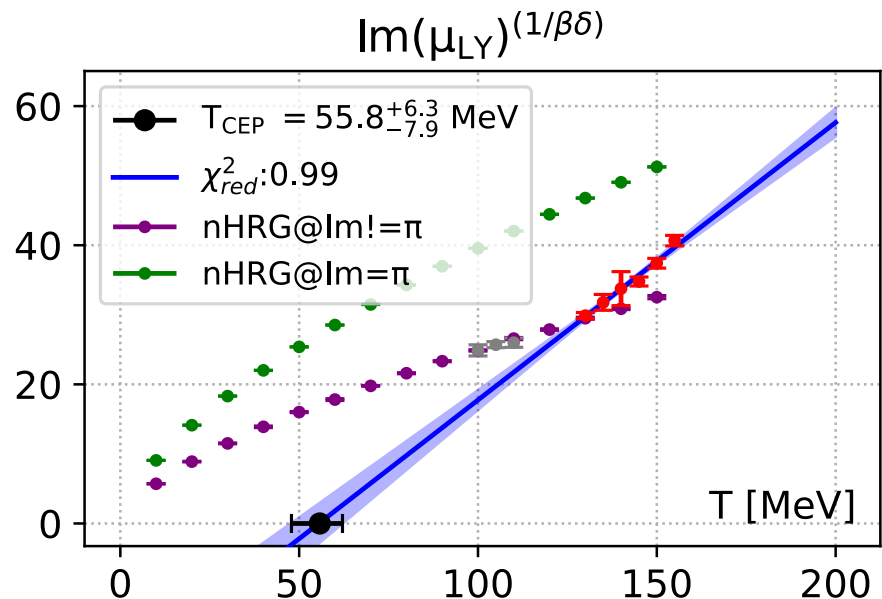
Pure imaginary chemical potential
+ Pade approximation

+Lee-Yang zeros /edge singularity

➤ Clarke, *et al.*(2024)

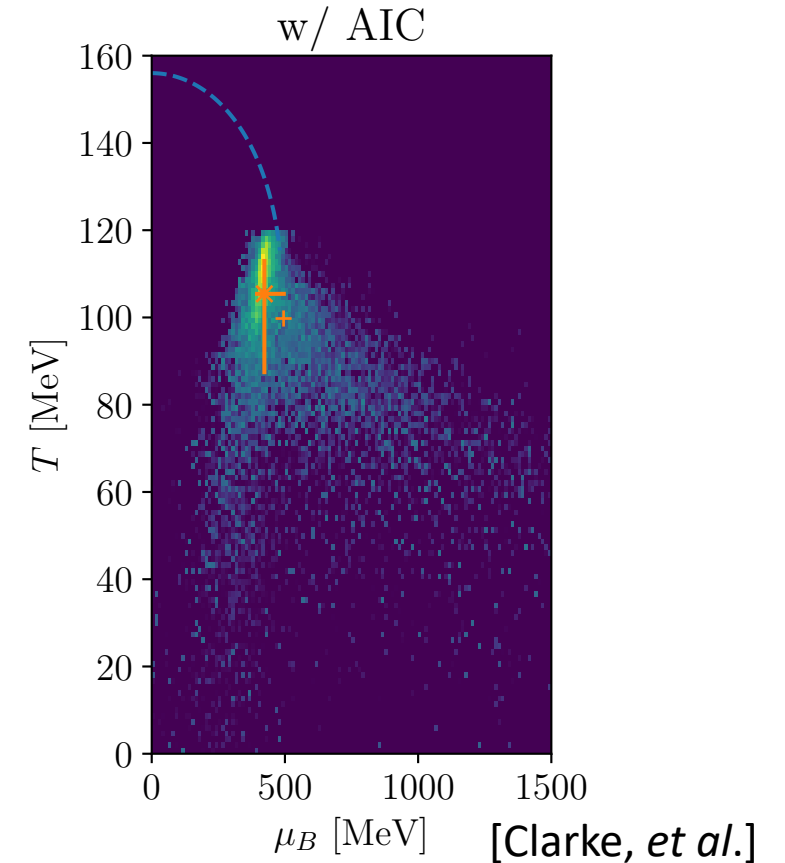
➤ Basar (2024)

➤ Alexander, *et al.*(2025)



$T_c \sim 70\text{MeV}$
or CP does not exist!
[Alexander, *et al.*]

Estimation of location of QCD CP



$$\begin{cases} \mu^{CEP} = 422^{+80}_{-35} \text{ MeV} \\ T^{CEP} = 105^{+8}_{-18} \text{ MeV} \end{cases}$$

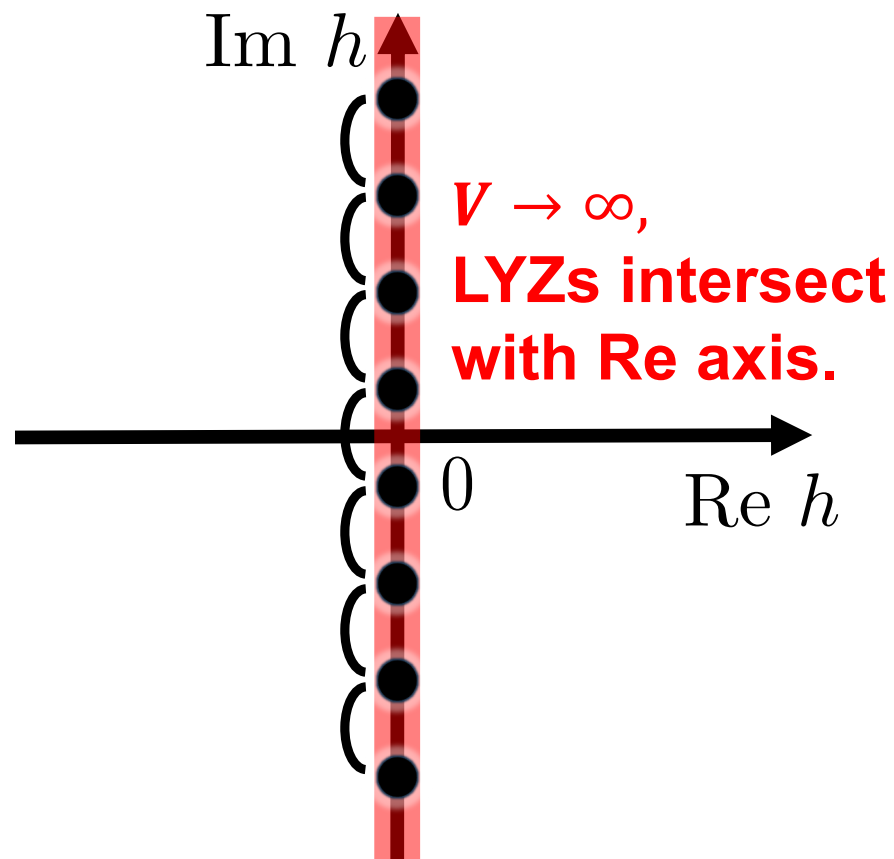
Review : Lee-Yang Zero

Lee, Yang 1952

Lee-Yang Zeros

= Zeros of Z in Complex Parameter space

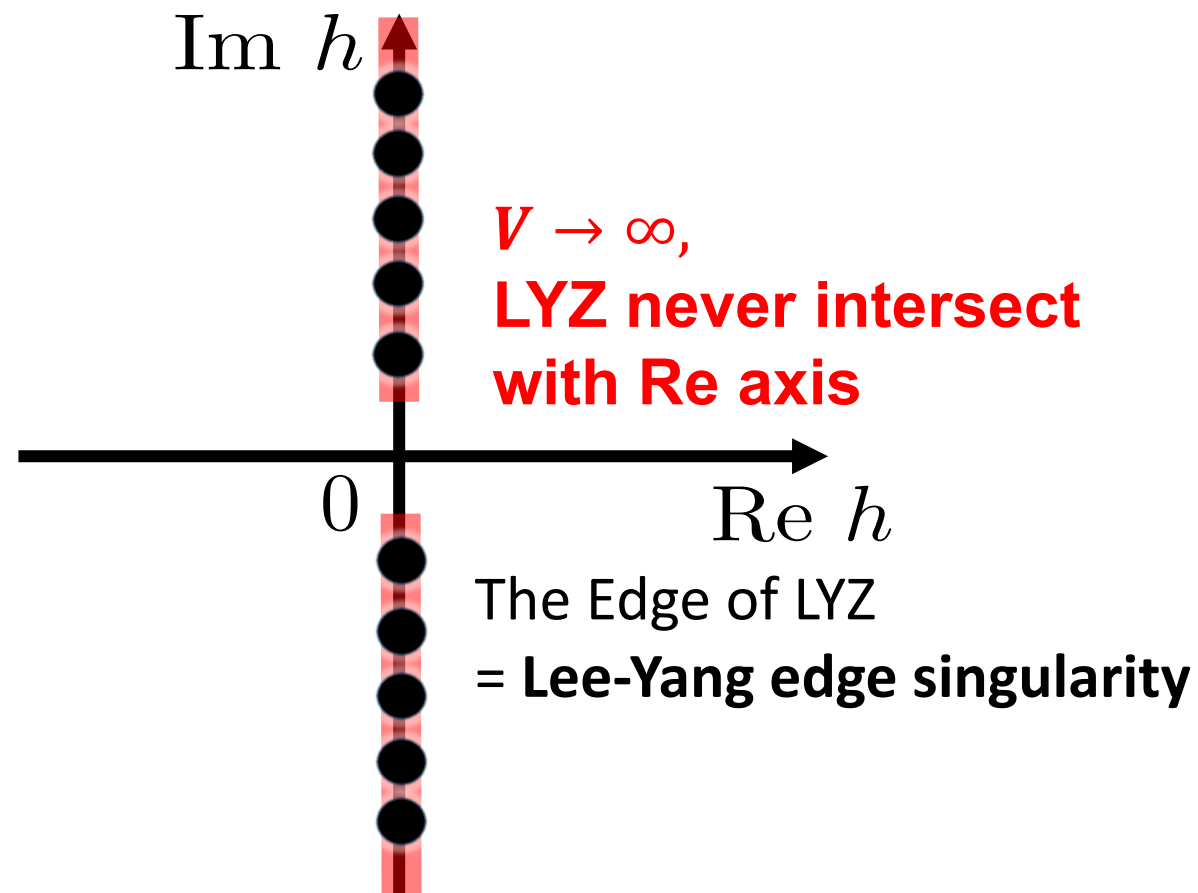
$t < 0$: 1st order



Ex.) d-dimensional Ising model $s_i = \pm 1$

$$Z(t, h, L) = \text{Tr} \exp \left(\frac{1}{T} \sum_{\langle i, j \rangle} s_i s_j + \frac{h}{T} \sum_i s_i \right)$$

$t > 0$: Crossover



Lee-Yang Zero Ratio

TW, Kitazawa, Kanaya (2025)

LYZ Ratio

$$R_{nm}(t, L) \equiv \frac{h^{(n)}(t, L)}{h^{(m)}(t, L)}$$



$$R_{nm}(t, L) \xrightarrow{L \rightarrow \infty} \begin{cases} \frac{2n-1}{2m-1} & (t < 0) \\ 1 & (t > 0) \end{cases}$$

◆ $t < 0$: 1st order phase transition

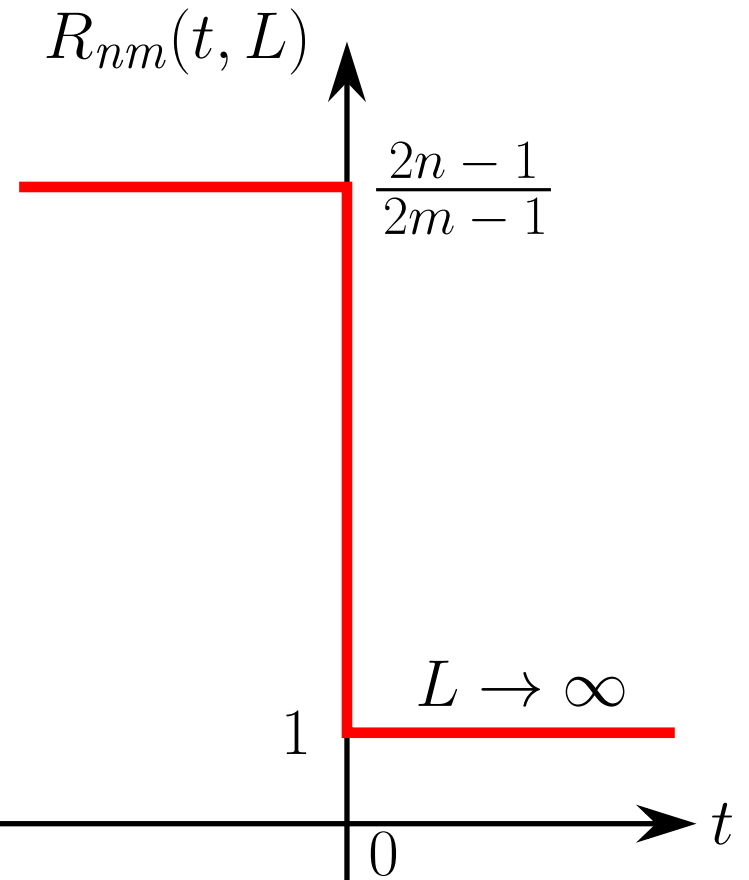
The distance between LYZ and Re-axis is 1:3:5... in $V \rightarrow \infty$

→ Ratio = $(2n-1)/(2m-1)$

◆ $t > 0$: Crossover

All LYZ accumulate at LYES in $V \rightarrow \infty$

→ Ratio = 1



How is the behavior in finite volume?

Finite Size Scaling (FSS)

FSS of Free Energy $\beta F_{\text{sing}}(t, h, L^{-1}) = -\log Z(t, h, L^{-1})$

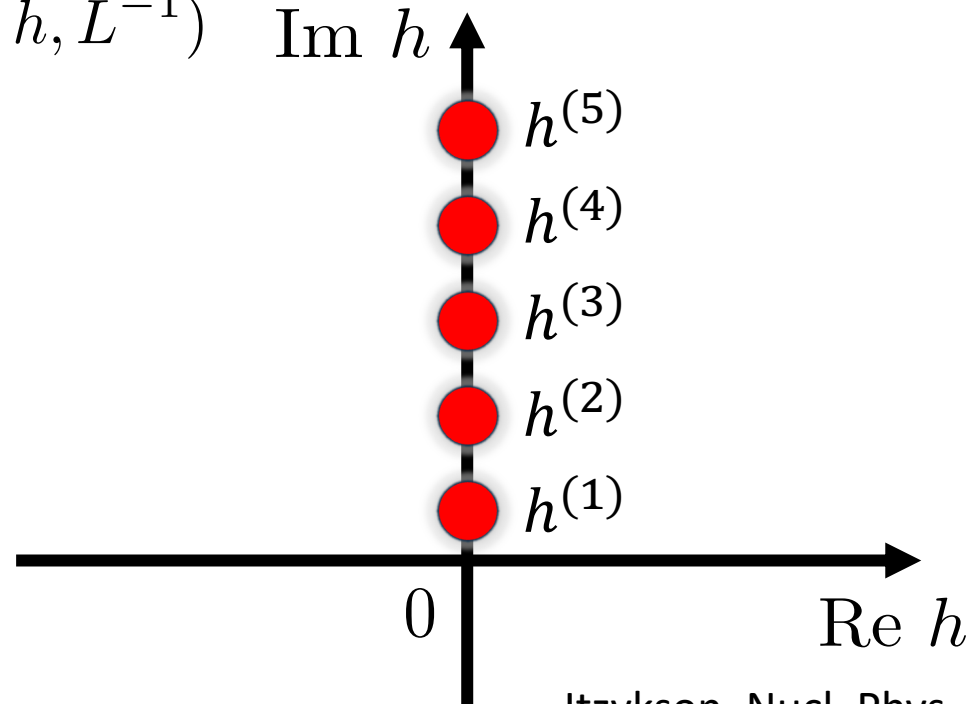
$$F_{\text{sing}}(t, h, L^{-1}) = \tilde{F}_{\text{sing}}(L^{y_t} t, L^{y_h} h)$$

FSS of Partition function

$$Z_{\text{sing}}(t, h, L^{-1}) = \tilde{Z}_{\text{sing}}(L^{y_t} t, L^{y_h} h)$$

Factorization of Partition function

$$Z_{\text{sing}}(t, h, L^{-1}) = \prod_{n=1}^N (h - h^{(n)}(t, L^{-1}))$$



Itzykson, Nucl. Phys. **220** (1983)

LYZ function

$$h = h_{\text{LY}}^{(n)}(t, L^{-1})$$

FSS of LYZ

$$L^{y_h} h_{\text{LY}}^{(n)}(t, L^{-1}) = \tilde{h}_{\text{LY}}^{(n)}(L^{y_t} t)$$

We can treat finite volume using FSS in the vicinity of the CP!

LYZR near CP

TW, Kitazawa, Kanaya (2025)

FSS & Taylor Expansion near the CP

$$L^{y_h} h_{LY}^{(n)}(t, L^{-1}) = X_n + Y_n L^{y_t} t + \mathcal{O}(t^2)$$

$$\gg R_{nm}(t, L) \equiv \frac{h^{(n)}(t, L)}{h^{(m)}(t, L)} = \frac{X_n + Y_n L^{y_t} t + \mathcal{O}(t^2)}{X_m + Y_m L^{y_t} t + \mathcal{O}(t^2)}$$

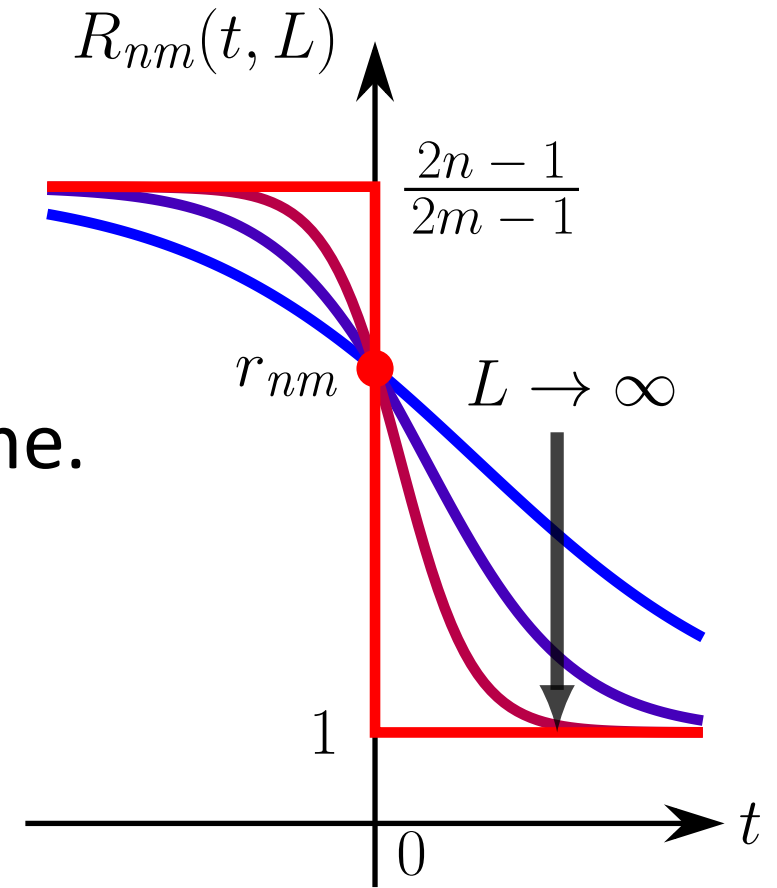
$$\gg R_{nm}(t, L) = (r_{nm} + c_{nm} L^{y_t} t + \mathcal{O}(t^2))$$

$R_{nm}(t = 0)$ does not depend on the volume.

→ Useful to locate CP

Cf. Similar to Binder Cumulant $B_4(t = 0) = b_4$

Crossing Point = Critical Point



Lattice Simulation in HQ-QCD

Lattice Simulation

- Wilson Fermion + Hopping parameter expansion $\rightarrow$ Heatbath method

$$\kappa \sim \frac{1}{2am} \text{ Hopping parameter}$$

- Lattice size : $N_s^3 \times (N_t = 4)$

Aspect Ratio **$LT = N_s/N_t = 8, 9, 10, 12$**

- # measurement = 6×10^5

$\rightarrow$ **Reweighting method**

- For fixed $\beta \in \mathbb{R}$, search LYZ **complex λ -plane**

$$\lambda \equiv 64N_c N_f \kappa^4$$

$$S_G \sim \square$$

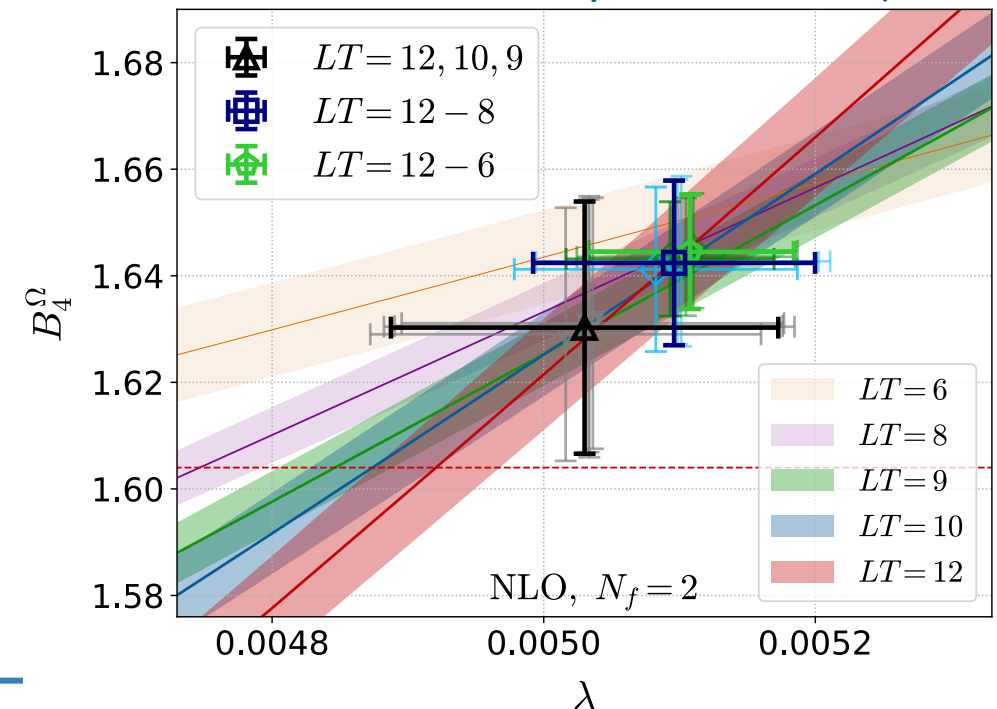
$$S_{LO} \sim \square + \text{cylinder}$$

Generation of Config.

$$S_{NLO} \sim \text{two squares} + \text{cube} + \text{cylinder} \quad \text{Reweighting method}$$

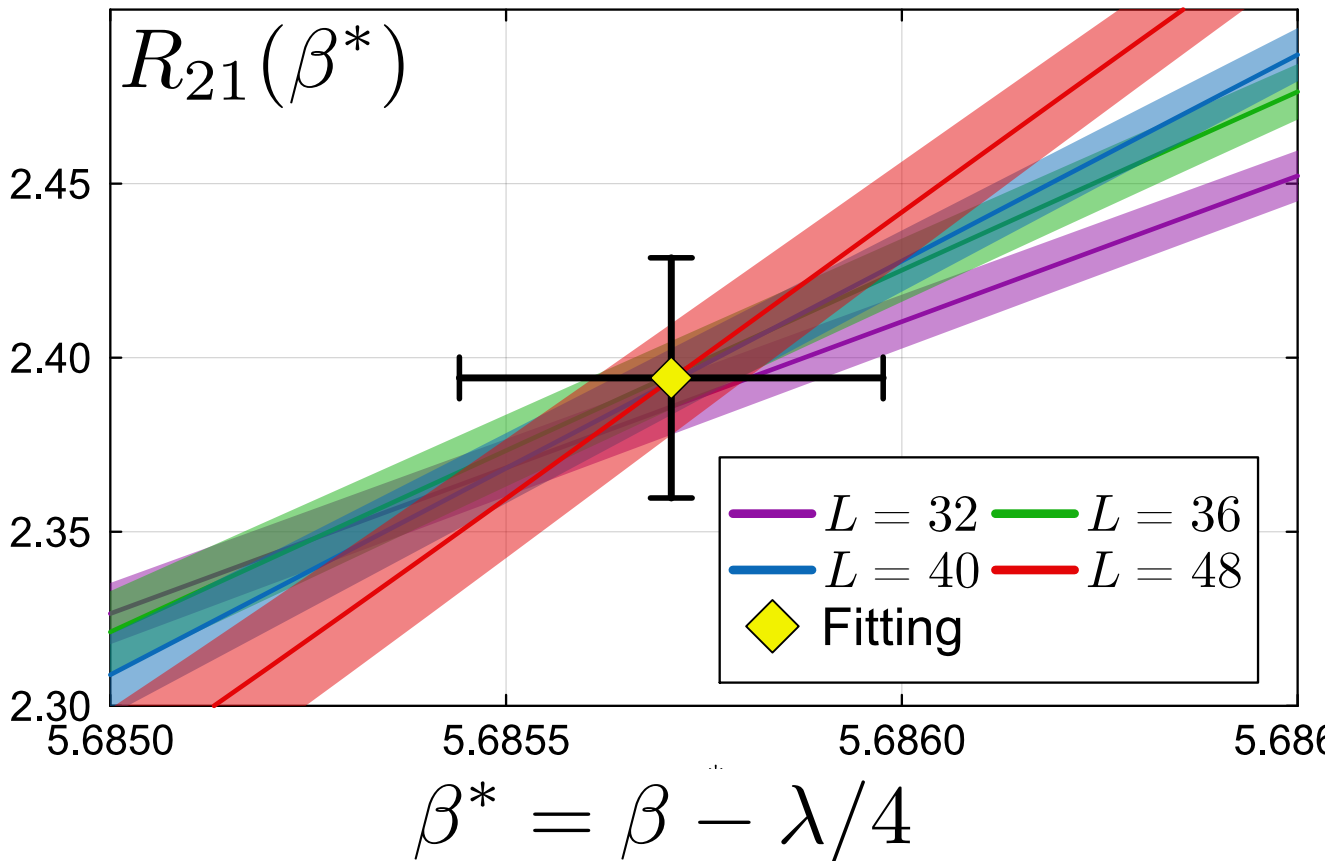
Binder Cumulant analysis

Kiyohara, et al. (2021)



Heavy-Quark QCD CP

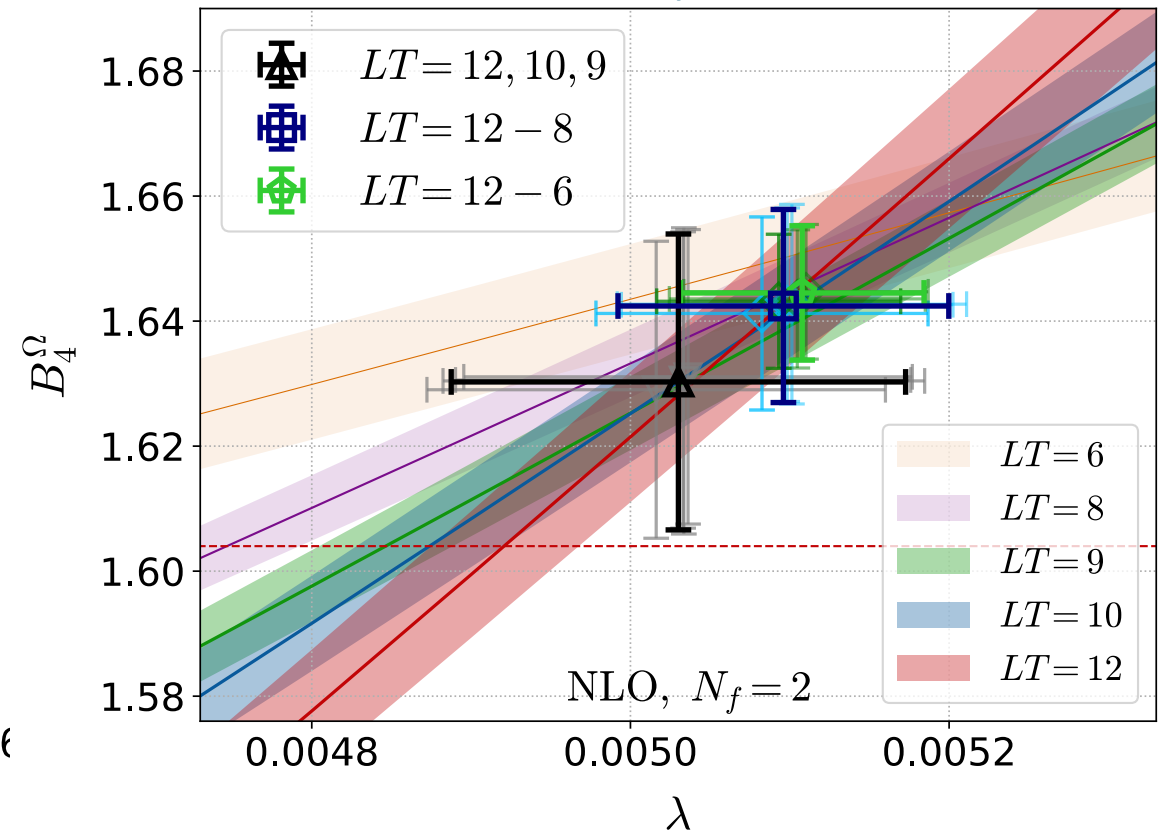
2nd/1st LYZ Ratio



Binder : $\beta^* = 5.68579(26), \lambda = 0.00503(14)(2)$
 LYZR21 : $\beta^* = 5.68571(27), \lambda = 0.00472(14)$

Binder Cumulant analysis

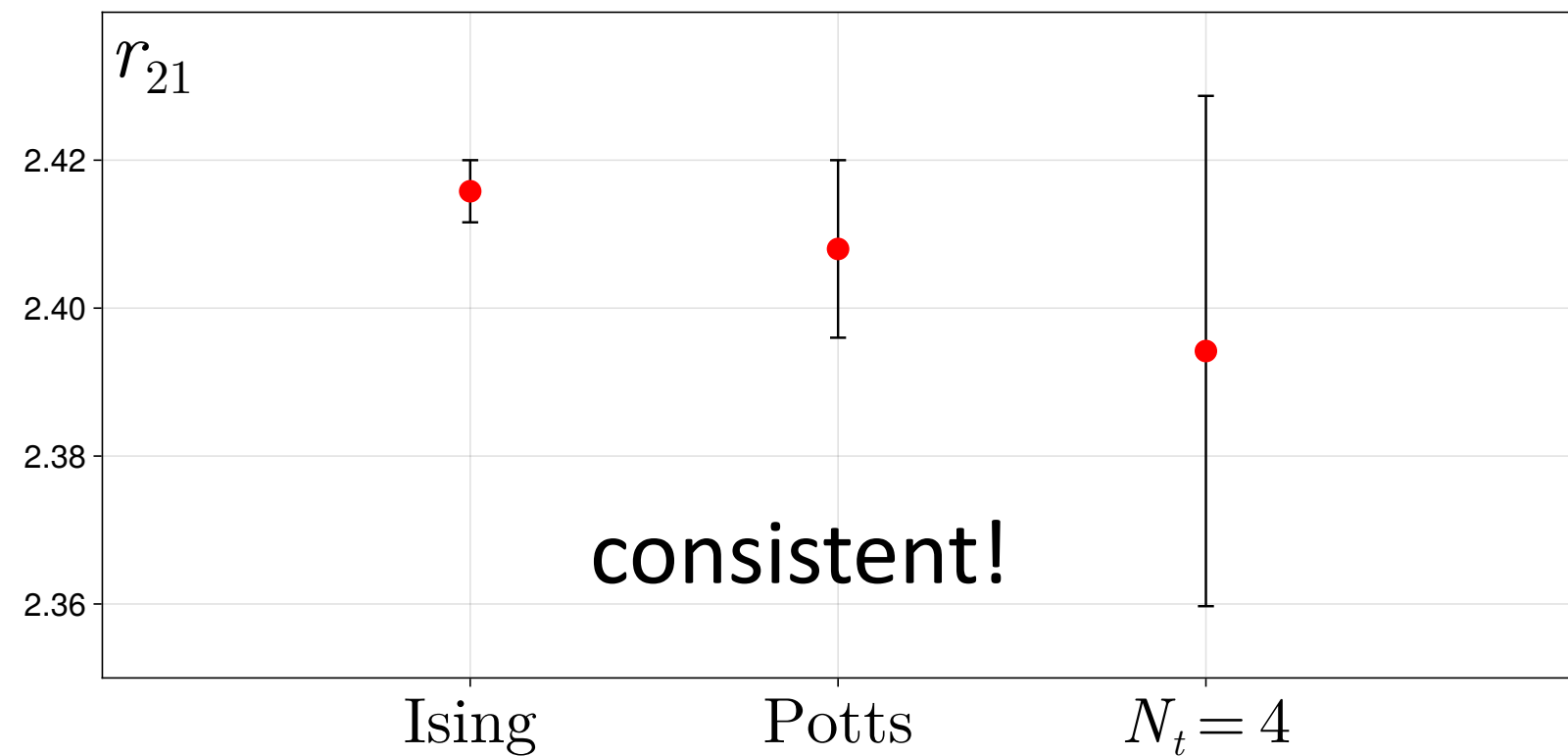
Kiyohara, et al. (2021)



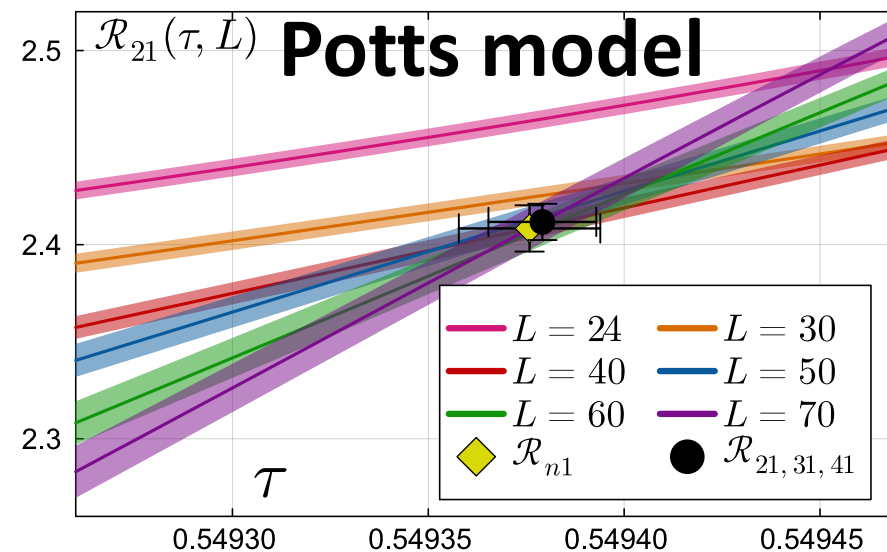
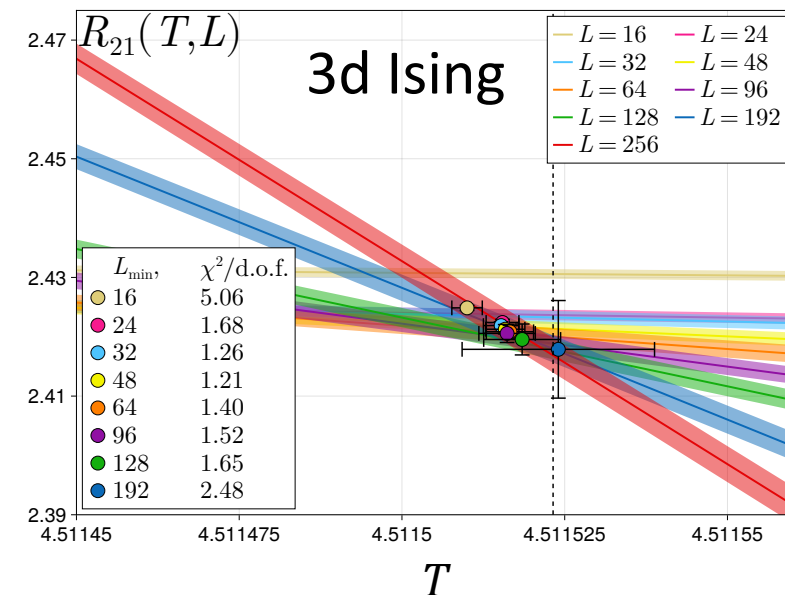
Consistent Result!

Comparing the r_{21}

TW, Kitazawa, Kanaya (2025)

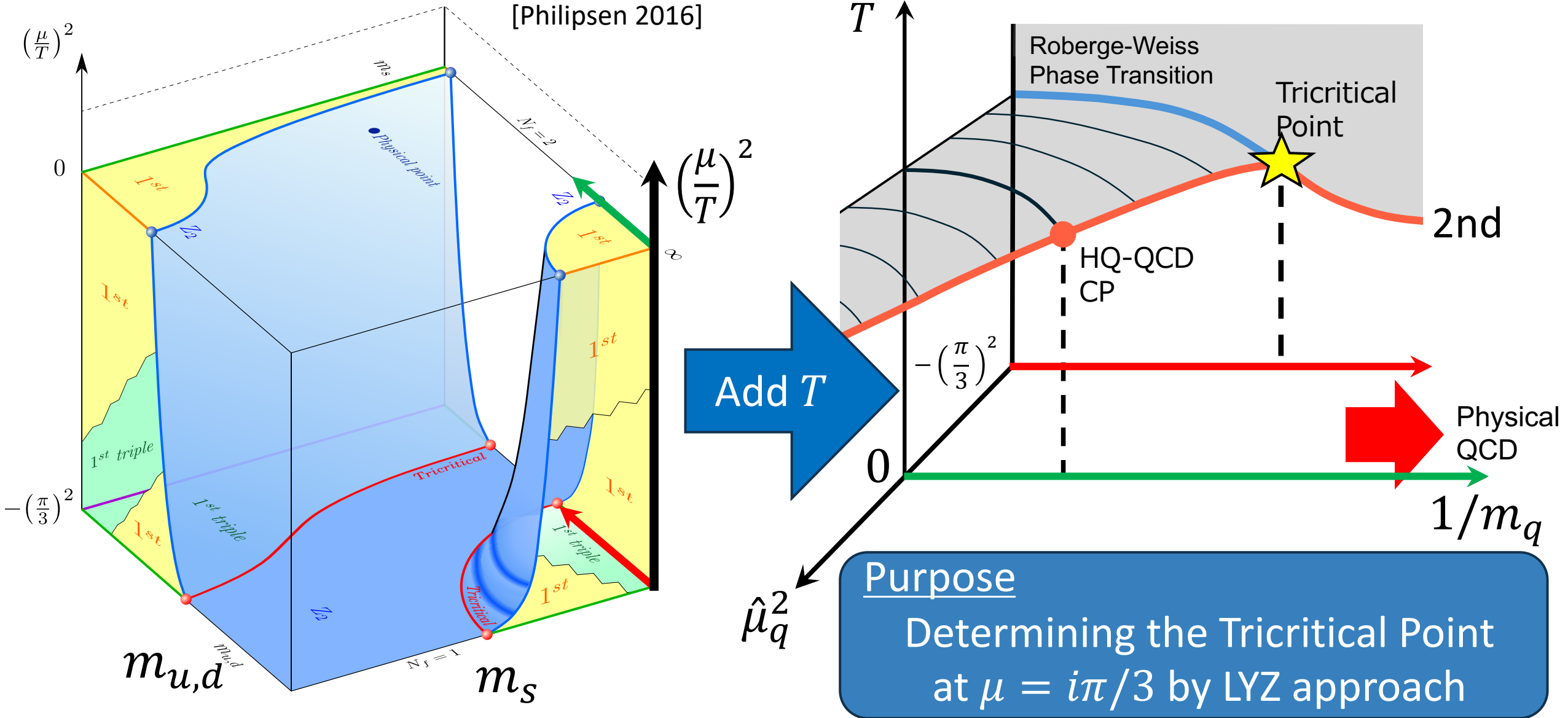


✂ Ising Result is extrapolated to infinite volume.
The others are not.

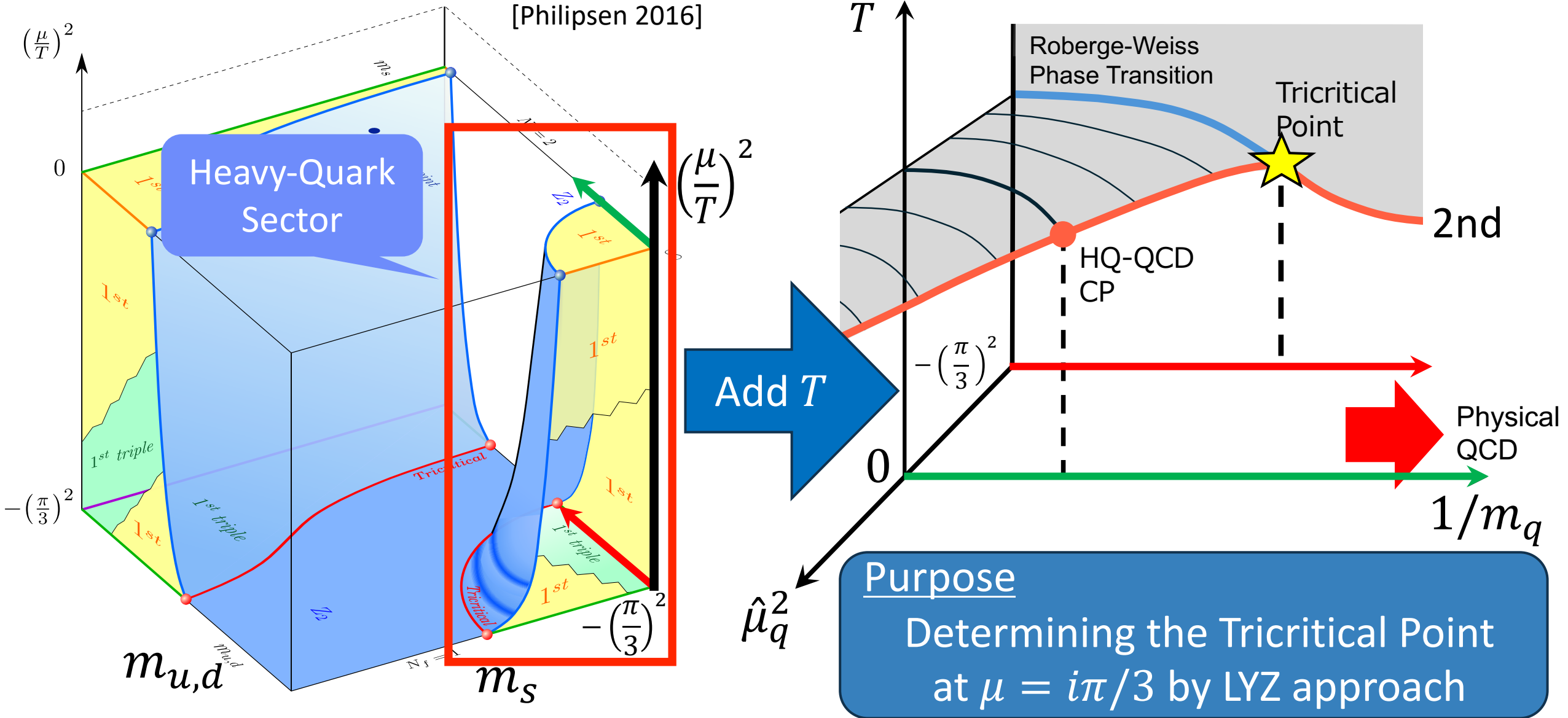


Columbia Plot at non-zero μ

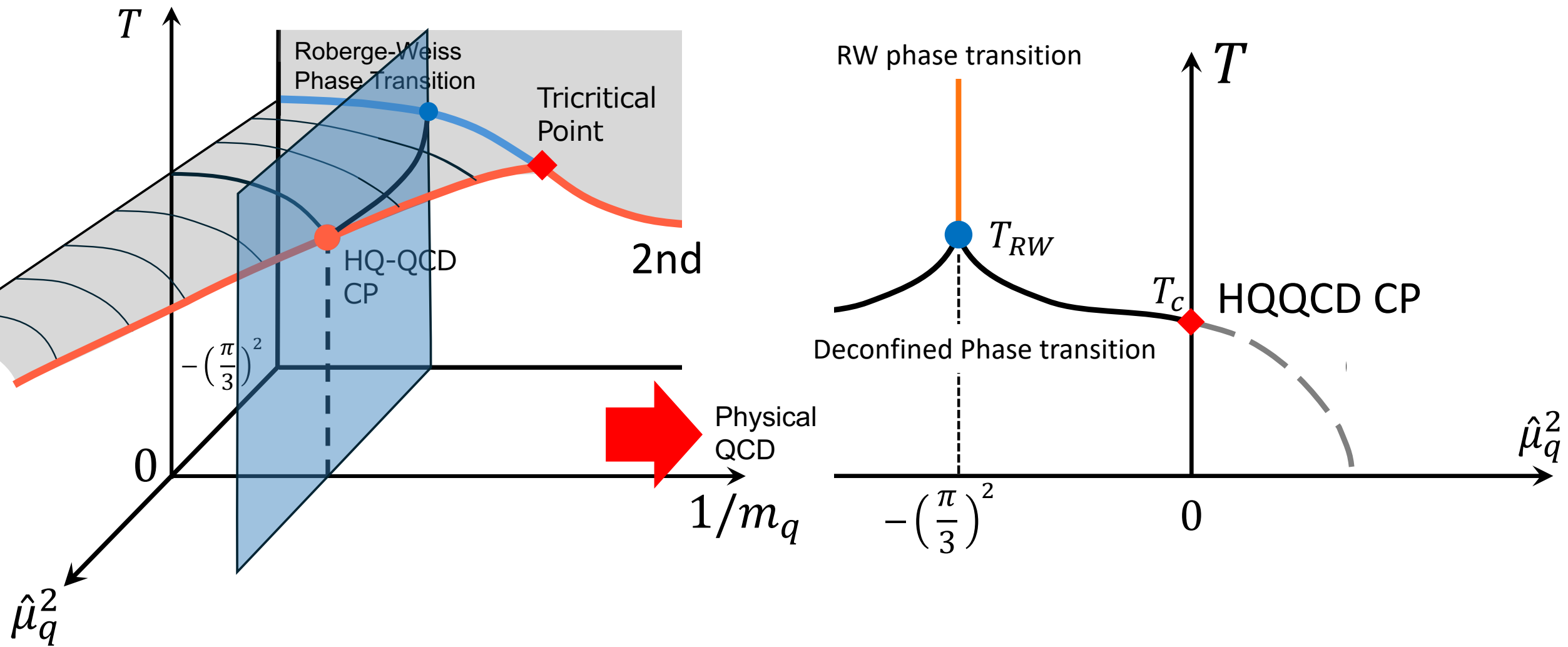
3d-Columbia Plot



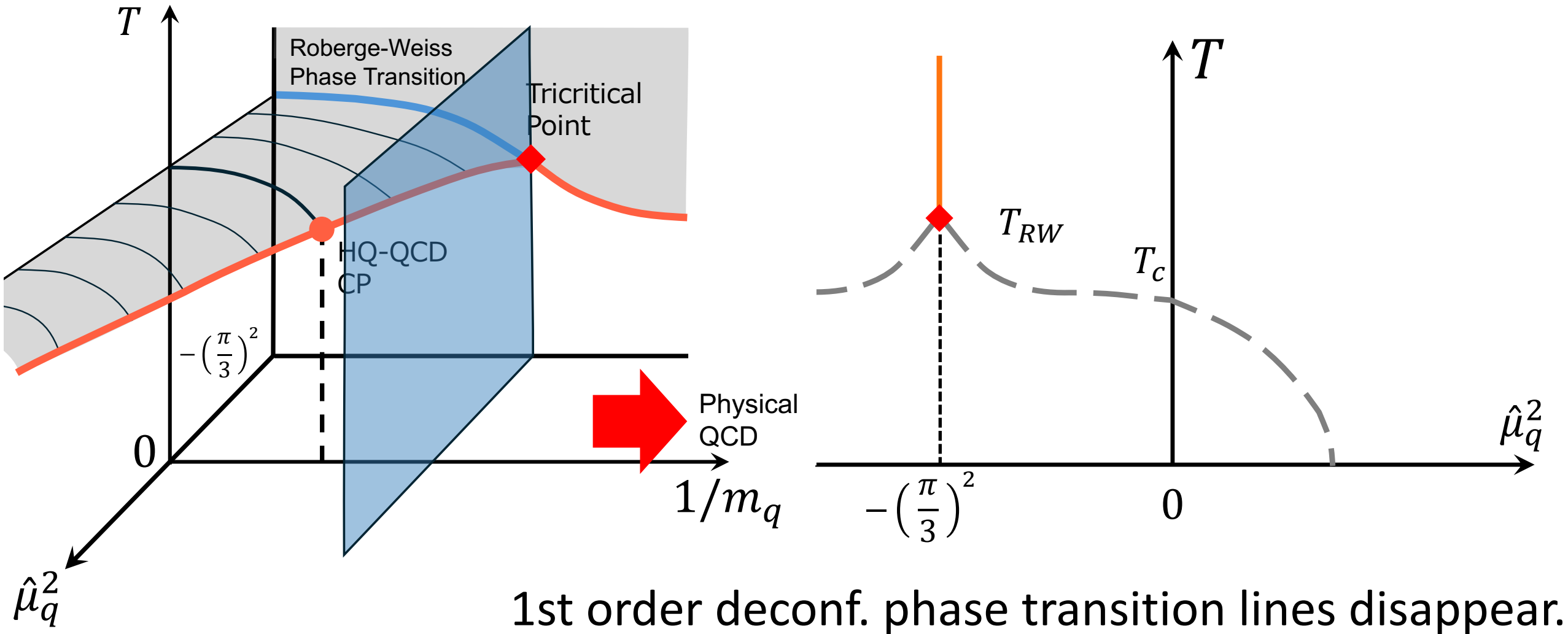
3d-Columbia Plot



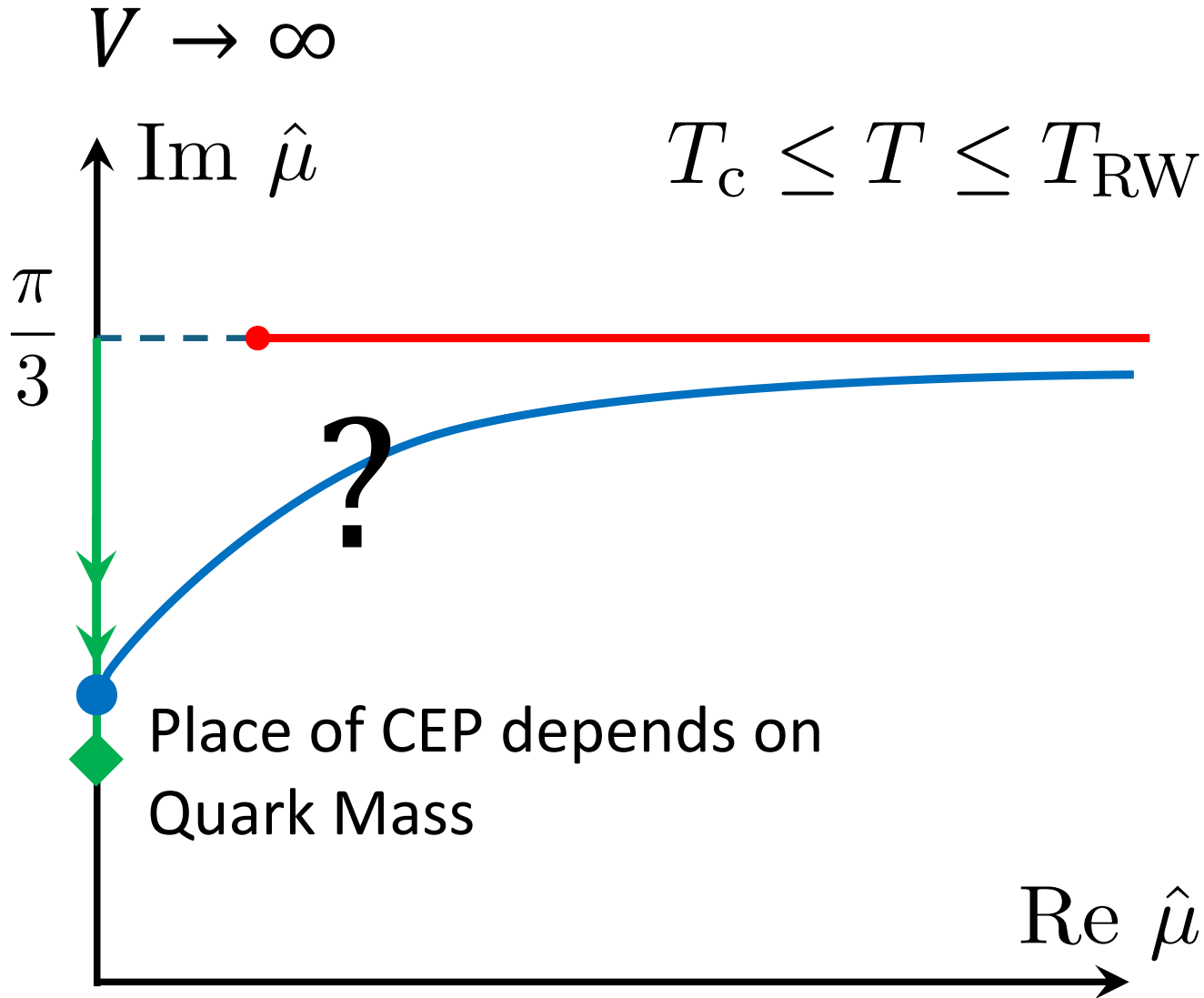
Columbia Plot at non-zero μ



Columbia Plot at non-zero μ



Behavior of LYZ



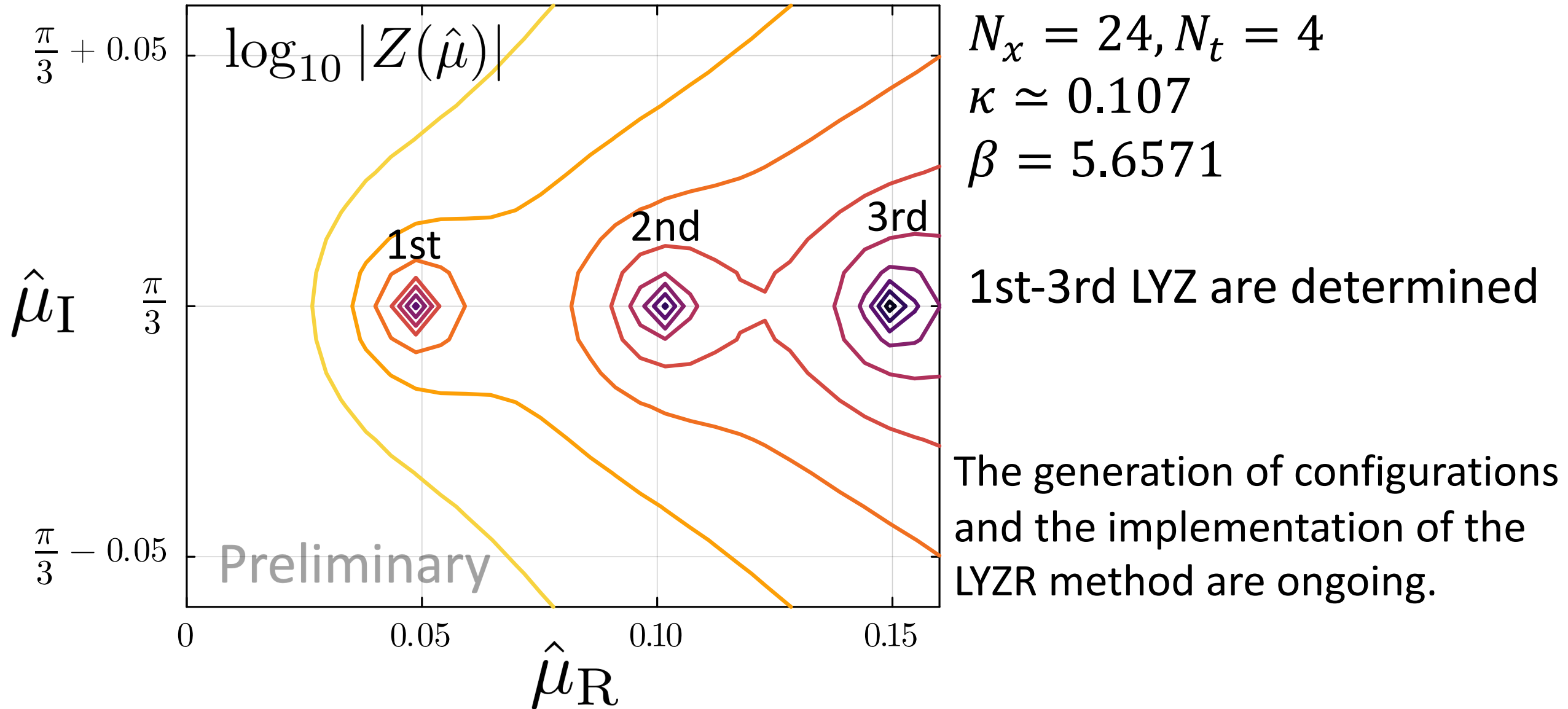
Series 1 —————
 Roberge Weiss Phase transition

Red line extends to the Re-axis.
 Due to the Symmetry of Polyakov loop,
 LYZ exists only at $\mu_I = \pi/3$.

Series 2 —————
 Deconfined Phase transition

The edge of the LYZ are predictable,
 the zero lines extend in $\mathbb{C}$ -plane
 remains unclear

Contour of Partition Function



Summary

- ✓ We detect the critical point in Heavy-Quark QCD using Lee-Yang zero Ratio method. The location of CP is consistent with Binder cumulant method.
- ✓ We conjecture the behavior of Lee-Yang zeros in the non-zero chemical potential.
- ✓ We make the configuration at $\mu = i\pi/3$ and detect the LYZ from Roberge Weiss Phase transition.

Future Works

- ❑ Calculating the Lee-Yang zero Ratio at $\mu = i\pi/3$ and detect the tricritical point in the Heavy-Quark region.
- ❑ Detect Critical end point in the imaginary and real μ
→ They give insight into QCD at m_{phys}
- ❑ Final Goal = Application to 2+1 flavor QCD
→ QCD critical point (We need 2nd LYZ)