

# Millisecond proto-magnetars as gamma-ray burst central engines

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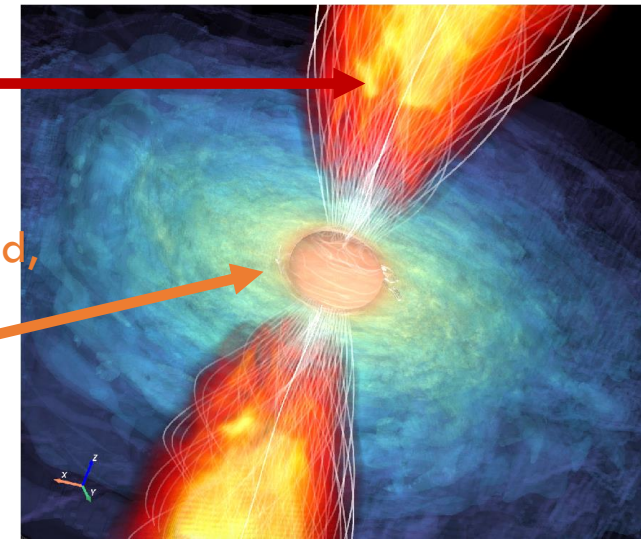


# Gamma-ray bursts (GRBs)

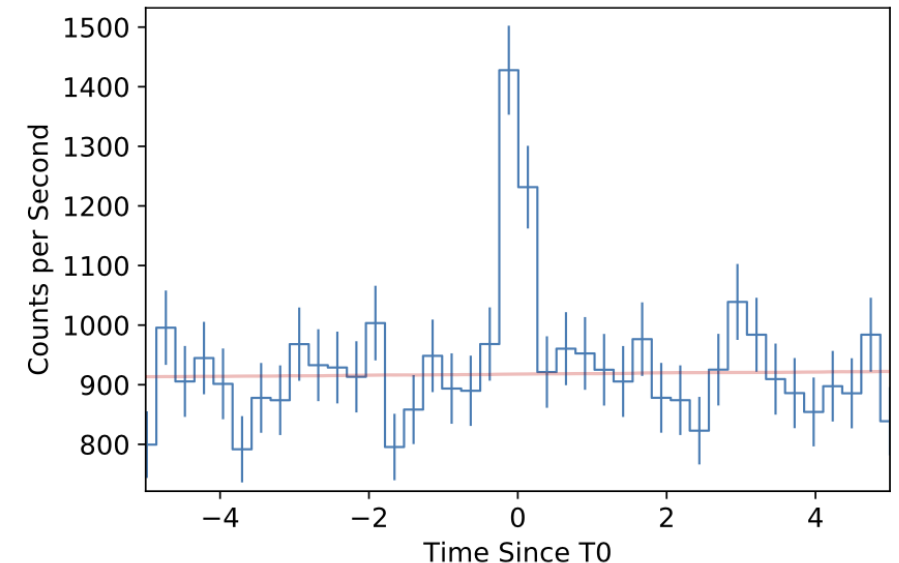
- Seconds-long duration gamma-ray emission observable, e.g.
- Powered by dissipation of energy within a collimated relativistic jet
- Astrophysical sites:
  - Central engine – proto-magnetar or black hole
    - Remnants of compact object mergers or collapsed massive star cores
- Often discovered alongside other signals

Rapidly rotating,  
strongly magnetized,  
hot neutron star

Jet



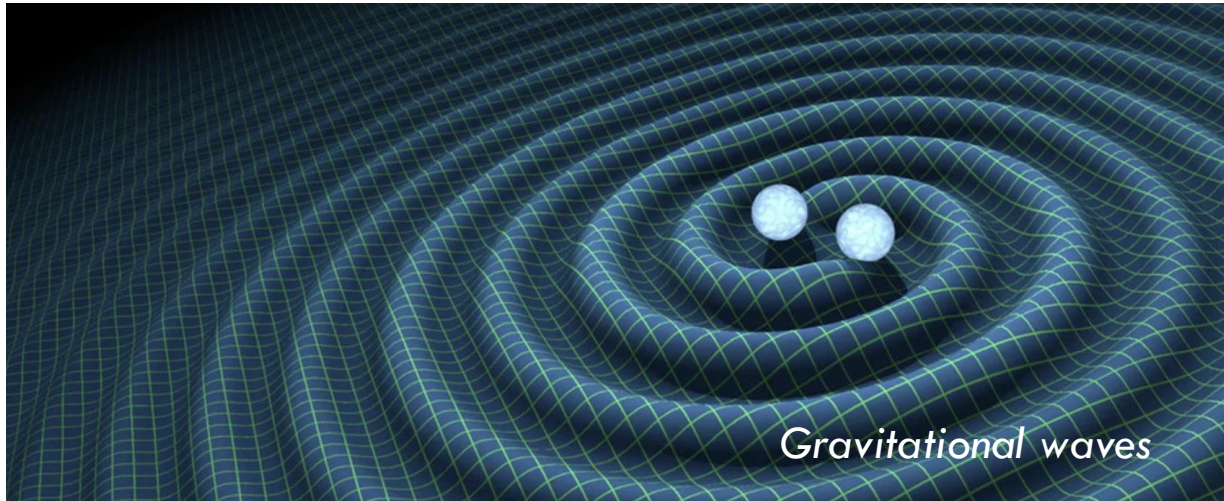
Desai+2026 (submitted,  
ApJL)



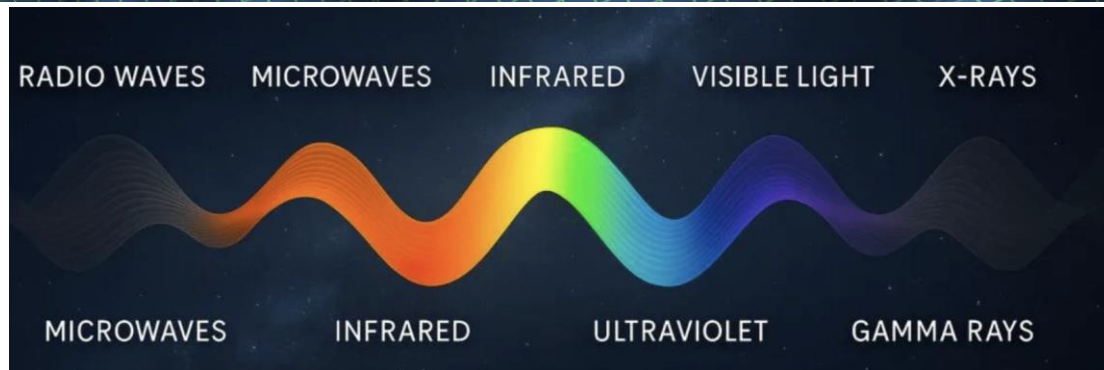
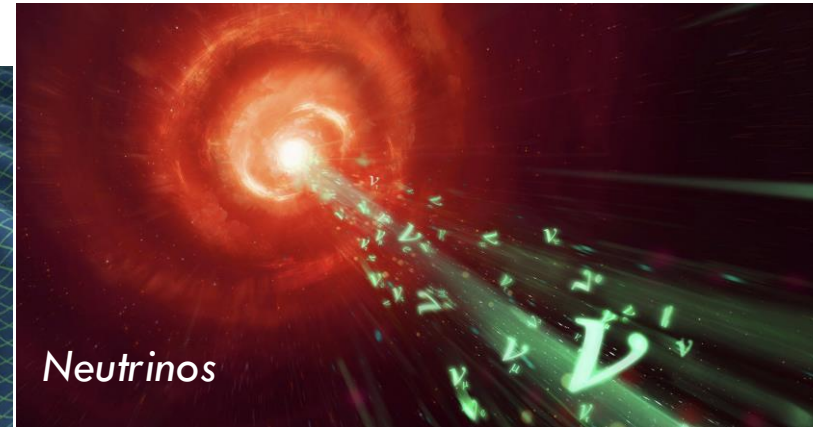
Goldstein+2017

# GRBs: one of many *multi-messenger* observations...

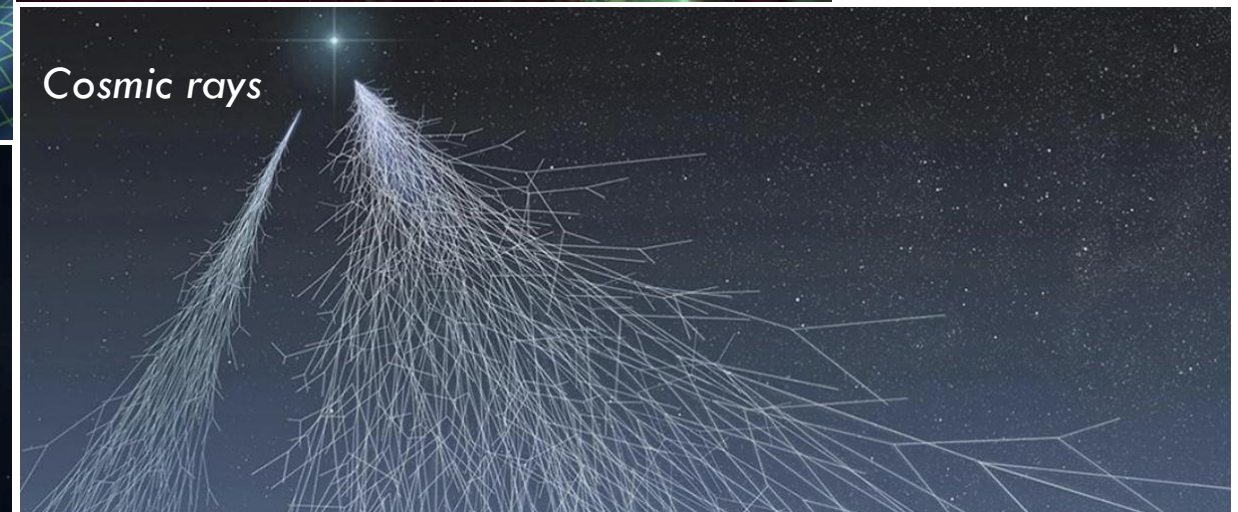
Caltech



Maciej Rebisz, Quanta Magazine



solarblog.pl



ASPERA

Highlight: short duration GRB associated with GW170817

# GRBs: one of many *multi-messenger* observations...

...which help us probe fundamental physics

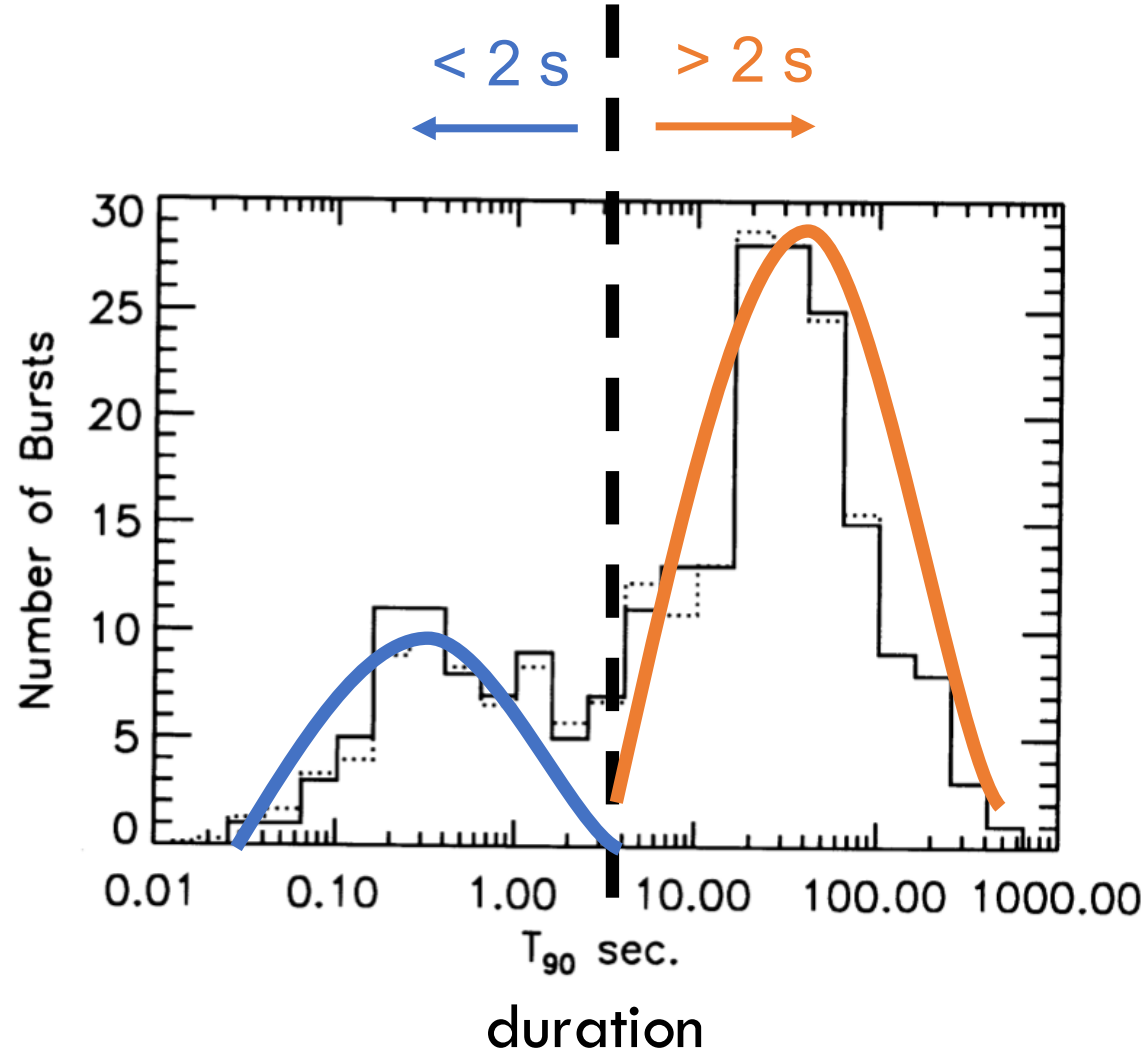
- Sites of r-process element production (among the heaviest elements in the periodic table)
  - Properties of neutron-rich matter
- Physics of neutrinos (e.g. flavor conversions)
- Tests of general relativity

Highlight: short duration GRB associated with GW170817

# GRBs historically classified by duration: long vs. short

- Short GRBs

mergers



- Long GRBs

core-collapse  
supernovae

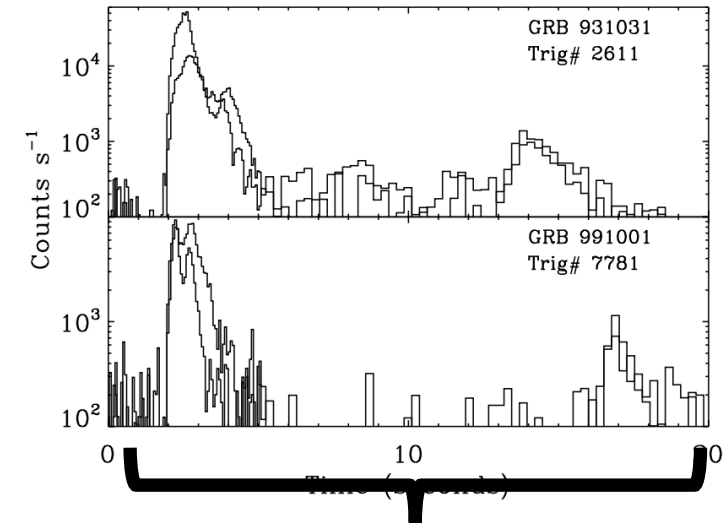
Duration thought to  
be set by accretion  
time-scales/disk  
size

# The line between long vs. short may be fuzzier than we thought

Though 1 Short GRB is linked directly to neutron star merger (GW170817),

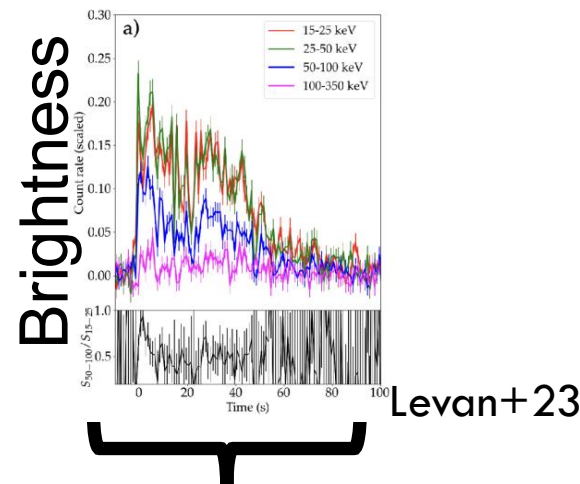
**An Ordinary Short Gamma-Ray Burst with Extraordinary Implications:  
*Fermi*-GBM Detection of GRB 170817A**

recent long GRBs linked to kilonovae (associated with neutron star mergers)



20 s

Short GRB +  
Extended emission



Levan+23

60 s

# Big open questions

- What are the dominant astrophysical sites responsible for the production of the heaviest elements?
- What are the central engines of gamma-ray bursts? How can they power a jet?

# Central engine: proto-magnetar or black hole?

## Black hole

powered by accretion

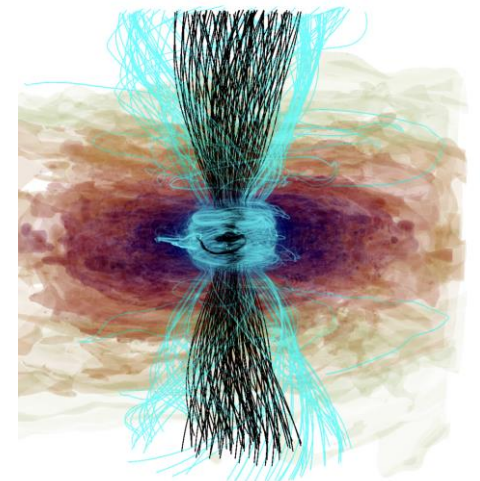
Operates via Blanford-Znajek mechanism



Nasa

## Proto-magnetar

Powered by its rotational energy



# Central engine: proto-magnetar or black hole?

## Black hole

powered by accretion

Operates via Blanford-Znajek mechanism

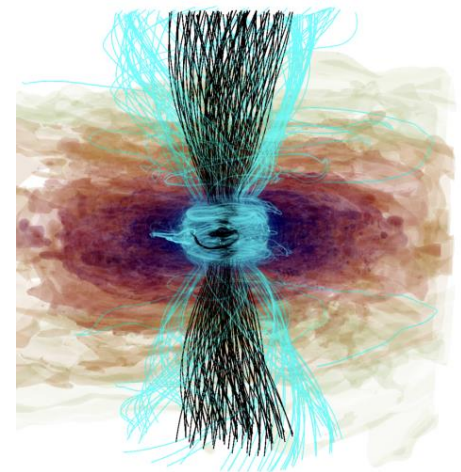


Nasa

## Proto-magnetar

Powered by its rotational energy

Favored because BH cannot baryon load polar jet



# Central engine: proto-magnetar or black hole?

## Black hole

powered by accretion

Operates via Blanford-Znajek mechanism

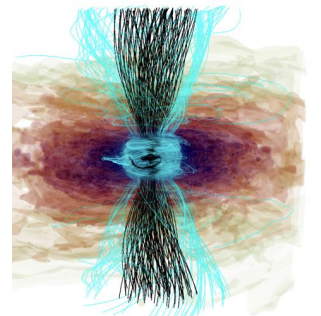


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## Proto-magnetar

Powered by its rotational energy

I focus on viability of **proto-magnetars** as central engines



Central question of our latest work:

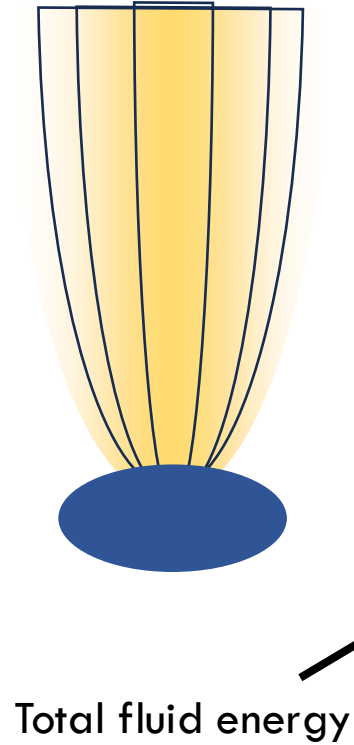
# **Can a proto-magnetar power a jet?**

i.e., can it produce collimated, relativistic  
outflows?

Does such a steady-state solution exist?

# Jets have large inferred Lorentz factors ( $>10$ ). How?

$$\Gamma Mc^2 = E_{\text{kin}} + E_{\text{mag}} + E_{\text{therm}} + Mc^2$$



This is the energy budget for outflows

# Jets have large inferred Lorentz factors ( $>10$ ). How?

$$\Gamma_{\infty} Mc^2 = E_{\text{kin}} + E_{\text{mag}} + E_{\text{therm}} + Mc^2$$

If magnetically dominated...

$$E_{\text{mag}}/Mc^2 = \sigma \equiv b^2/4\pi\rho c^2$$

$$\Gamma_{\infty} = \sigma + 1 \approx \sigma$$

Asymptotic Lorentz factor

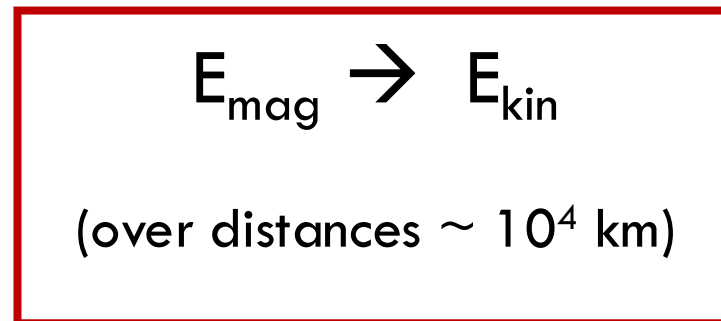
# Jets have large inferred Lorentz factors ( $>10$ ). How?

$$\Gamma_{\infty} Mc^2 = E_{\text{kin}} + E_{\text{mag}} + E_{\text{therm}} + Mc^2$$

If magnetically dominated...

$$\Gamma_{\infty} = \sigma + 1 \approx \sigma$$

Not my  
problem



$$\sigma \equiv b^2 / 4\pi\rho c^2$$

**Magnetization is key**

reformulated:

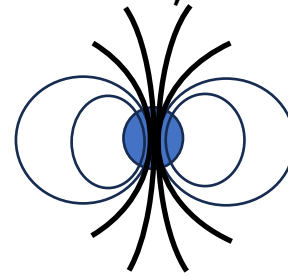
**Can proto-magnetars have  
magnetically-dominated (high -  $\sigma$ )  
outflows?**

# Previous work + open questions

- **Local/global simulations:** How do proto-magnetar properties develop self-consistently?

e.g. Mösta+20, Kuroda+20, Combi+2023, Hayashi+2024, Kiuchi+2024, Kalinani+2025, Gutierrez+2026...

How to form strong, large-scale, ordered B fields???

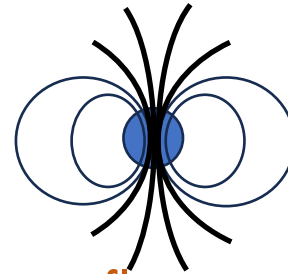


# Previous work + open questions

- **Local/global simulations:** How do proto-magnetar properties develop self-consistently?

e.g. Mösta+20, Kuroda+20, Combi+2023, Hayashi+2024, Kiuchi+2024, Kalinani+2025...

Assume we have strong, large-scale, ordered B fields



- **Proto-magnetar outflows as central engines for GRBs? Are outflow properties compatible?** e.g., Bucciantini+06, Metzger+07, +08, Prasanna+2024...

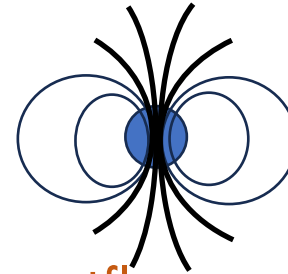
**They are 1D or assume neutron star boundary conditions**

# Previous work + open questions

- **Local/global simulations:** How do proto-magnetar properties develop self-consistently?

e.g. Mösta+20, Kuroda+20, Combi+2023, Hayashi+2024, Kiuchi+2024, Kalinani+2025...

Assume we have strong, large-scale, ordered B fields



- **Proto-magnetar outflows as central engines for GRBs? Are outflow properties compatible?** e.g., Bucciantini+2006, Metzger+2007, +2008, Thompson+04, Prasanna+2024...

Assume we have outflow properties on scales  $\sim 1000$  km

- **Large scale jet studies** to connect simulations to observations

eg., Ciolfi+2020, Pavan+2025, Dreas+2026...

# **Ingredients of proto-magnetars**

# 3 key ingredients determine outflow properties

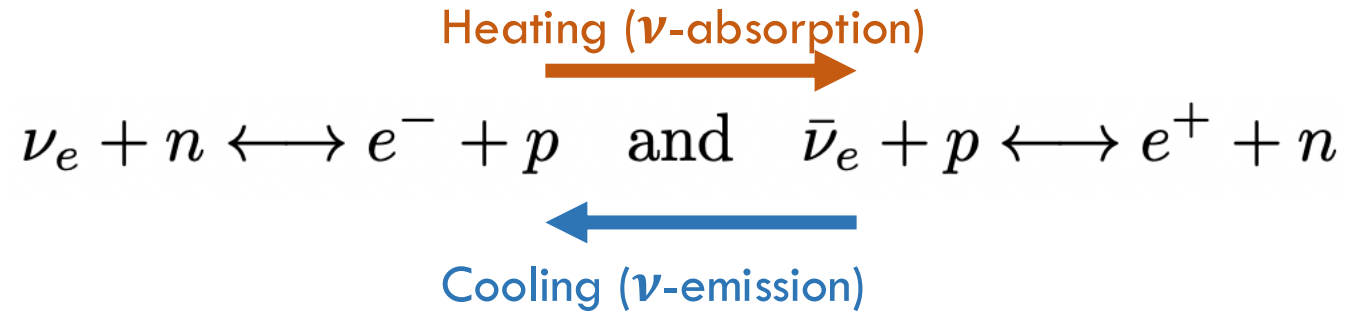
- Net neutrino heating
- Rotation rate (amplified from angular momentum conservation)
- Magnetic field strength/geometry (flux-freezing + dynamos)

# Following a successful explosion or merger...

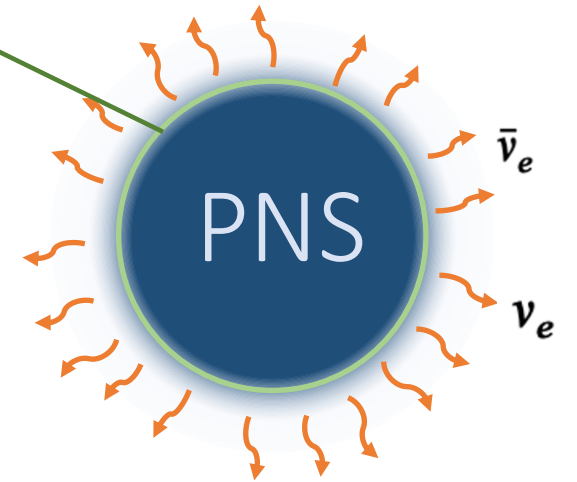
- There may be a phase where accretion disk is gone and other material is ejected
- Left with proto-magnetar winds

# Net neutrino absorption drives baryonic wind

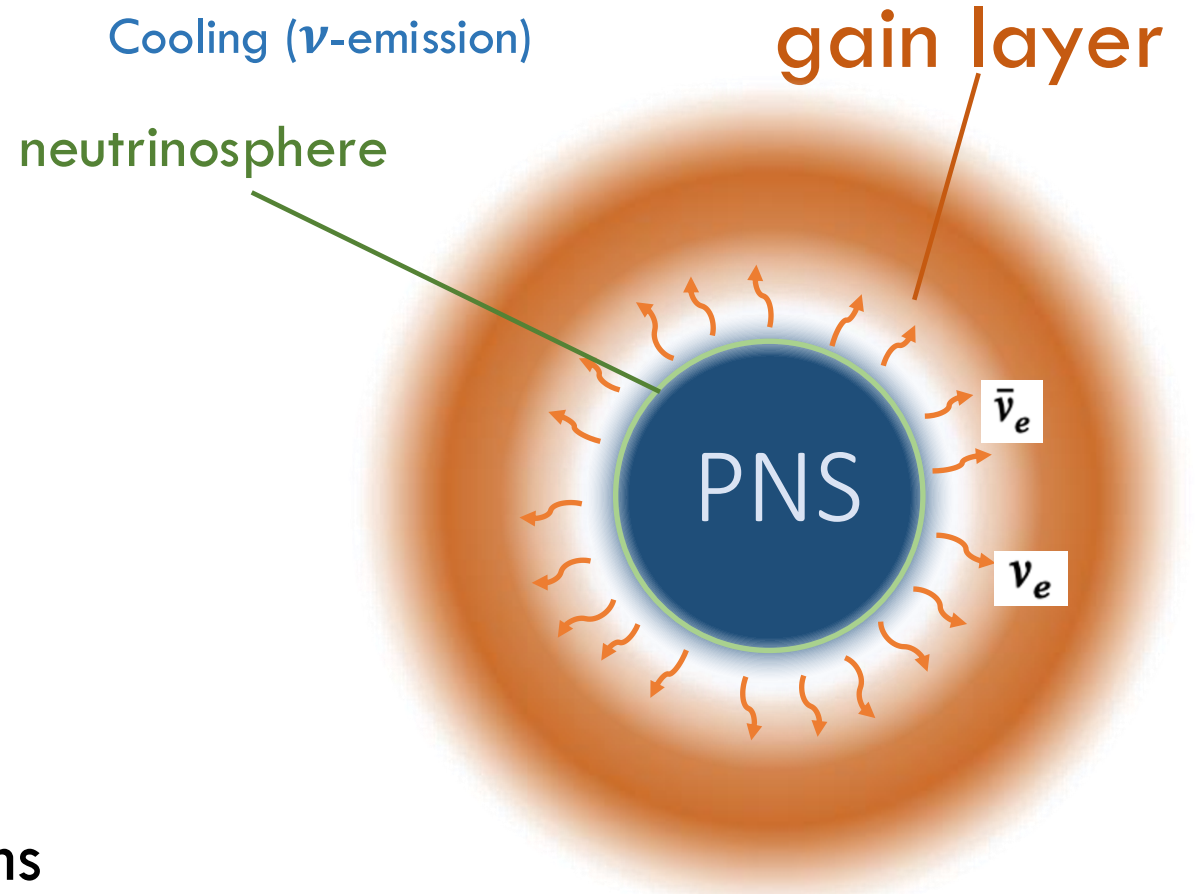
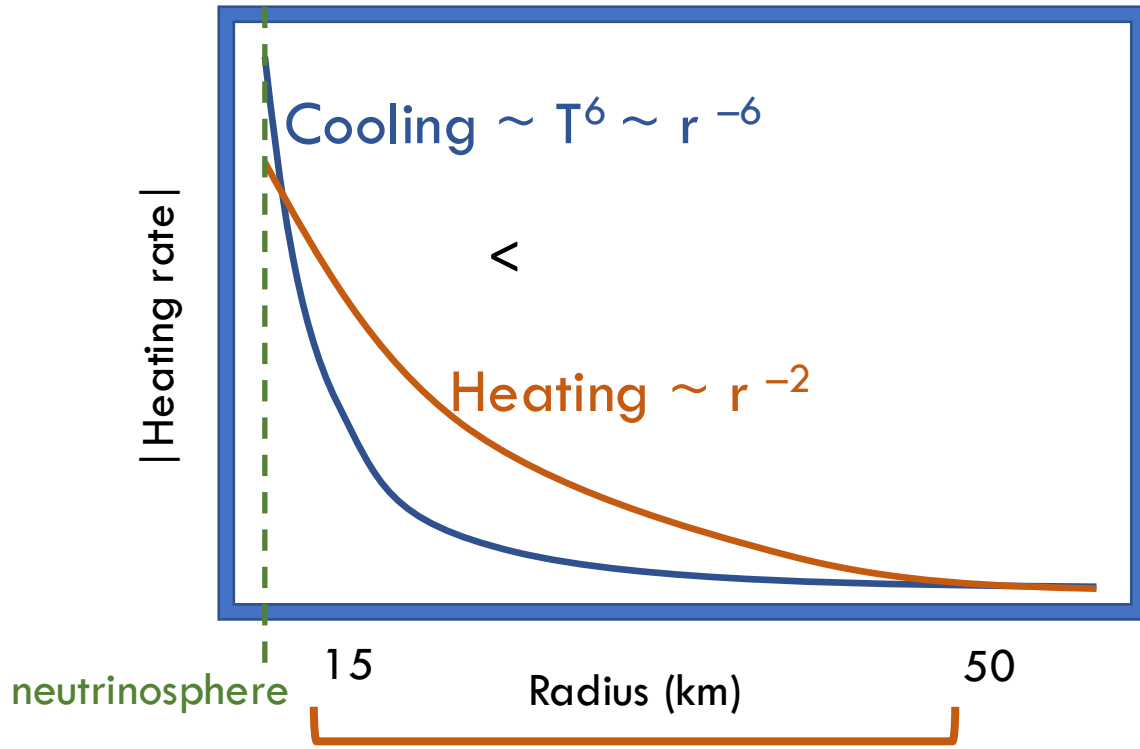
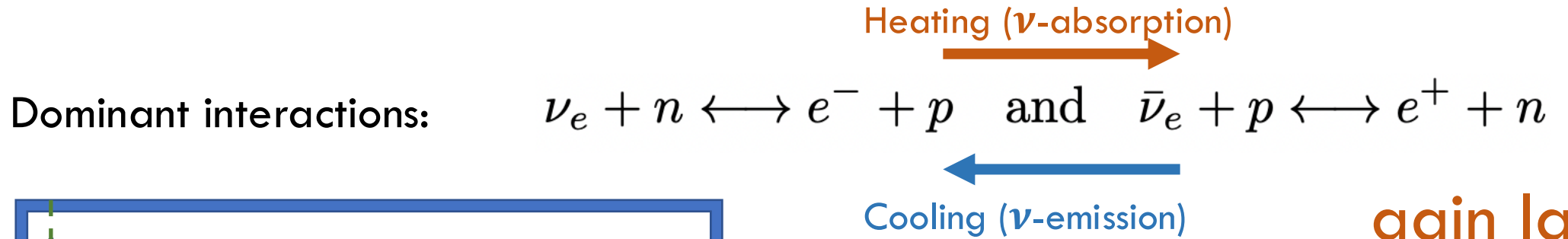
Dominant interactions:  
(charge-current)



neutrinosphere



# Net neutrino absorption drives baryonic wind



gain layer with net positive heating forms

# Net neutrino absorption drives baryonic wind

Dominant interactions:



Heating ( $\nu$ -absorption)



Cooling ( $\nu$ -emission)



gain layer

neutrinosphere

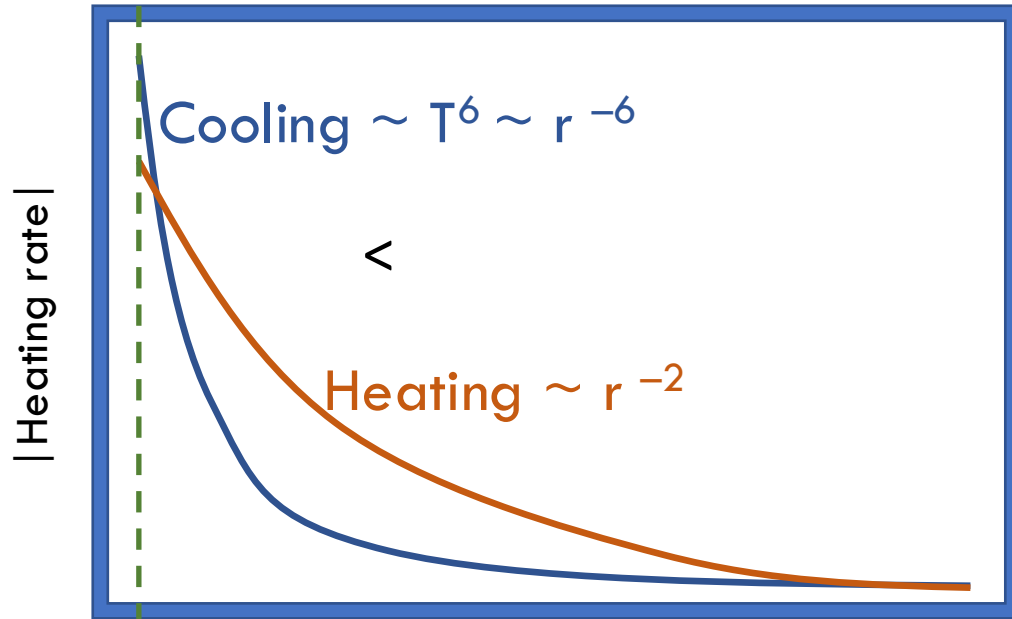
**High neutrino luminosities?**

**High mass loss rate**

PNS

$\bar{\nu}_e$

$\nu_e$



neutrinosphere

15

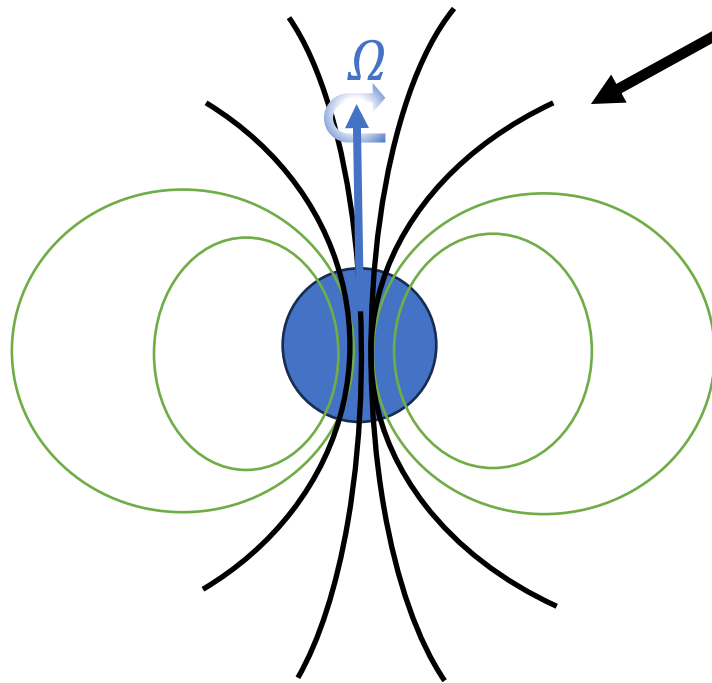
Radius (km)

50

gain layer with net positive heating forms

**magnetic fields and rotation can be amplified within ~100s of ms after collapse**

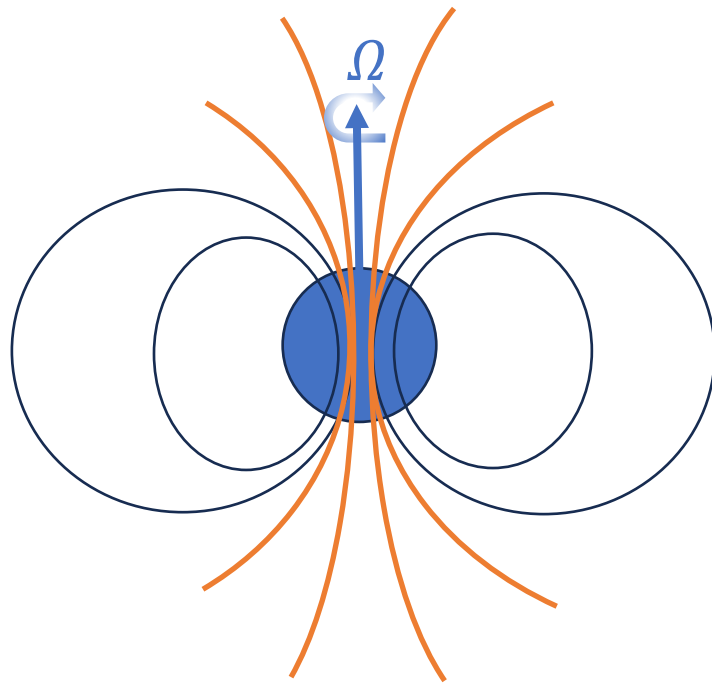
# Neutrino heating loads field lines with baryonic matter



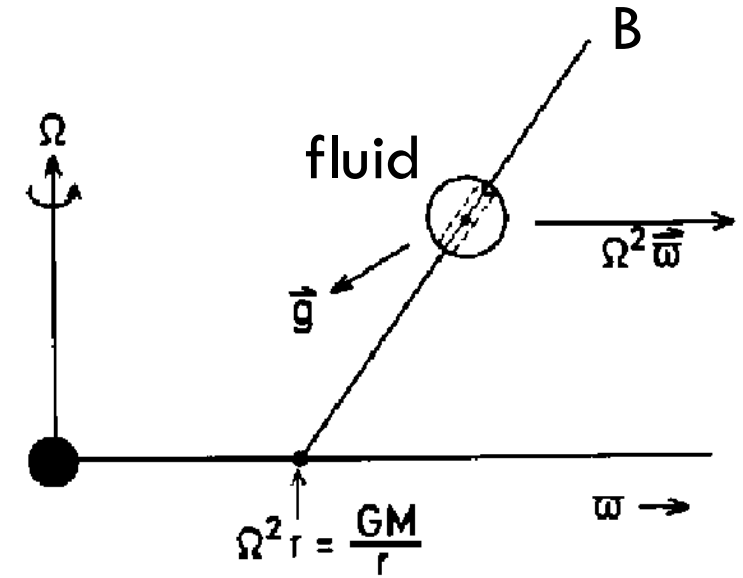
Matter along closed field lines are trapped

Physical intuition: The stronger the magnetic field the more rigid the wires along which fluid travels

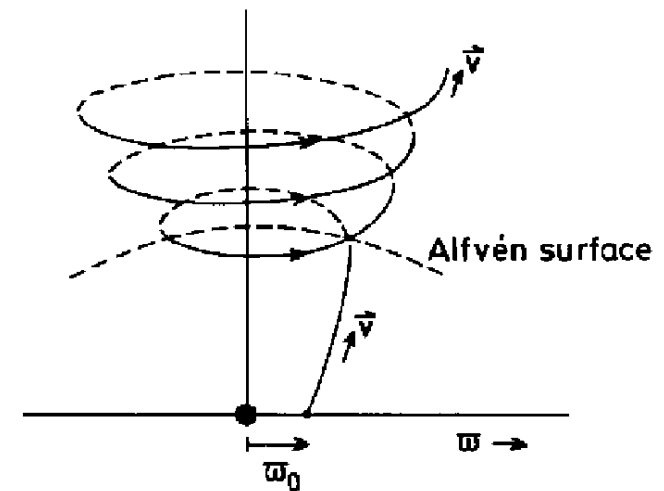
# Rotation + magnetic fields can sling fluid along open field lines



The faster (stronger) the rotation (magnetic field), the greater the acceleration



Spruit 96



# Long-term global evolution ( $\sim$ seconds) of merger/collapse is uncertain/ has diverse outcomes

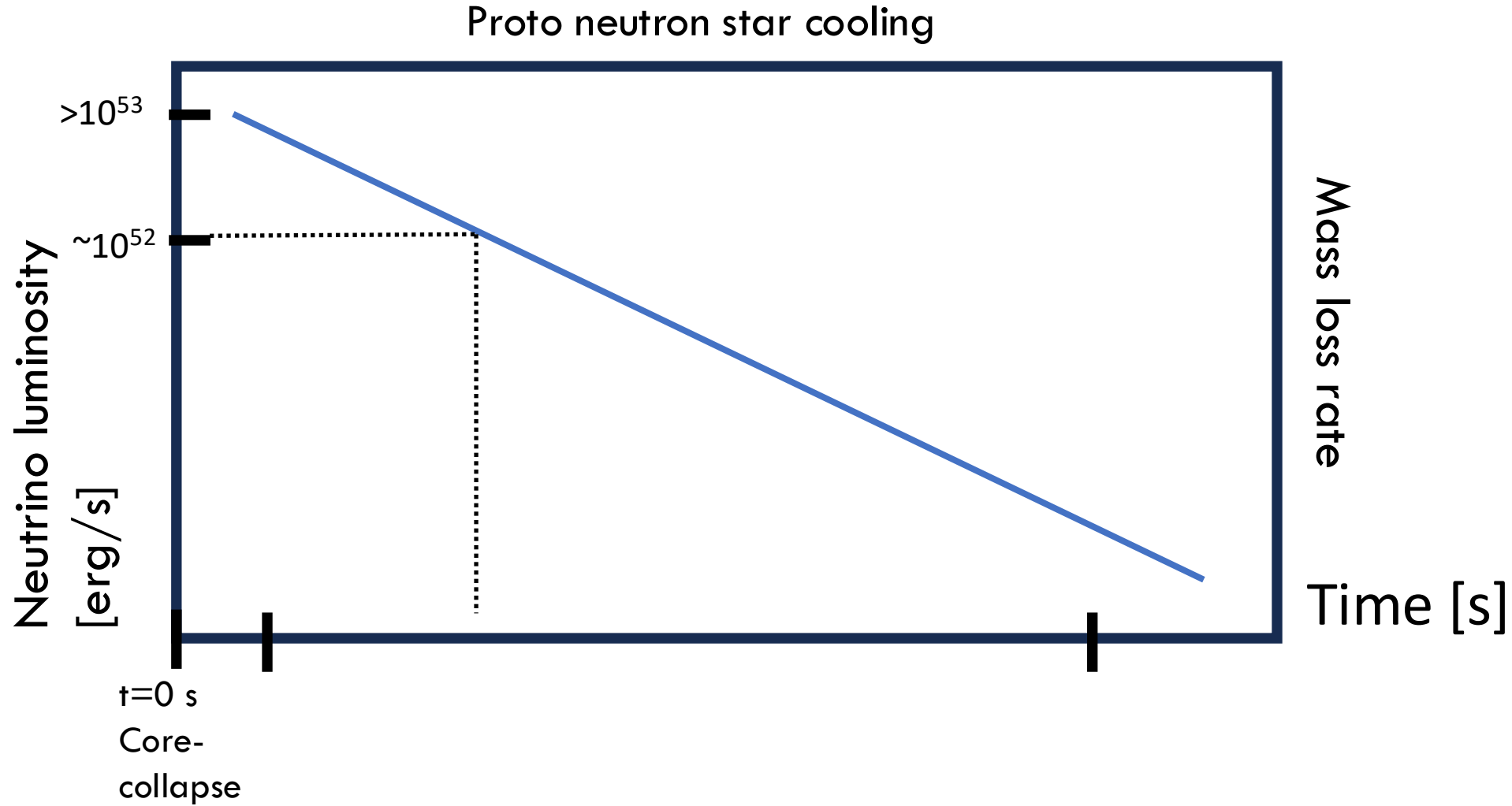
Uncertainties in the physics (e.g. EOS, progenitor models...)

Computationally challenging to simulate with all the key physical ingredients for  $\sim$ seconds

Imagine the following *successful* supernova scenario producing a long-lived proto-magnetar...

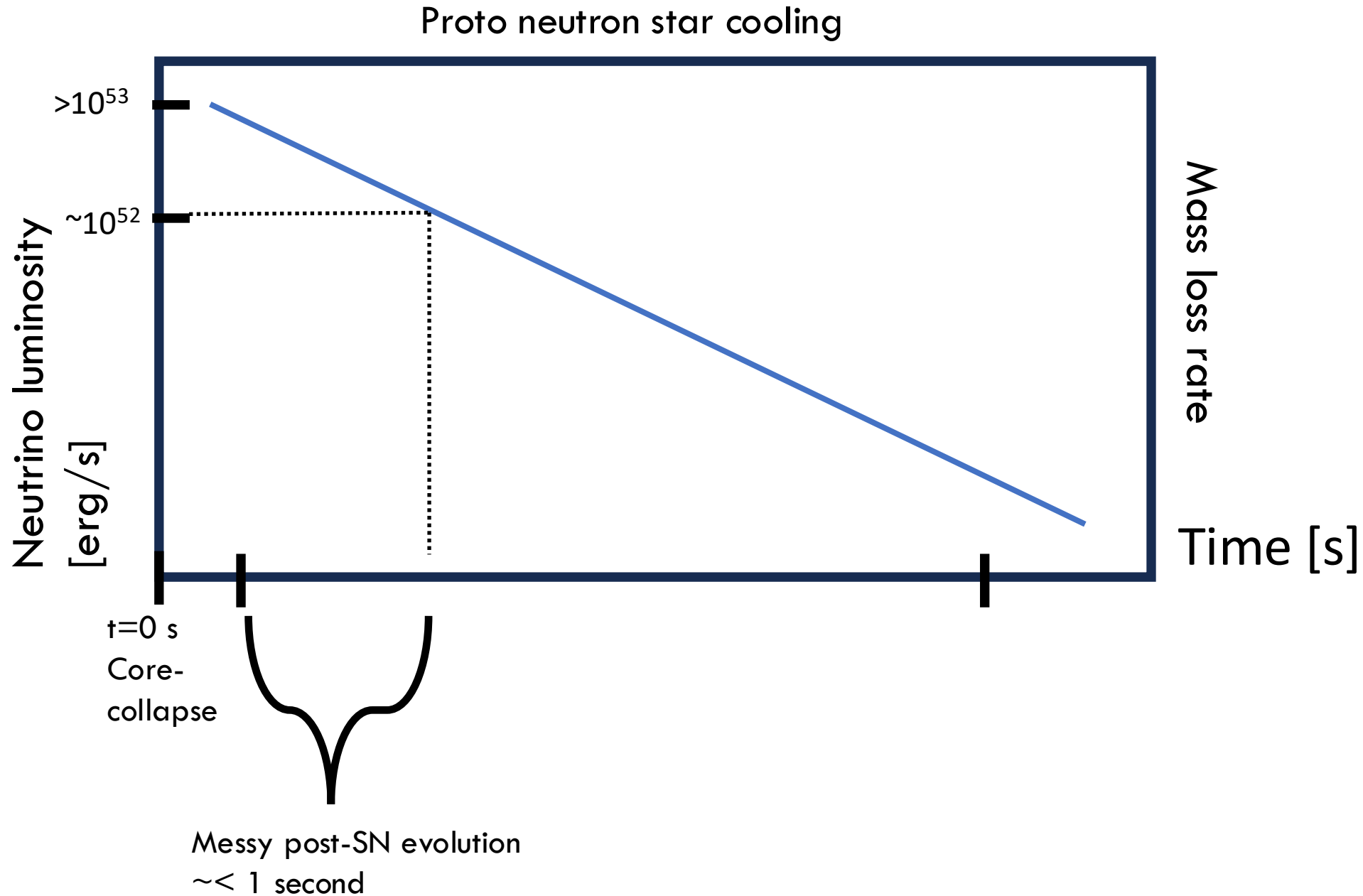
# Cooling time scale is ~seconds

Clean, **baryonic wind** may develop following **successful, energetic supernova explosion** after ~few seconds



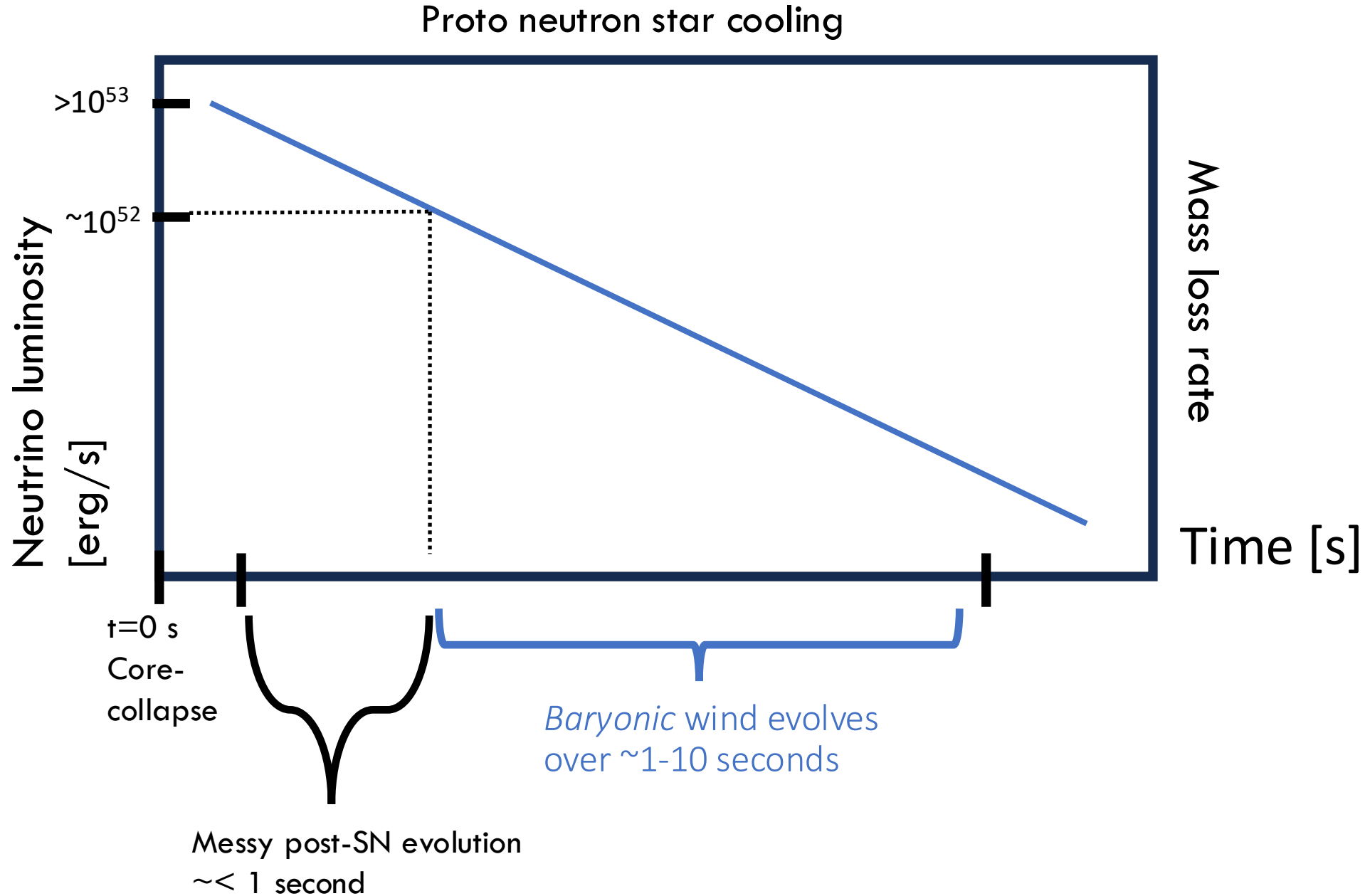
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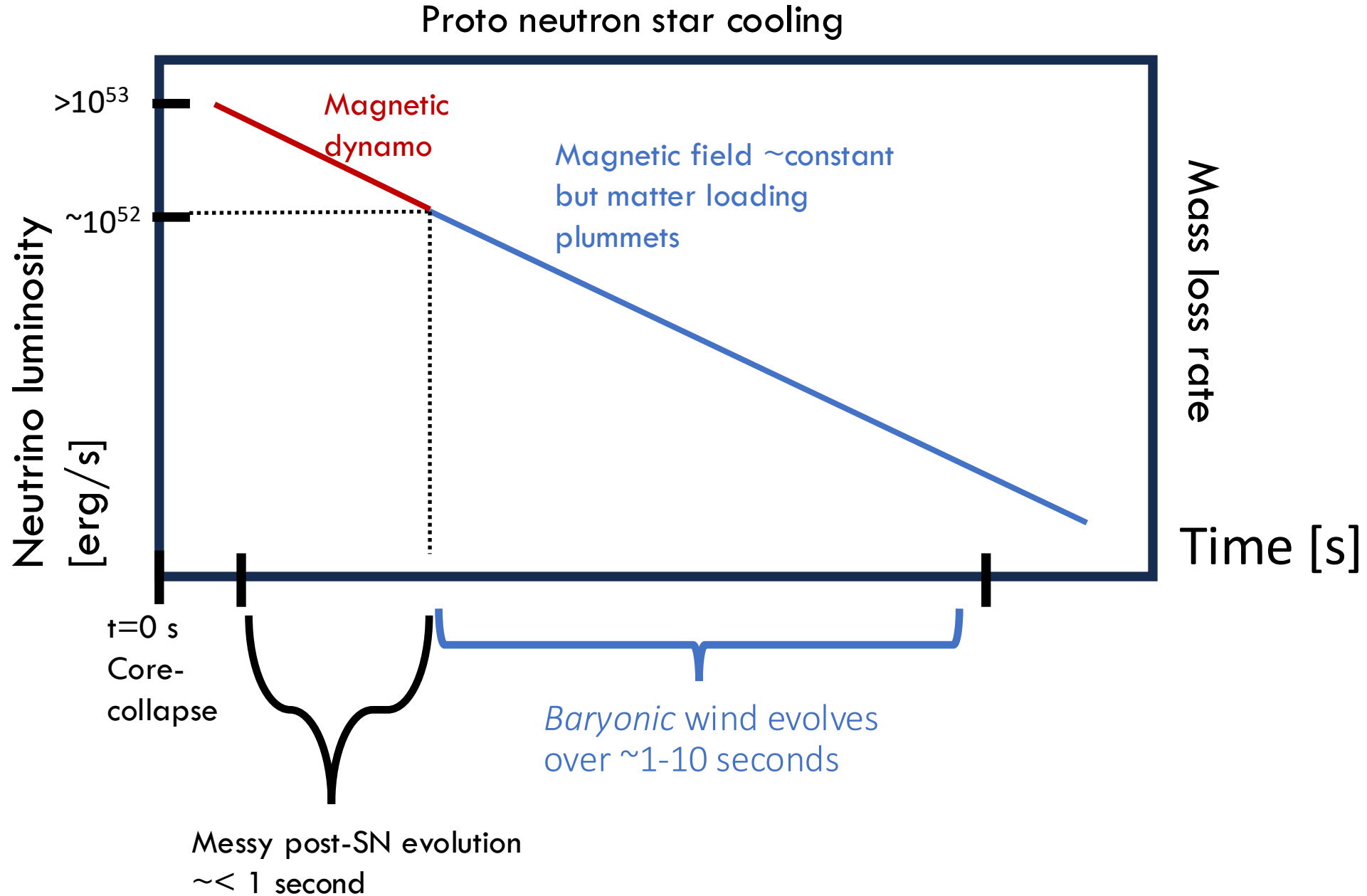
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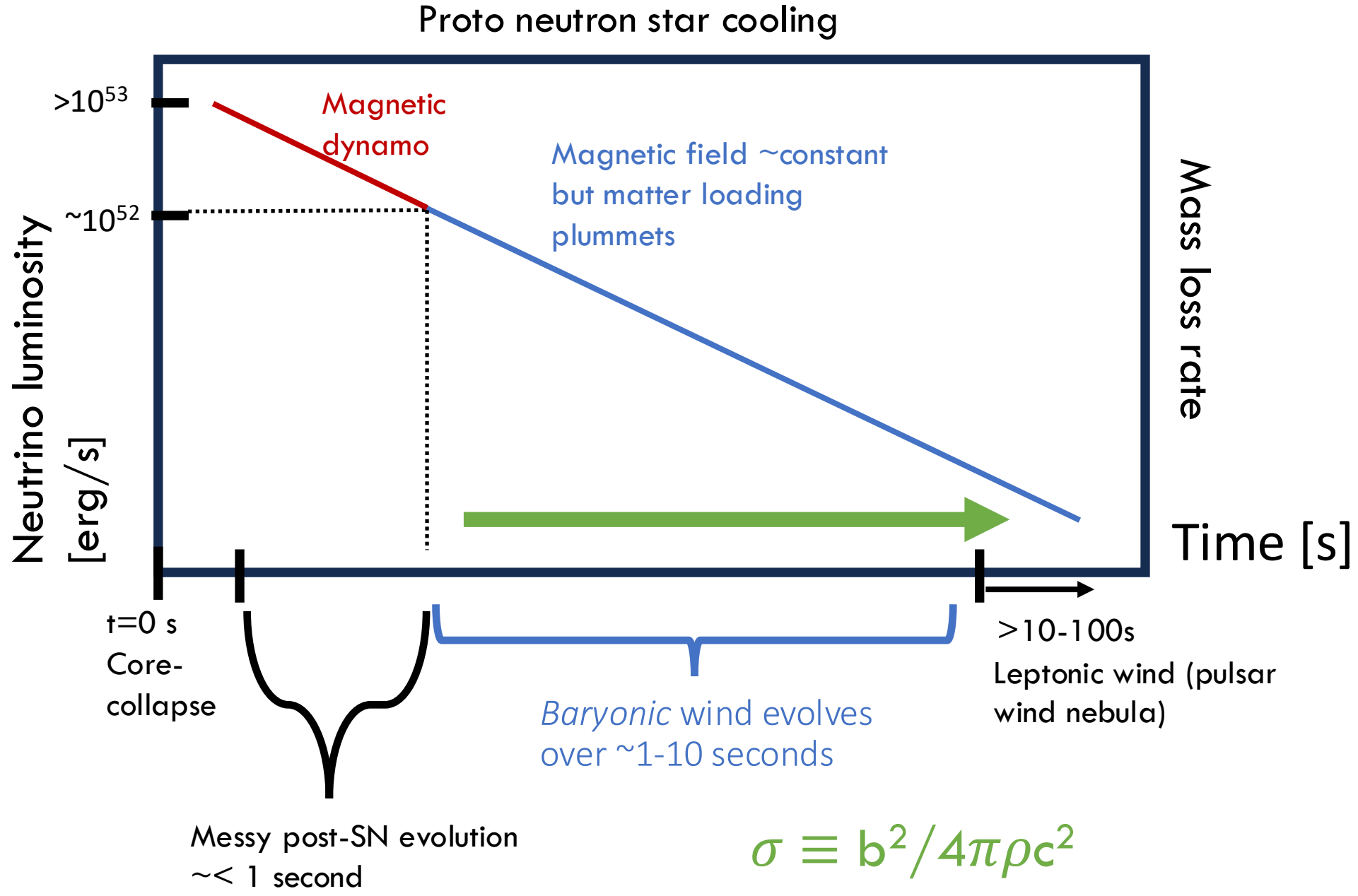
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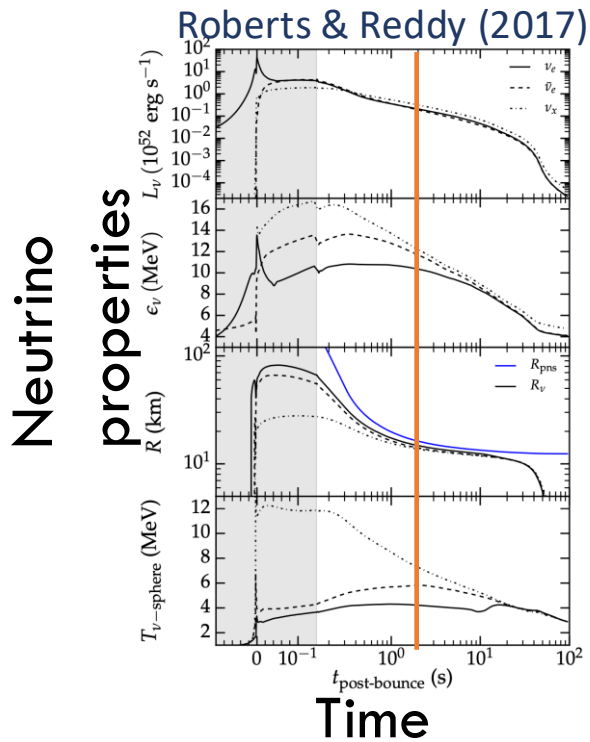


$$\sigma \equiv b^2 / 4\pi\rho c^2$$

grows in time. Can reach ~1000

# Aim: capture steady-state outflows for a variety of configurations

Previously: Isolated effects of rotation, magnetic fields (Desai+22, Desai+23)



+ Magnetic fields

+ rotation



Wind properties,  
eg:

- Mass loss rate
- entropy
- Composition
- magnetization

**The first 3D GRMHD simulations  
of neutrino-heated winds from  
rapidly rotating, strongly  
magnetized proto-magnetars**

# GRMHD simulations: toy models to better understand outflows

proto-magnetar not formed self-consistently from core-collapse simulation

*\*idealized, controlled environment\**

# GRMHD simulations as toy models allow us to...

Be agnostic of progenitor

WD collapse

massive star collapse

merger

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Be agnostic of progenitor

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massive star collapse

merger

Compare to observations - diverse

# GRMHD simulations as toy models allow us to...

Be agnostic of progenitor

WD collapse

massive star collapse

merger

Compare to observations - diverse

Disentangle physical mechanisms

# Our 3D GRMHD toy model

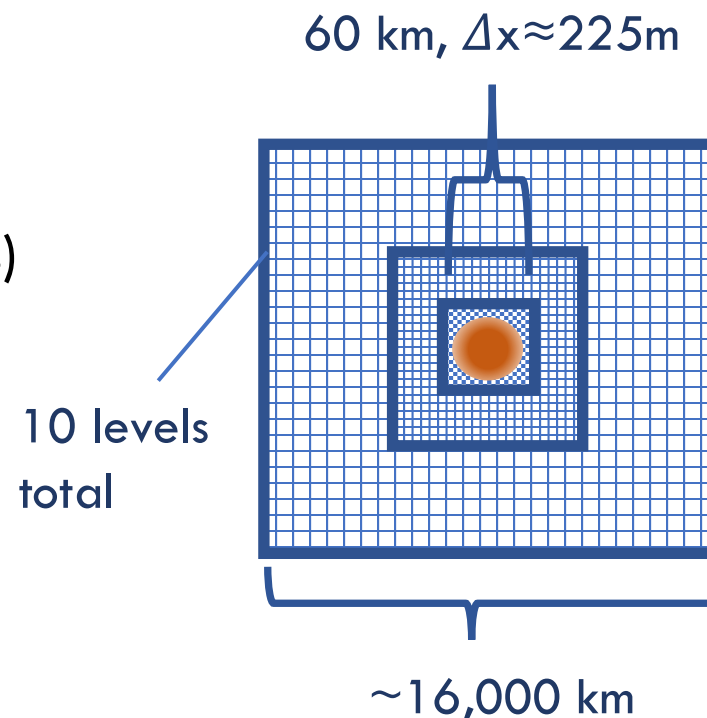


## Evolution code

- GRMHD code (Siegel & Metzger 2018)
  - B-field evolution now using **vector potential formalism**
  - Hydro solver: HLLE + weno-z (adaptive scheme in **high gradient regions, high magnetizations**)
- Tabulated **SFHo** nuclear equation of state (Steiner+2013)
- 1-moment (**MO**) scheme for neutrino transport (Radial equations evolving neutrino mean energies/number densities along radial rays)
  - Emission + absorption (Radice+16, with modifications)
- Rotating models use RNS (Stergioulas+95)

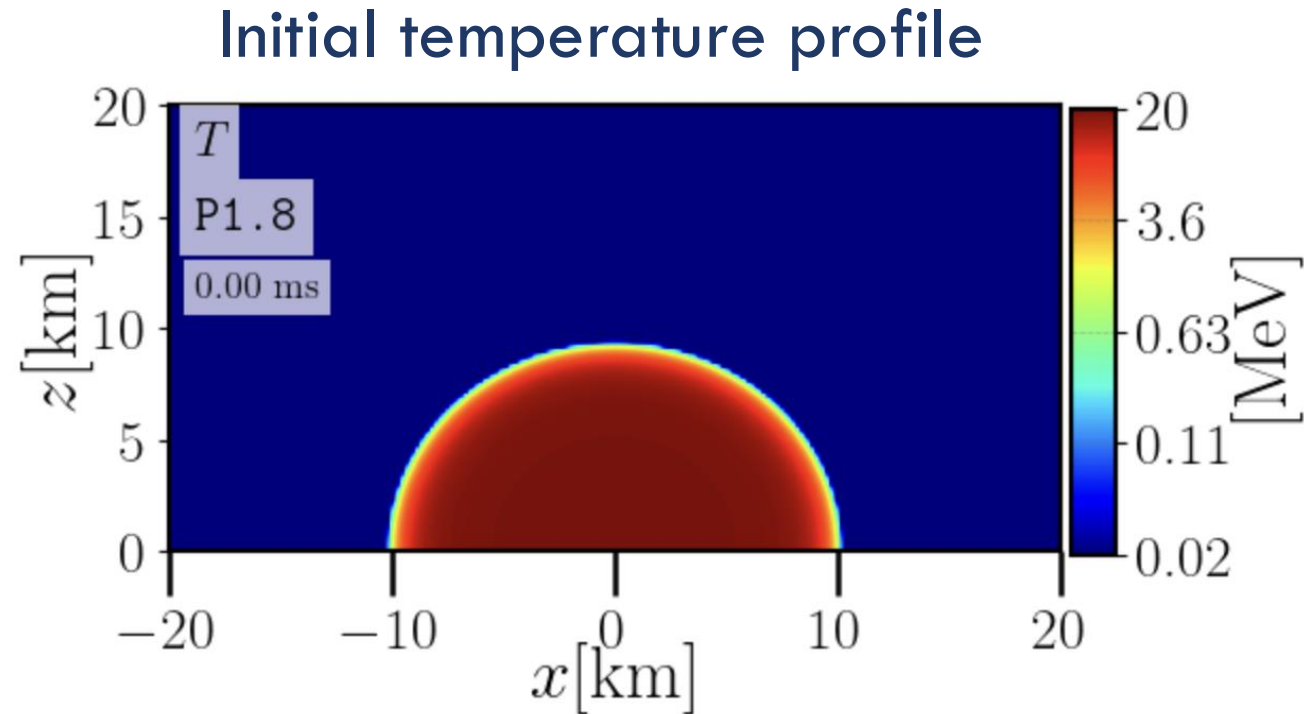
## Nested Grid setup

- Cartesian 3D grid with static mesh refinement. Reflection symmetry.

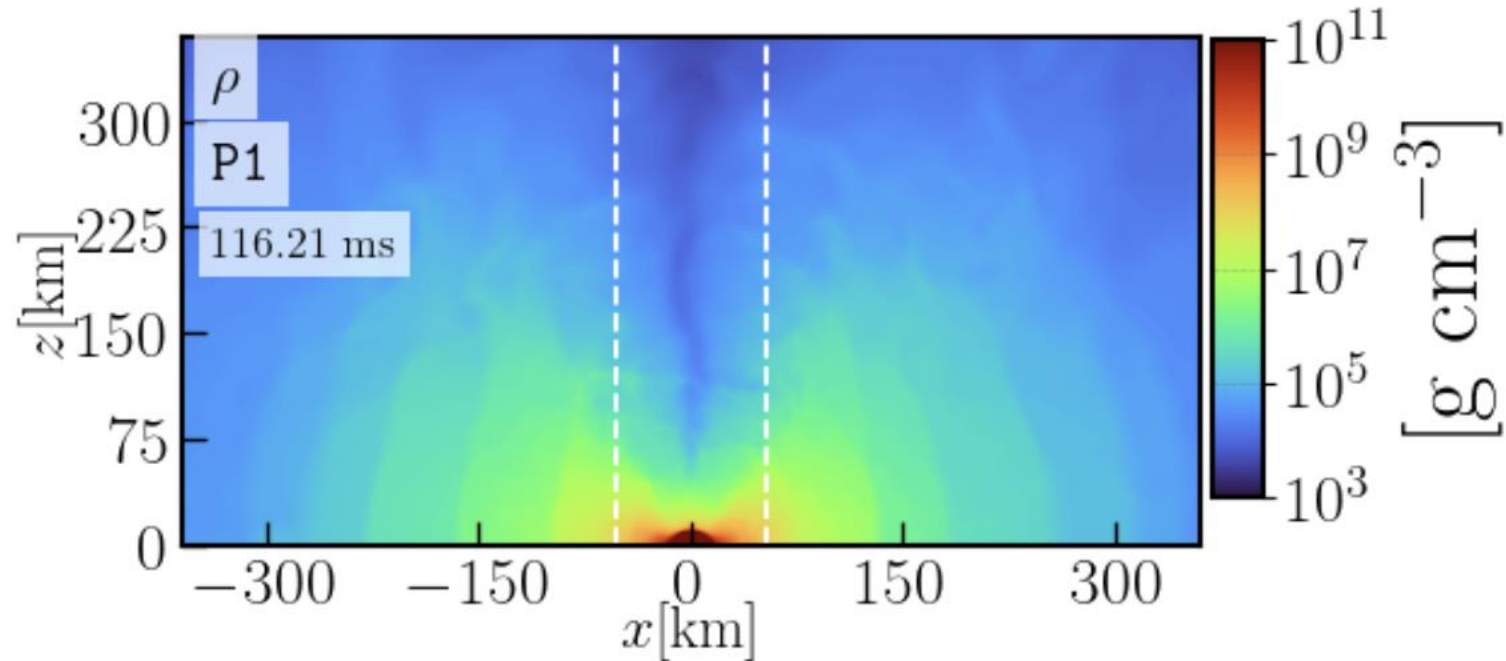


# Our GRMHD toy model: Initial Conditions

- 1.4 solar mass NS
- Hot:  $T_c \sim 20$  MeV, chosen to reproduce steady state neutrino emission
- Metric fixed after initial time step (low mass loss rate)



# First: models evolved to steady-state outflows

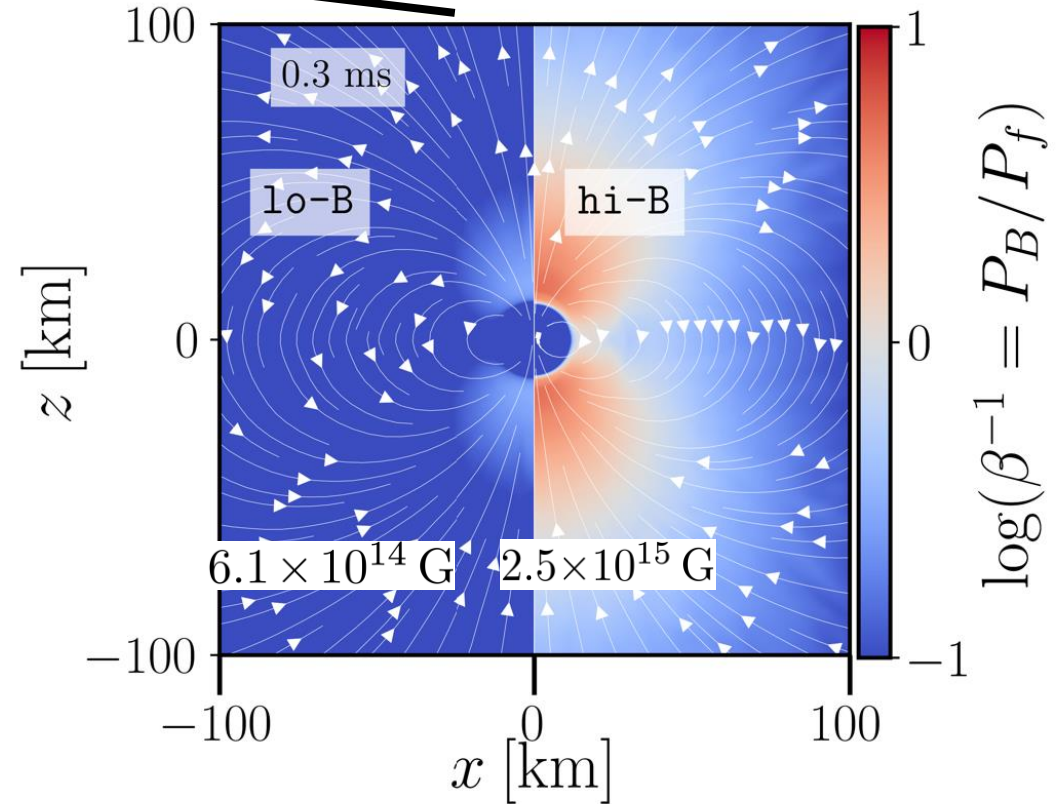


$t \sim 50-100$  ms

# Then, large-scale dipole magnetic field initialized

Again we evolve to steady-state outflows  
( $\sim 50$  ms)

We present these outflows.



# Criterion for rotation to impact the wind

Rotation speed at  
neutrinosphere radius is  
comparable to the  
sound speed

$$P_c \approx 2\pi \frac{R_\nu}{c_s}$$

Speed of sound  
(isothermal):

$$c_s \approx (kT/m_p)^{1/2}$$

Blackbody neutrino  
luminosity:

$$L_\nu \approx 4\pi \cdot 7/8 R_\nu^2 \sigma T^4$$

In terms of neutrino properties:

$$P_c \approx 2\pi \frac{R_\nu}{c_s} \approx \boxed{3.4 \text{ ms}} \left( \frac{R_\nu}{12 \text{ km}} \right)^{5/4} \left( \frac{L_\nu}{10^{52} \text{ erg s}^{-1}} \right)^{-1/8}$$

Wind dynamics impacted for rotation periods **less than ~3 ms**

# Criterion for magnetic fields to impact the wind

Magnetic pressure is comparable to fluid pressure at base of wind

$$P_B = B^2 / 8\pi$$

In terms of neutrino properties:

$$B_{\text{crit}} \approx \boxed{1 \times 10^{15} \text{ G}} \left( \frac{L_\nu}{4 \times 10^{51} \text{ erg s}^{-1}} \right)^{1/2} \left( \frac{R_\nu}{10^6 \text{ cm}} \right)^{-1}$$

Fluid pressure (radiation-dominated):

$$P_f \simeq P_{\text{rad}} = (11/12)aT^4$$

Blackbody neutrino luminosity:

$$L_\nu \simeq 4\pi 7/8 R_\nu^2 \sigma T^4$$

Wind dynamics impacted for magnetic field strengths **greater than  $\sim 10^{15} \text{ G}$**

# Criteria to impact the wind dynamics

rotation periods **less than  $\sim 3$  ms**

magnetic field strengths **greater than  $\sim 10^{15}$  G**

# Combining Magnetic fields and rapid rotation

Models  
tested

Surface  
poloidal  
Magnetic field  
strength

$B_0 \sim 3e15 \text{ G}$

$B_0 \sim 7e14 \text{ G}$

$B_0 \sim 3e14 \text{ G}$

$B_0 \sim 3e13 \text{ G}$

Most  
energetic

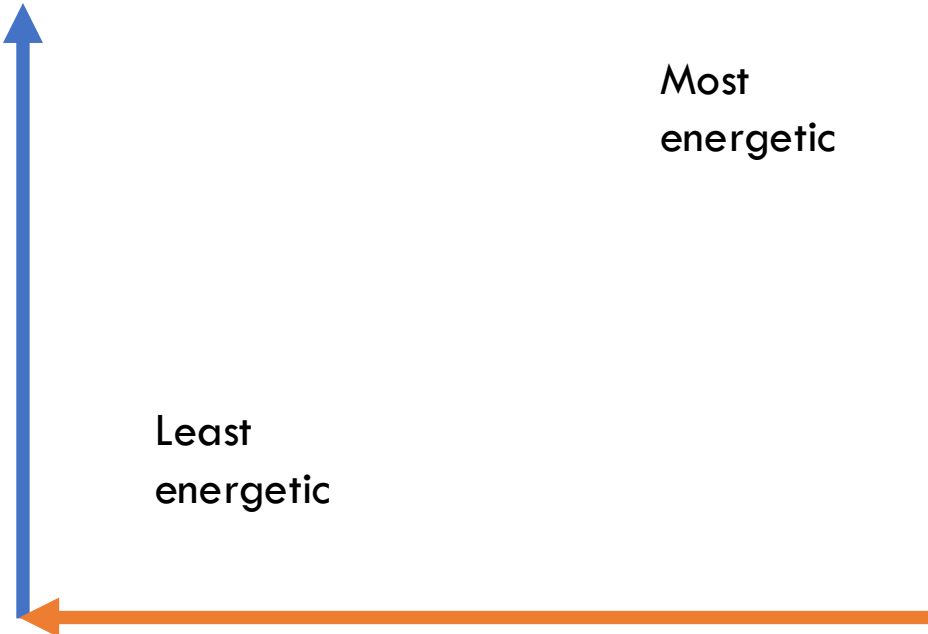
Least  
energetic

$P \sim 16 \text{ ms}$

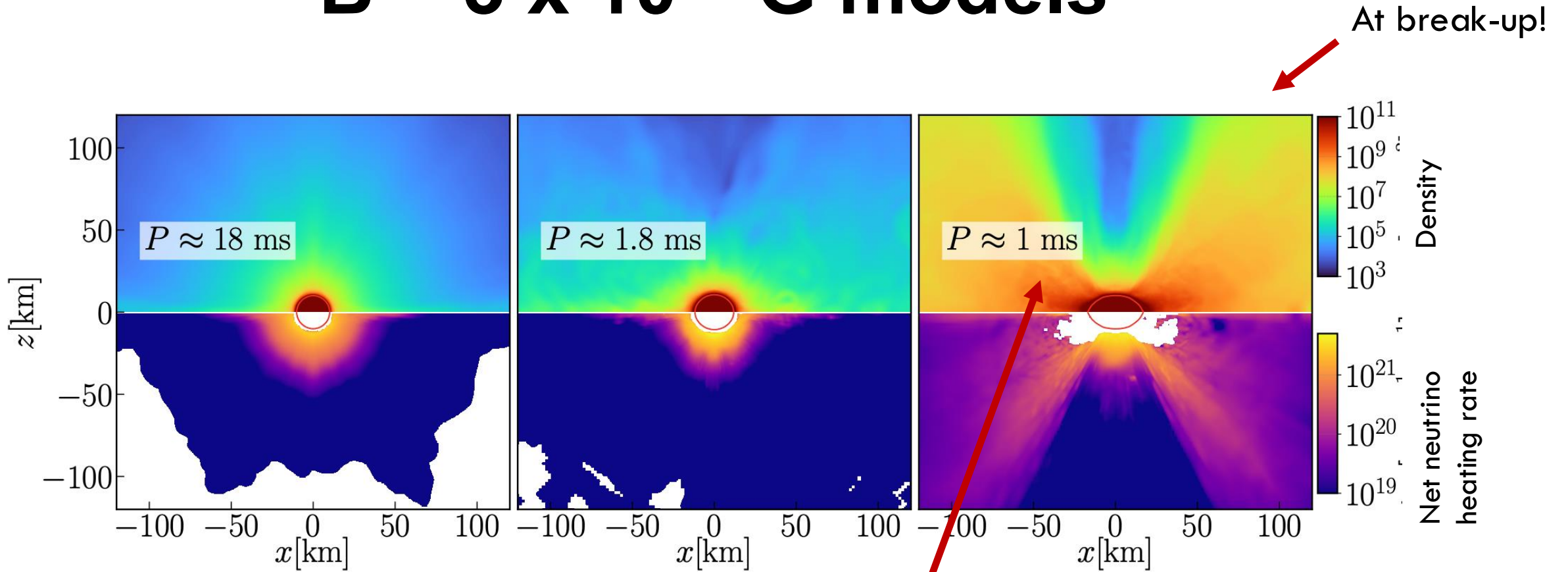
$P \sim 1.8 \text{ ms}$

$P \sim 1 \text{ ms}$

Period

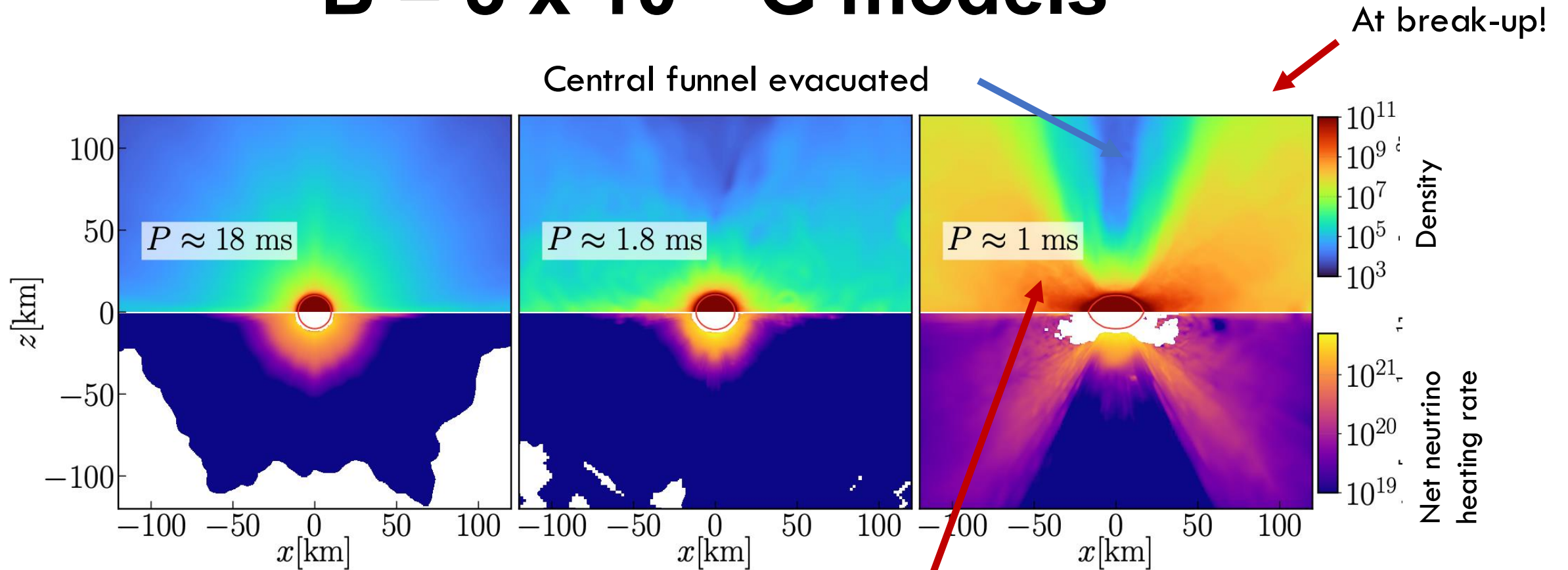


# $B = 3 \times 10^{15}$ G models



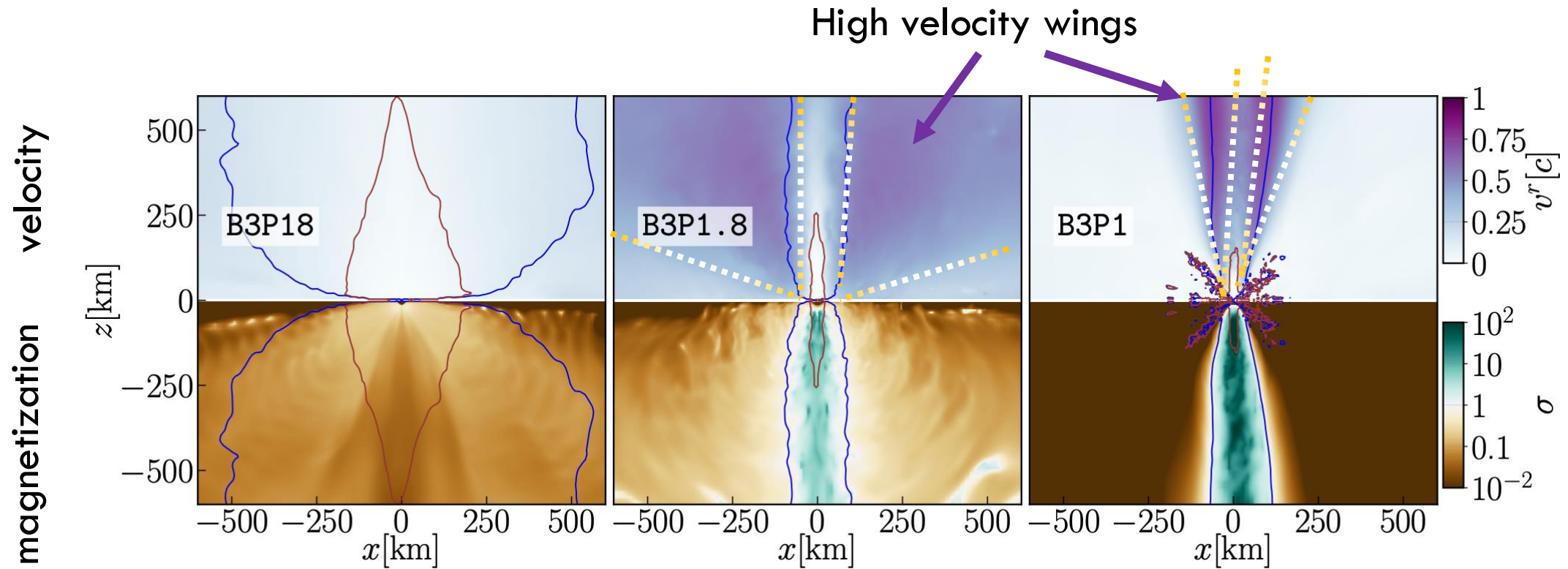
Dense equatorial outflows develop for  $P=1$  ms model

# $B = 3 \times 10^{15}$ G models



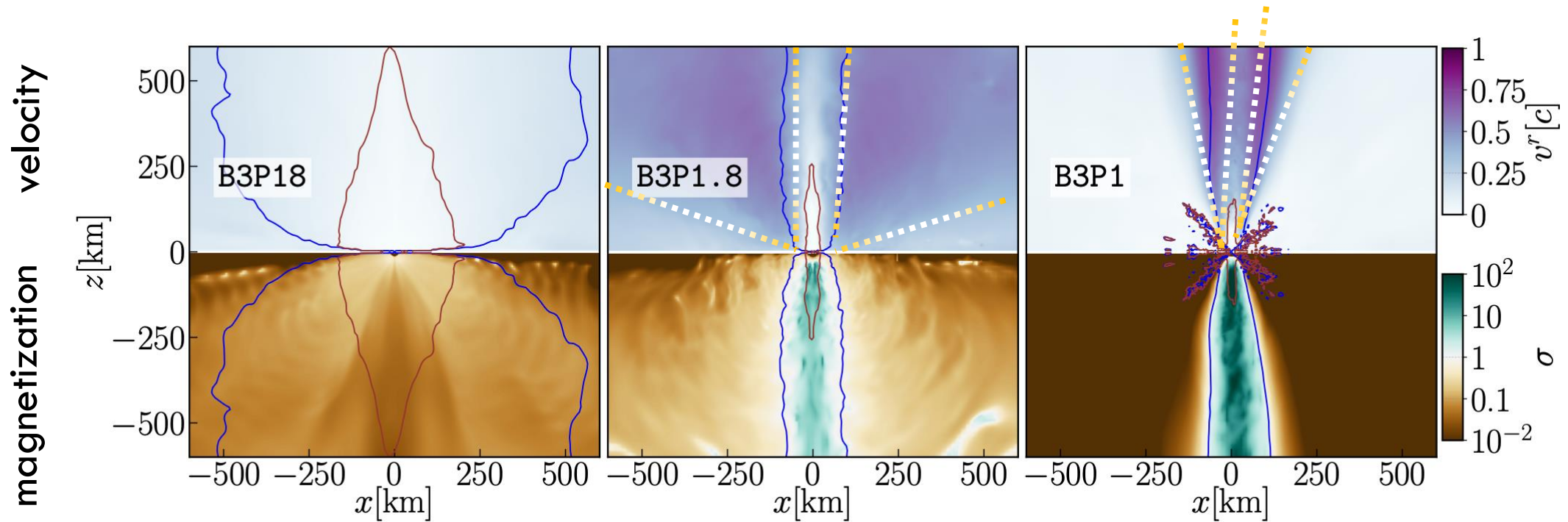
Dense equatorial outflows develop for  $P=1$  ms model

$$\mathbf{B} = 3 \times 10^{15} \text{ G}$$



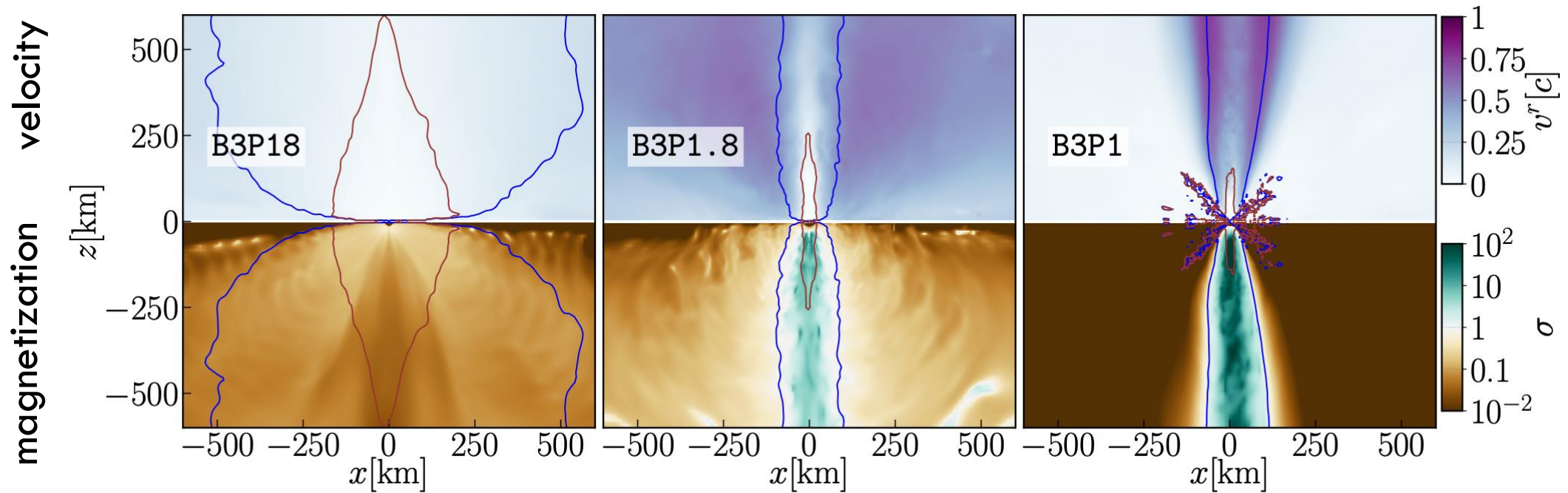
Smaller angular extent of high velocity region for  $P=1$  ms model

$$\mathbf{B} = 3 \times 10^{15} \text{ G}$$



$P=1$  ms polar outflows *collimated by equatorial outflows*

$$\mathbf{B} = 3 \times 10^{15} \text{ G}$$



High sigma 'jet' even at neutrino luminosities  $\sim 10^{52} \text{ erg/s}$ . Why??

# No centrifugal mass enhancement in polar direction

$$\sigma_{\text{outflow}} = \frac{\text{Poynting flux}}{\text{Mass flux}} = \frac{\Phi_B^2 \Omega^2}{\dot{M} c^3}$$

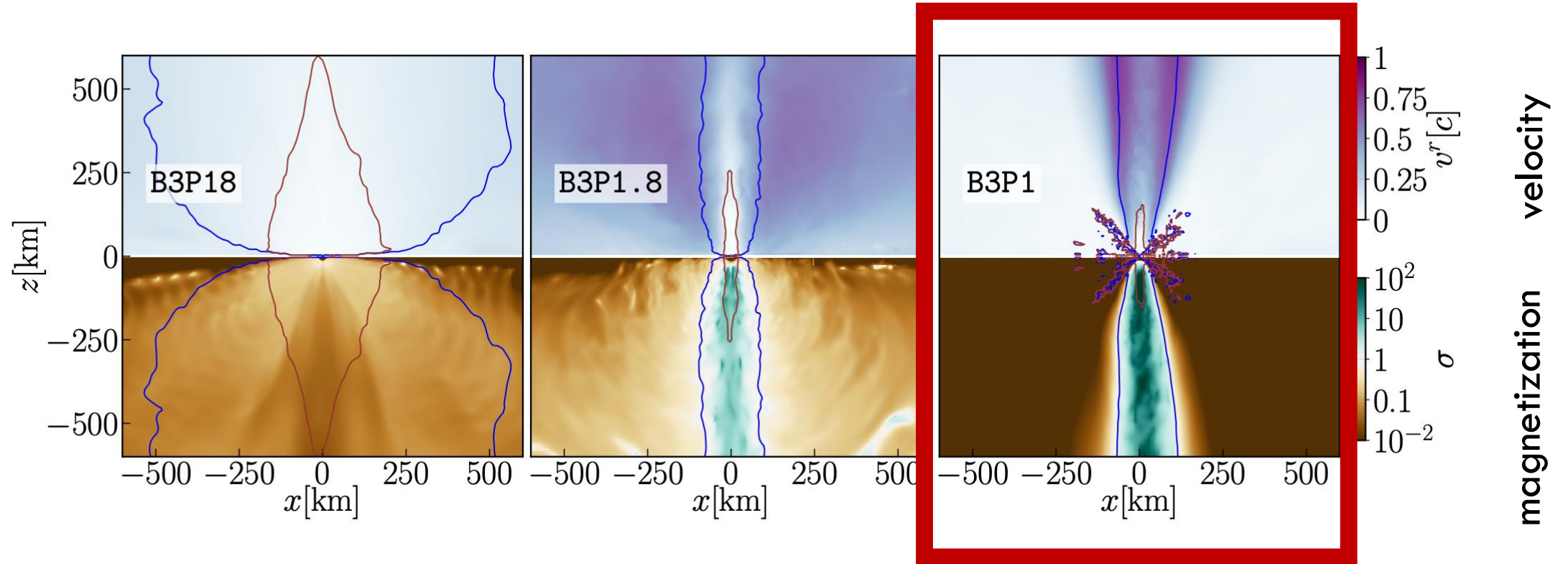
Magnetic flux

Rotation rate

mass flux  
Reduced in polar funnel

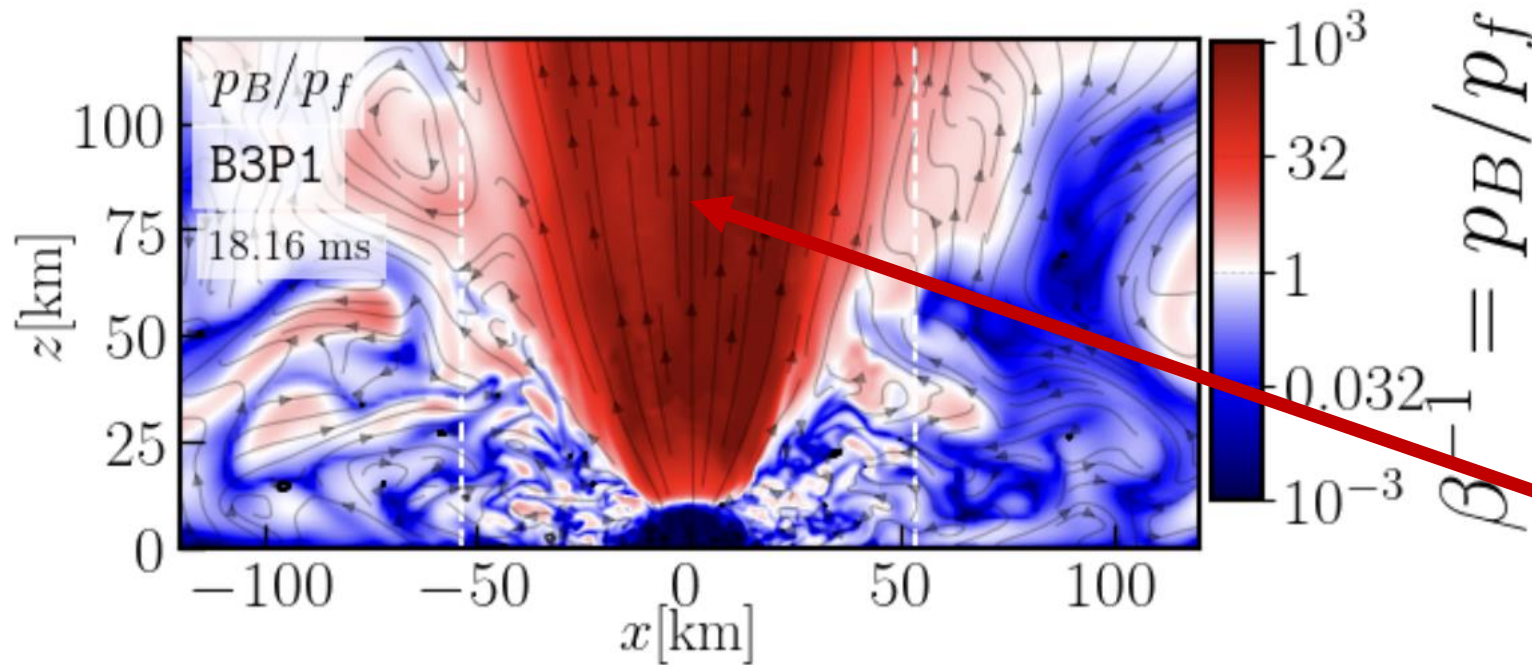
Not surprising!

$$B = 3 \times 10^{15} \text{ G}$$



High sigma 'jet' even at neutrino luminosities  $\sim 10^{52} \text{ erg/s}$

# Significant magnetic pressure support in central funnel



If present, massive disk/  
late time fall-back

must overcome

Centrifugal barrier +  
**magnetic pressure** in polar  
region

to impact outflows from such  
a protomagnetar

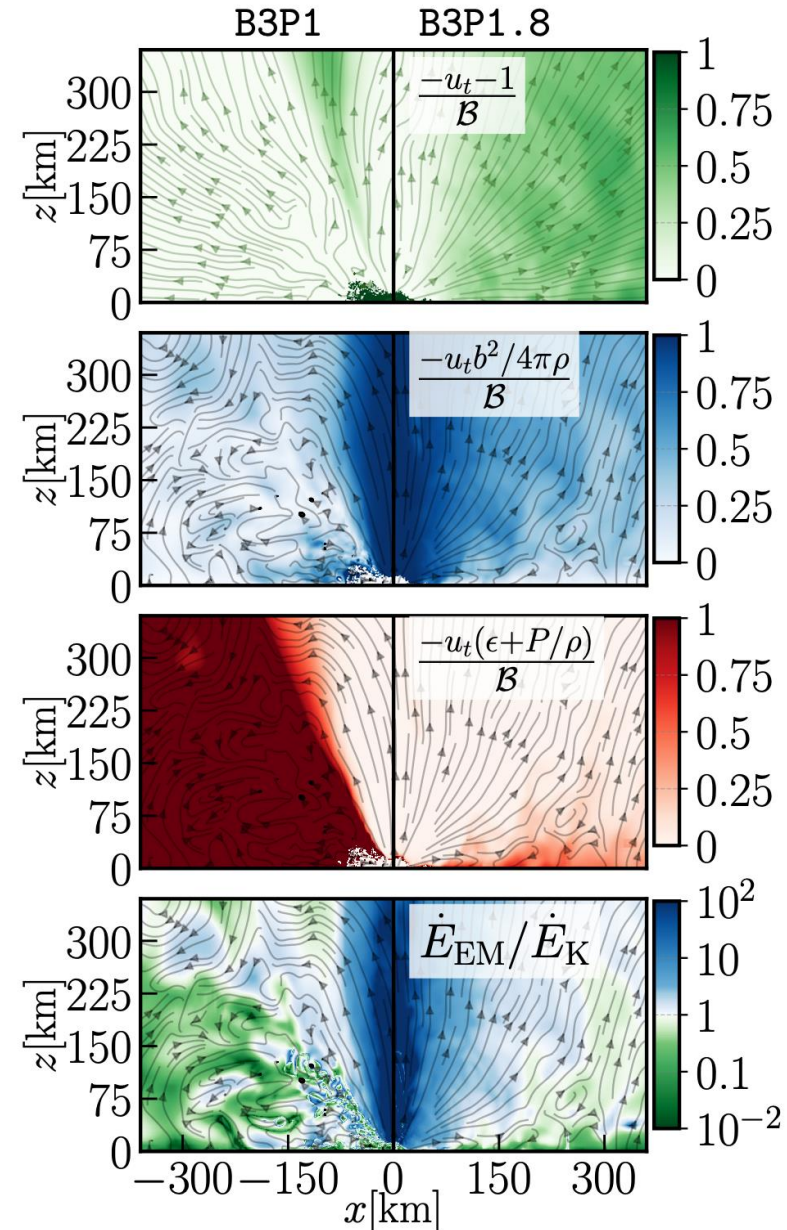
# Fractional energy contributions for $P = 1$ ms and $P = 1.8$ ms

$$B = 3 \times 10^{15} \text{ G}$$

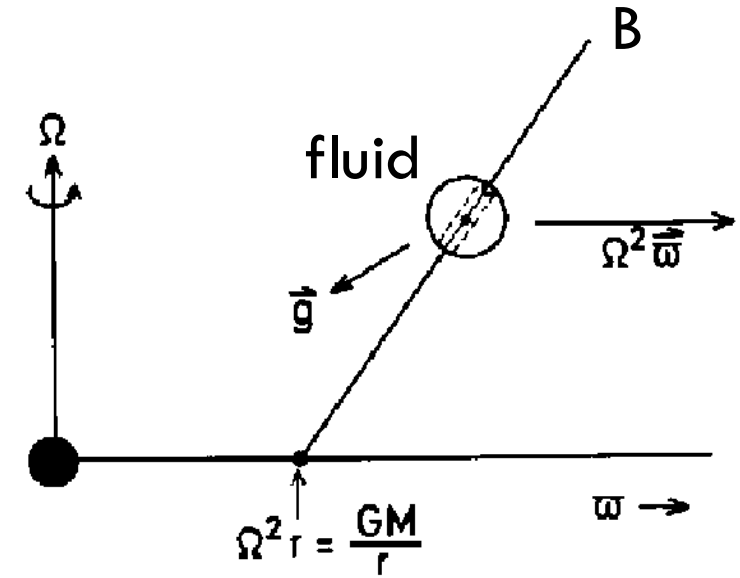
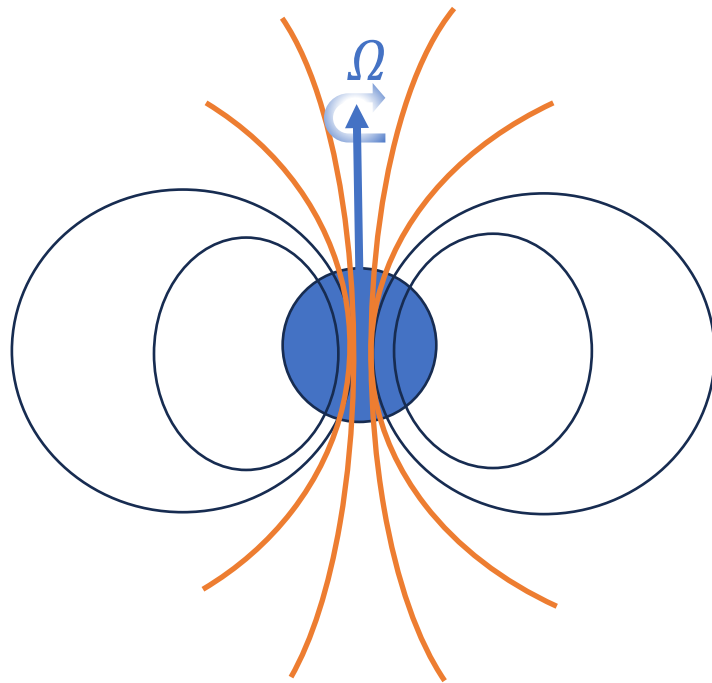
$$P = 1 \text{ ms}$$

$$B = 3 \times 10^{15} \text{ G}$$

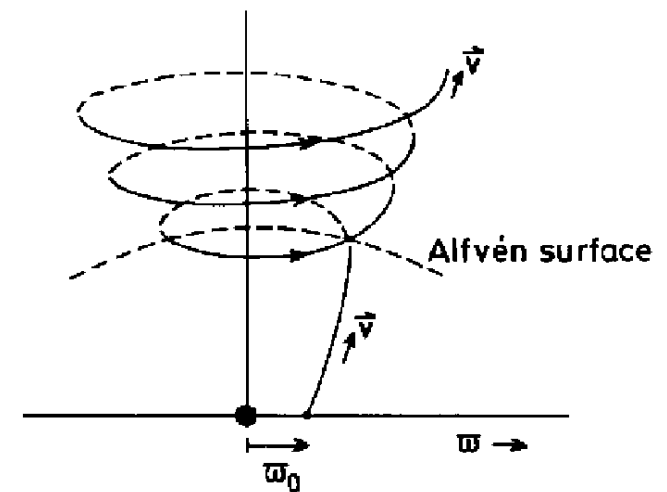
$$P = 1.8 \text{ ms}$$



# Rotation + magnetic fields can sling fluid along open field lines



Spruit 96



# Fractional energy contributions for $P = 1$ ms and $P = 1.8$ ms

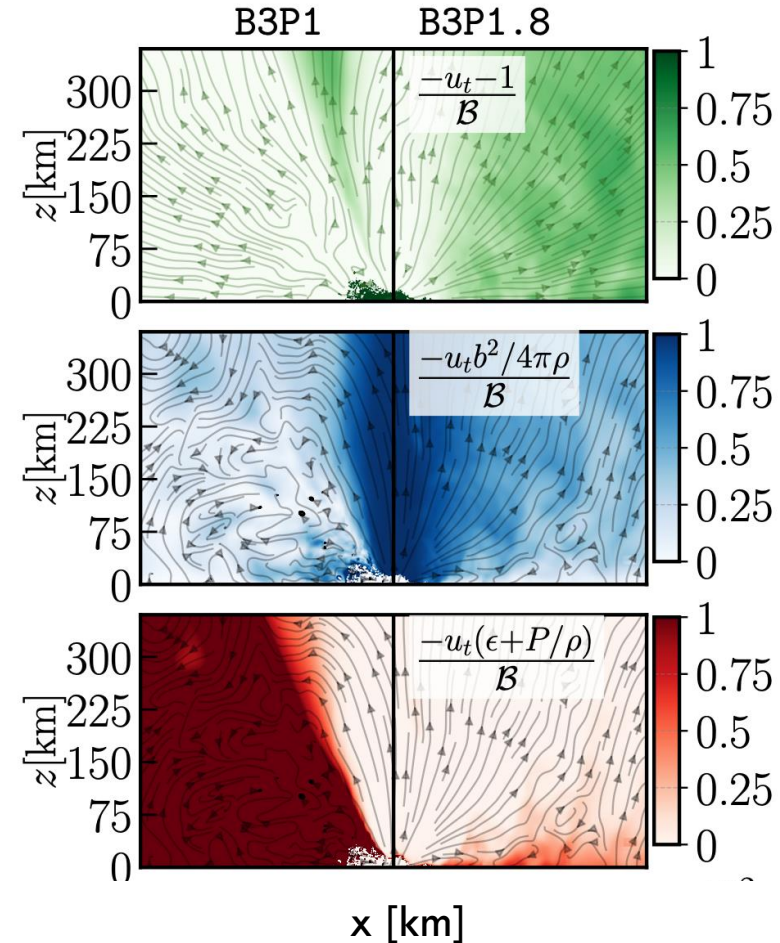
Magnetic energy dominates core of the jet

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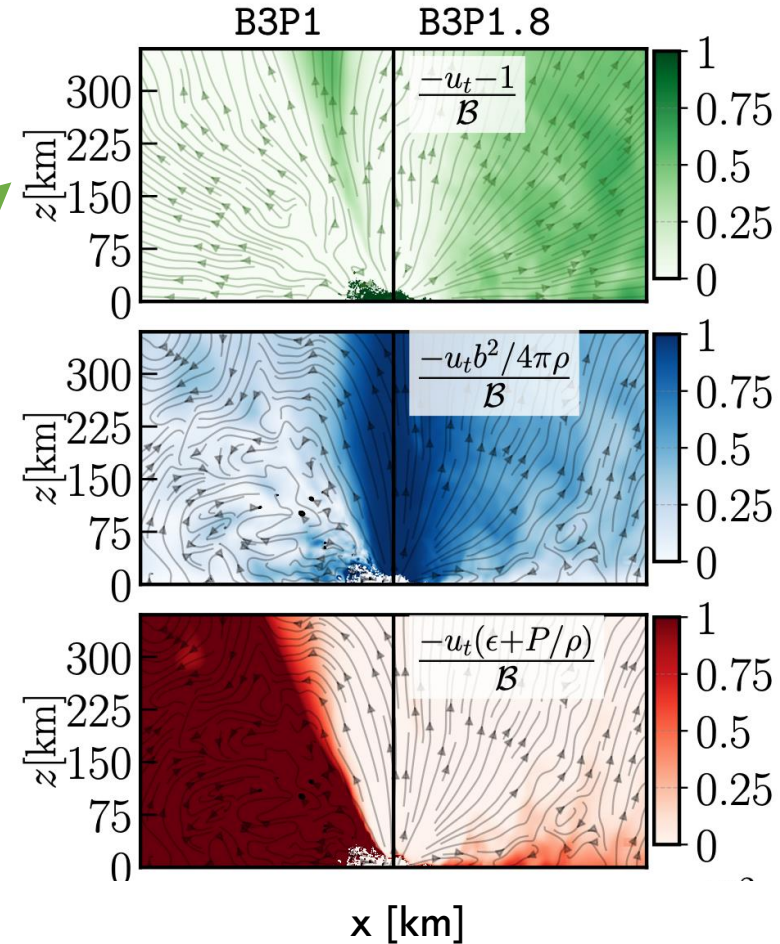
$B = 3 \times 10^{15}$  G  
 $P = 1$  ms

$B = 3 \times 10^{15}$  G  
 $P = 1.8$  ms

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1. Centrifugal slinging at play!

Energy converted to kinetic energy as fluid flows away from proto-magnetar



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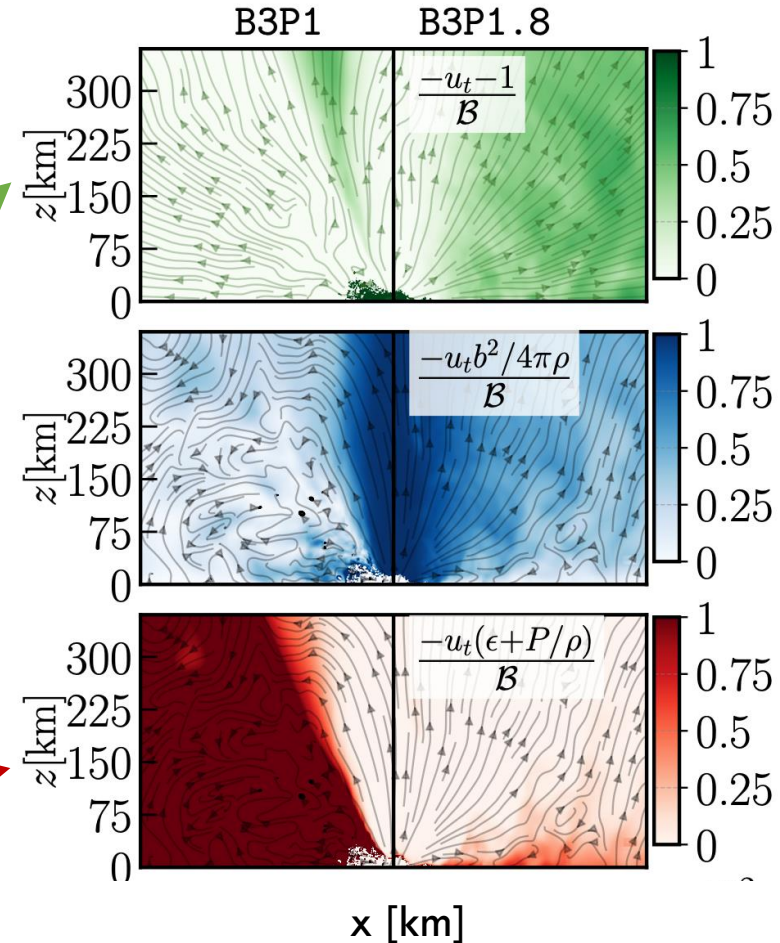


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thermal energy negligible in pole

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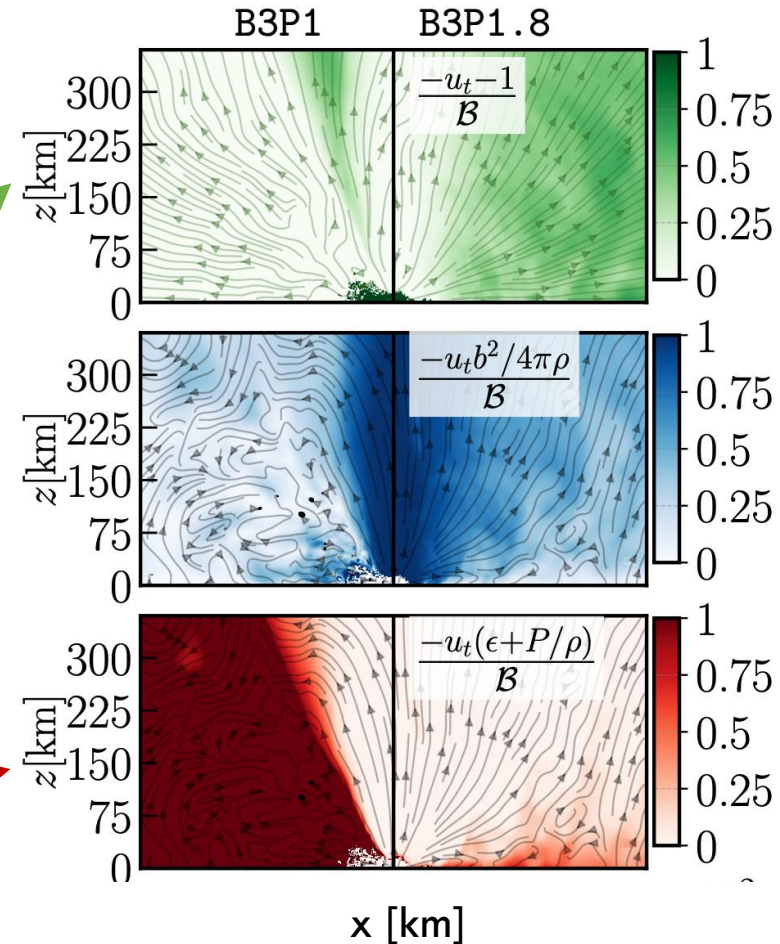
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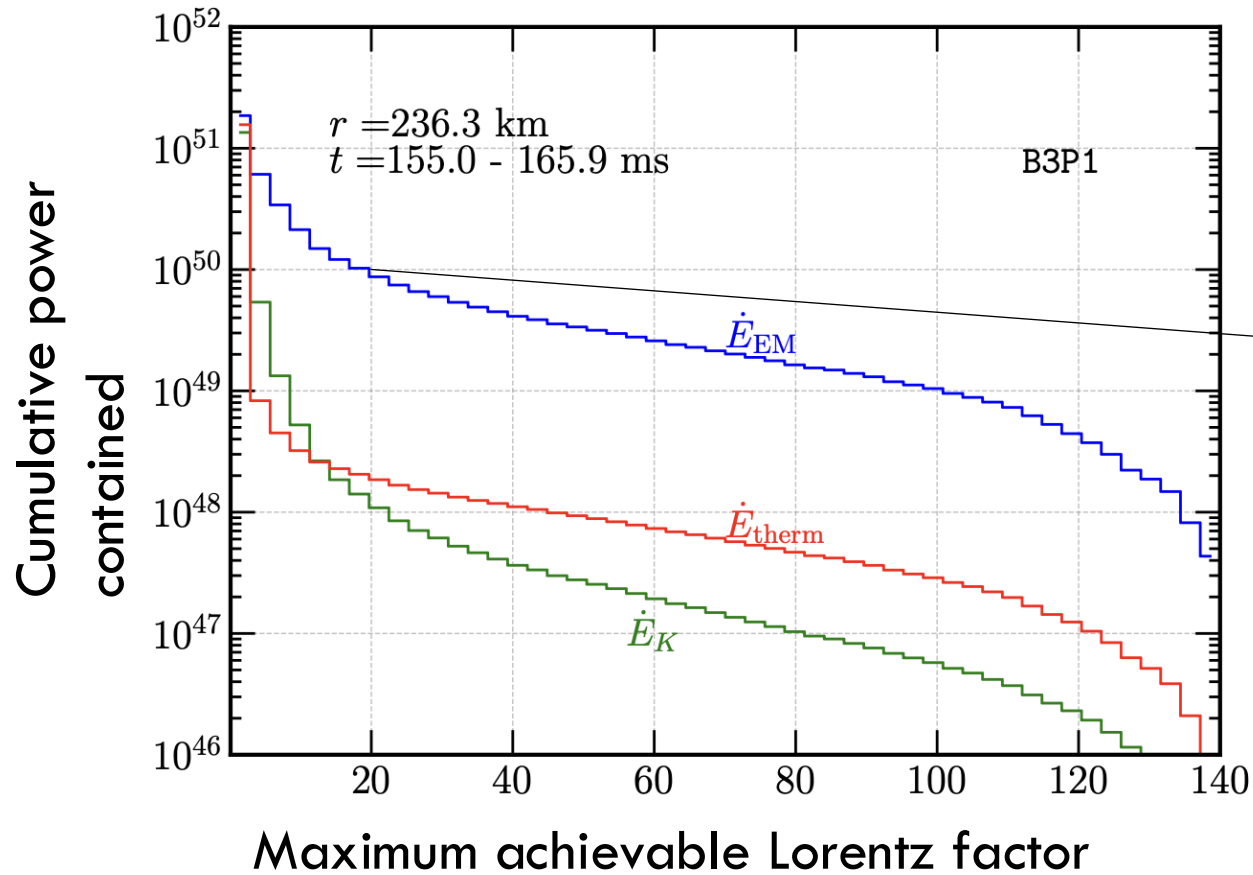
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2. Dense equatorial outflows focus acceleration into a polar jet



# Distribution of power to high Lorentz factors ( $\sim > 100$ )

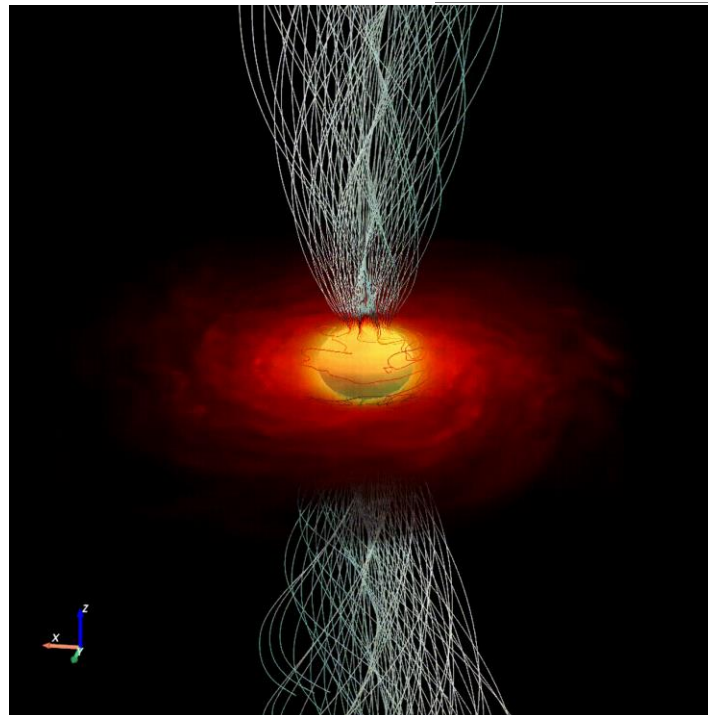
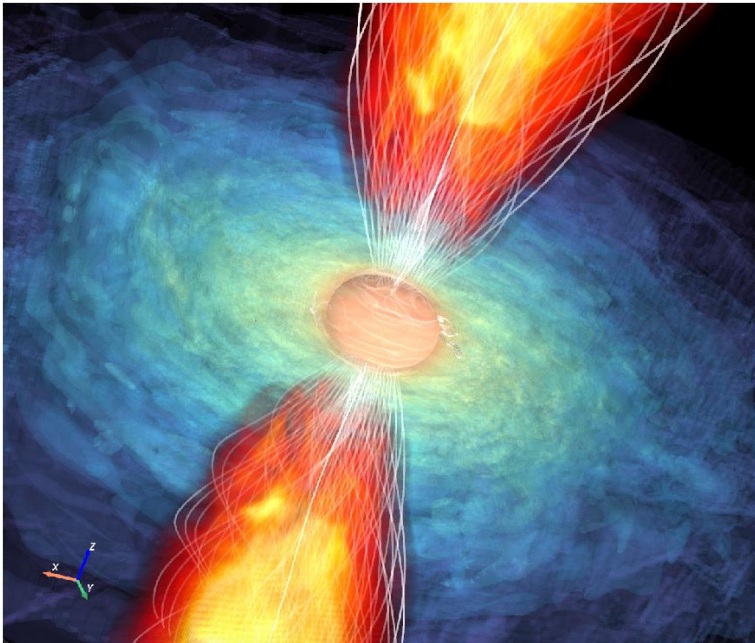
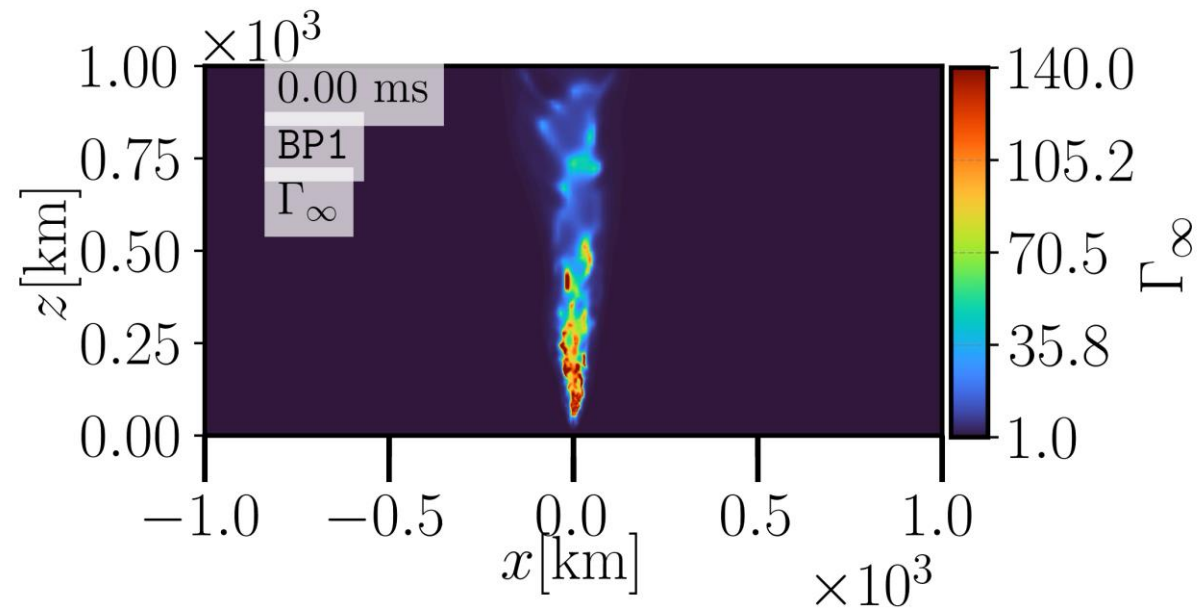


e.g.,  $10^{50} \text{ erg/s}$  power contained in fluid with Lorentz factor  $> 20$

These are powerful outflows!

$$\Gamma_{\infty} \approx \sigma$$

**Such outflows from a rapidly rotating proto-magnetar are physically possible**



$B \sim 3e15$  G

$P \sim 1$  ms

neutrino luminosities  
 $\sim 10^{52}$  erg/s

# Collimated, relativistic jet emergence is **not exceptional**

Rather a *natural consequence*  
of rotation + magnetization

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**Bonus:** **collimation** arises intrinsically at small radii,  
without requiring confinement by external envelope or  
supernova

Aids in  
 $E_{\text{mag}} \rightarrow E_{\text{kin}}$   
(over large distances)

# When & where might such remnants be formed?

- Energetic supernova  
(rapidly rotating progenitors)
- White dwarf collapse
- Neutron star mergers

# Magnetization evolves on neutrino cooling timescales

$$\sigma_{\text{outflow}} = \frac{\text{Poynting flux}}{\text{Mass flux}}$$

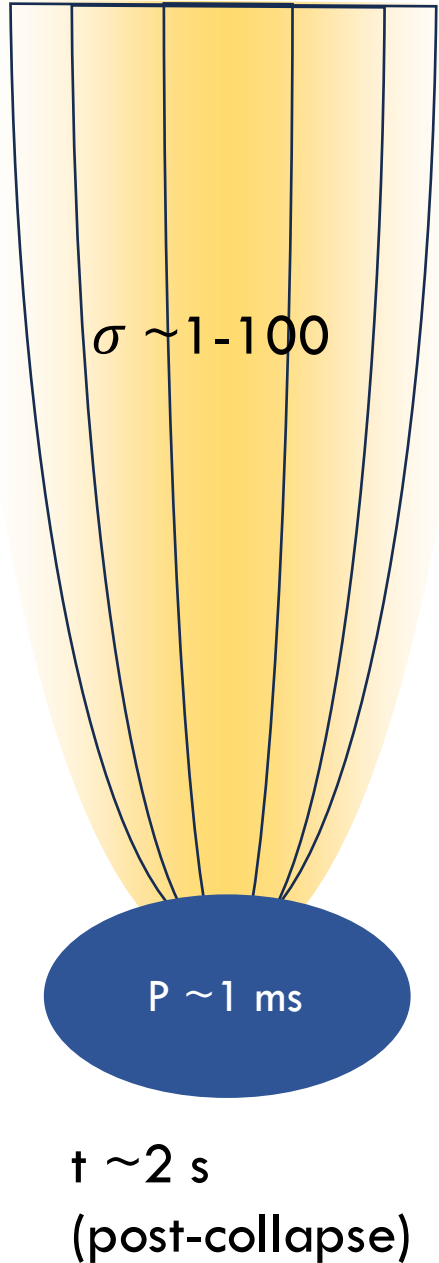
Magnetic field  
strength

Rotation  
rate slows

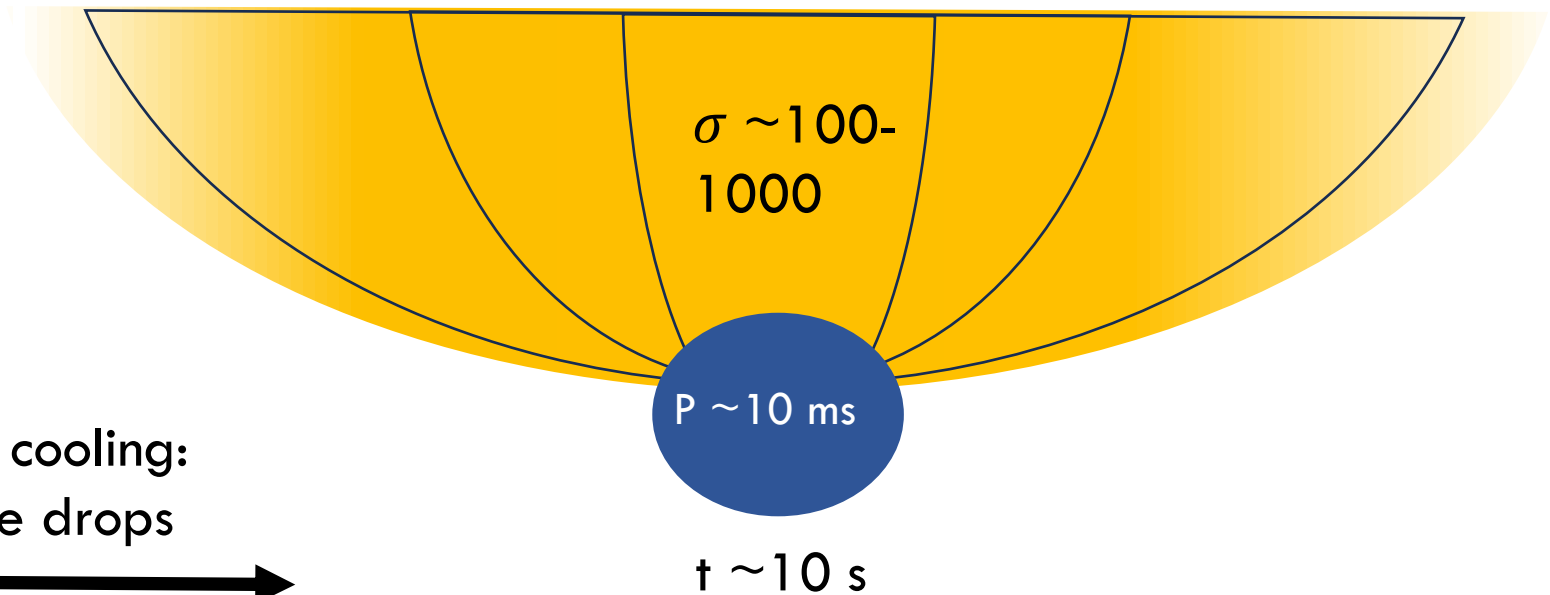
$$= \frac{\Phi_B^2 \Omega^2}{\dot{M} c^3}$$

**Mass loss rate drops  
drastically as proto-magnetar  
cools...**

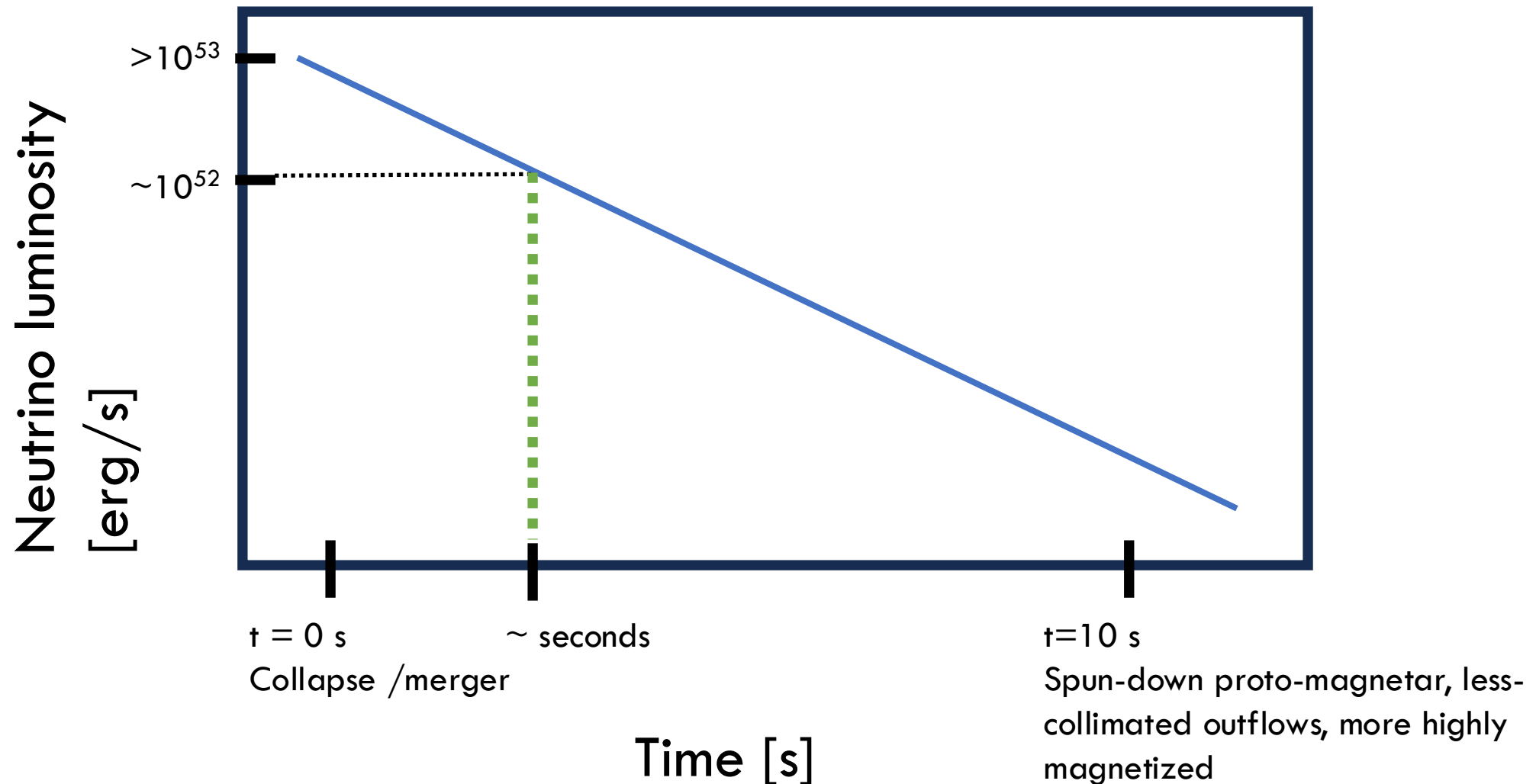
**...so the jet becomes less collimated, more magnetized over ~10 seconds**



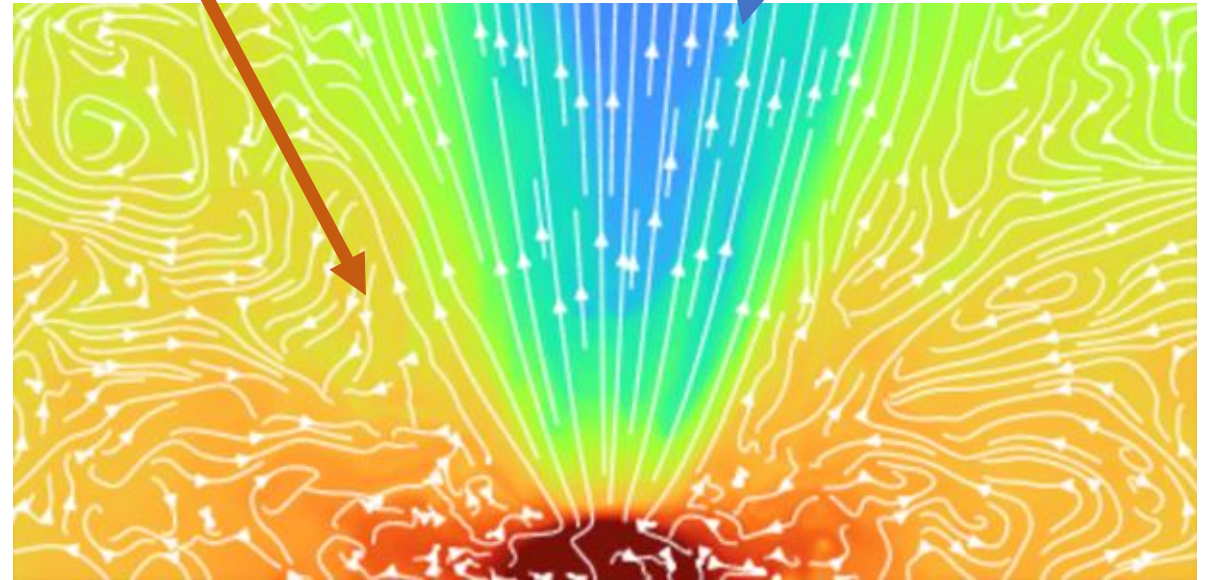
$t \sim 5-10 \text{ s}$  of cooling:  
mass loss rate drops



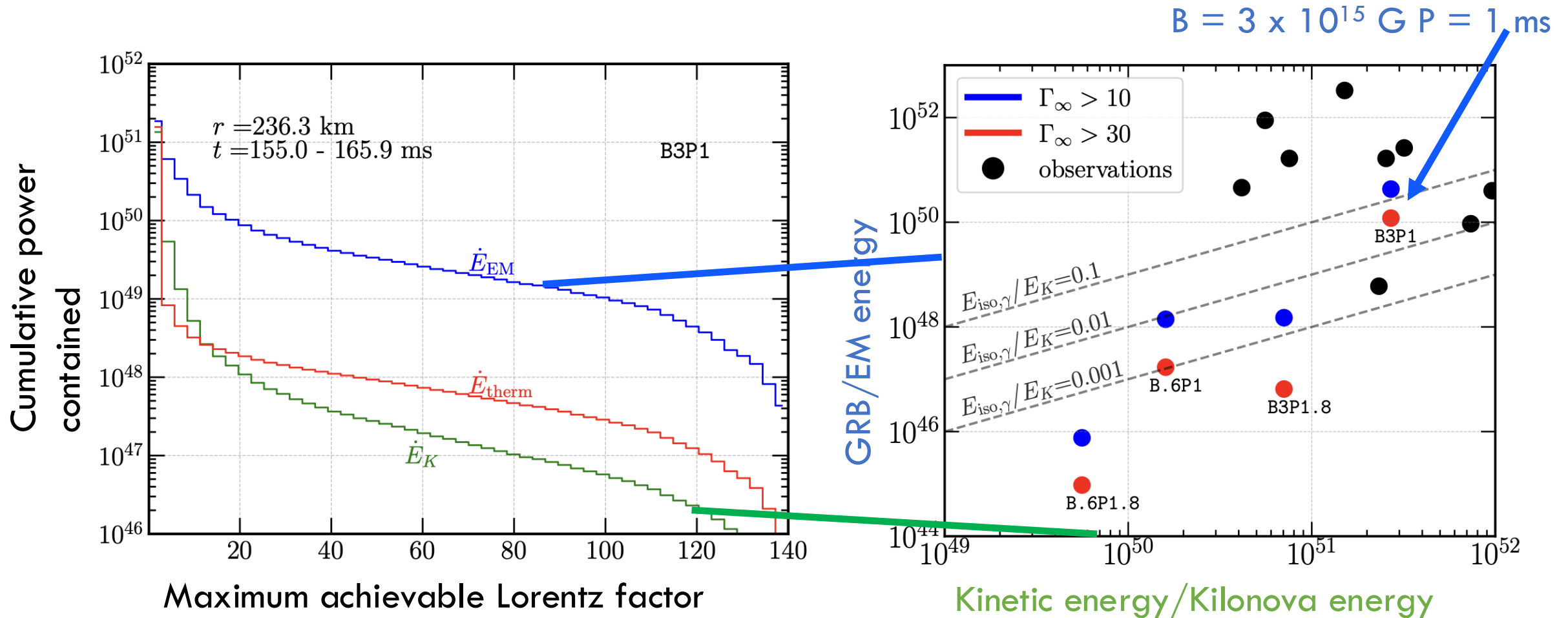
# A **jet** would occur within seconds of an event (prompt)



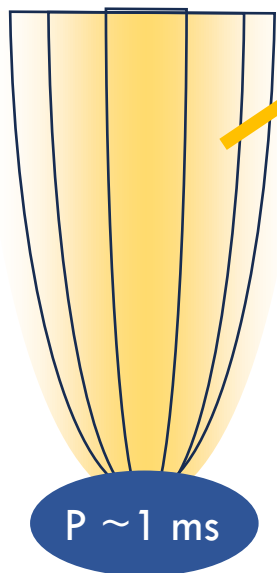
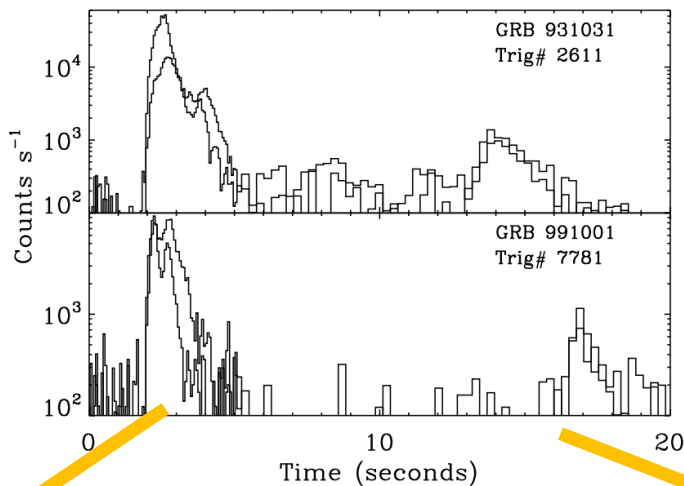
Outflows potentially account for short GRB (**jet**)  
and simultaneously long-lived counterparts  
(**equatorial outflows**)



# Outflows of millisecond proto-magnetar consistent with KN+ GRB observations

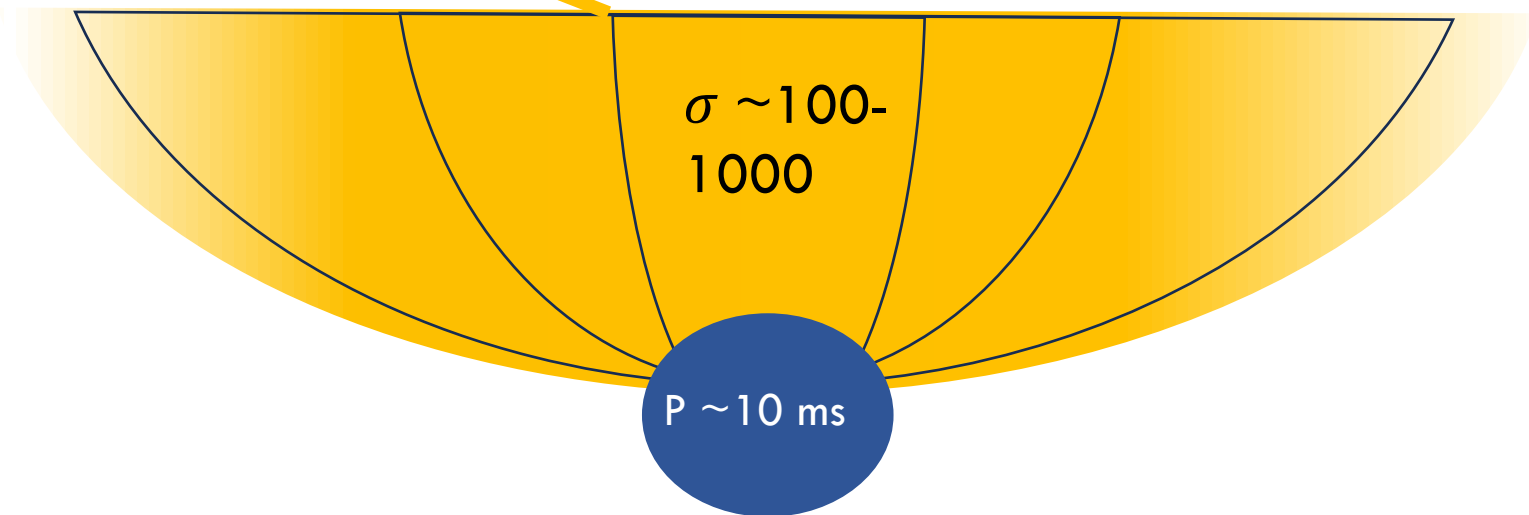


# Head-on observation: Short GRB + extended emission



$\sigma \sim 1-100$

$t \sim 2$  s (post-collapse)

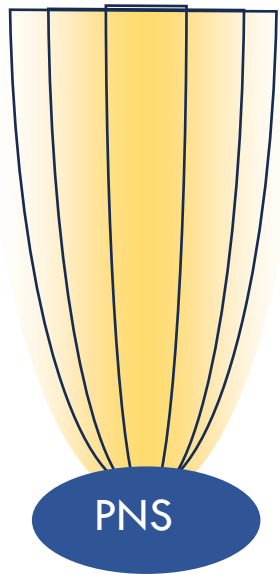
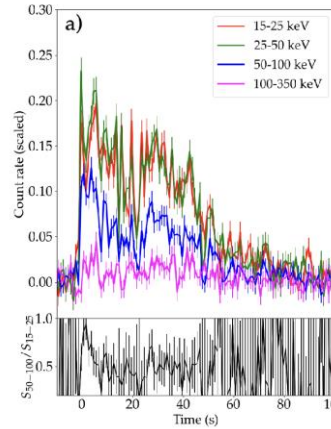


$\sigma \sim 100-1000$

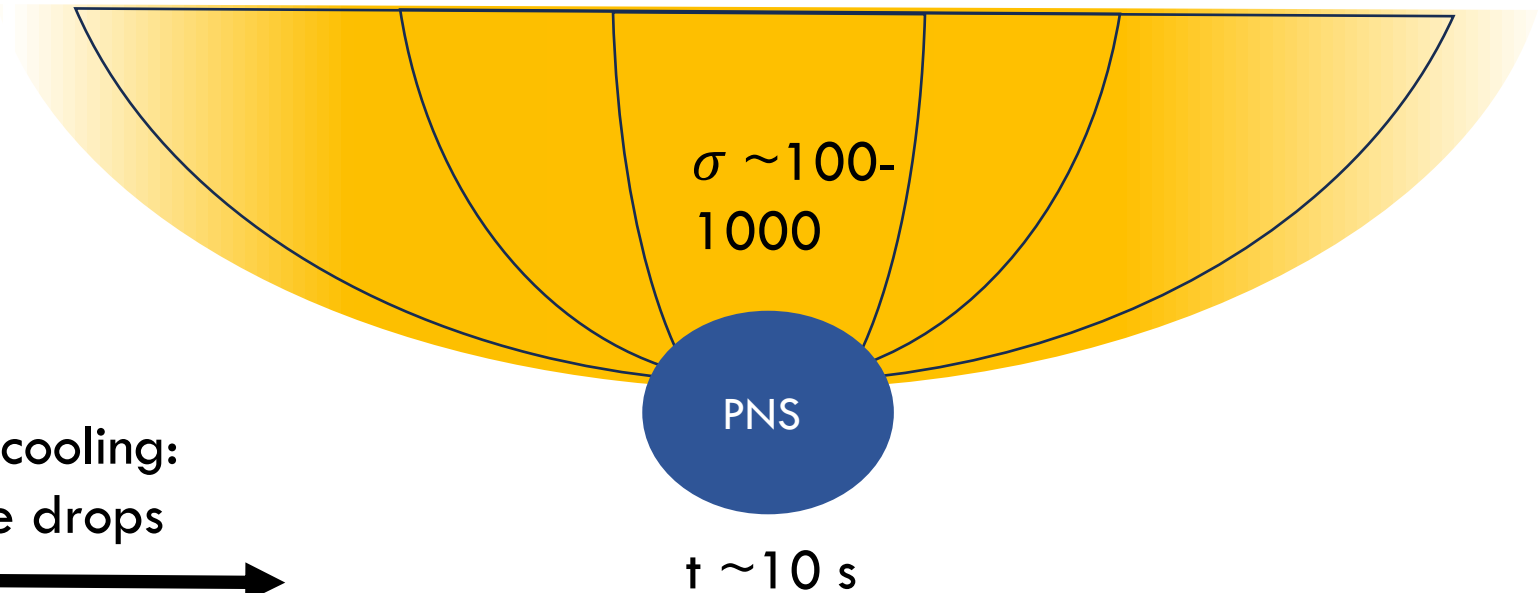
$P \sim 10$  ms

$t \sim 10$  s

# Off-angle observation: X-ray flash/long GRB



$t \sim 5-10$  s of cooling:  
mass loss rate drops



# Conclusions

- **continuous transition** in most proto-magnetar wind properties
  - low  $B \rightarrow$  high  $B$ , & low  $\Omega \rightarrow$  high  $\Omega$
- Rapidly rotating model –  $B \sim 3 \times 10^{15}$  G,  $P \sim 1$  ms (unique)
  - **massive equatorial outflows** which help to **naturally collimate jet**;  
enables efficient conversion from magnetic to kinetic energy at large distances
  - produces steady-state, fast polar outflows ( $\sim 0.7c$ ), and material with **high sigma ( $\sim 50-150$ )**
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# Next steps

- Interplay between disk and slower rotating models
- Nucleosynthesis
- Global simulations and better neutrino transport
- Input to large-scale jet simulations



*That's all Folks!*