

# Forward baryon production and meson structure in $ep$ collisions

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- High-energy DIS and forward baryon production: reviewing HERA results
- Interpretation with meson exchange model
  - One-Pion Exchange (OPE) models and non-OPE contribution
  - Pion flux measurement and extraction of pion structure
  - Yield of forward protons and neutrons
- Factorisation breaking in baryon production
  - Baryon yield for photoproduction and DIS
  - $t$  dependence of the baryon in inclusive processes
  - cf: exclusive processes (diffraction, vector meson production, DVCS ...)
- 1<sup>st</sup> year measurement at the EIC: what to measure first

# Forward baryon production and DIS

- Baryon number should conserve: there should exist either a proton or neutron
- A popular view: t-channel **particle (meson) exchange**
  - Diffractive production (Pomeron) exists!
- More generic: **fragmentation**
  - **Limiting fragmentation**, in particular: the fragmentation occurs independent of the upper vertex i.e.  $(x, Q^2)$

Kinematical variables for leading baryons (LNs):

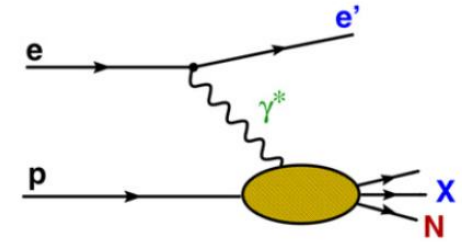
$z = p'_L/p = x_L$  : fraction of the longitudinal mom. carried by LN

$p_T(LN) = zp \sin \theta_{LN}$  : transverse momentum of the LN

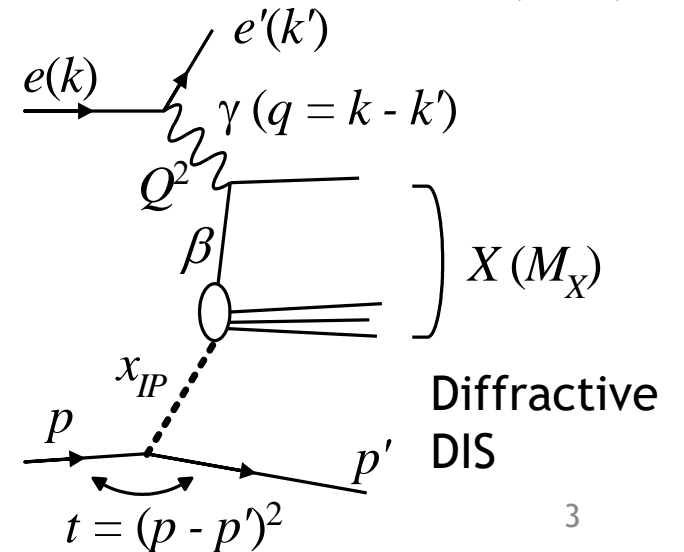
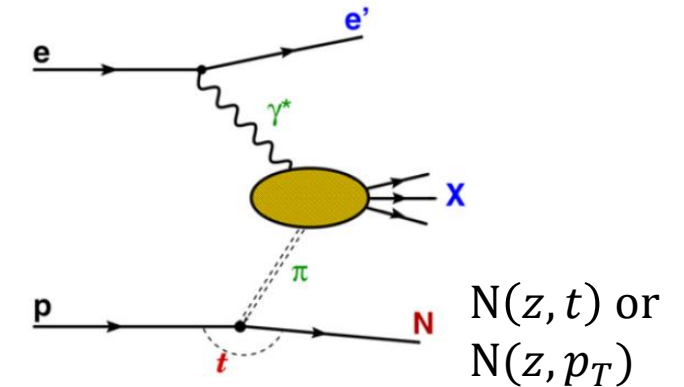
$t = (p - p')^2 \simeq -p_T^2/z = -zp \sin^2 \theta_{LN}$

: momentum transfer via the exchange

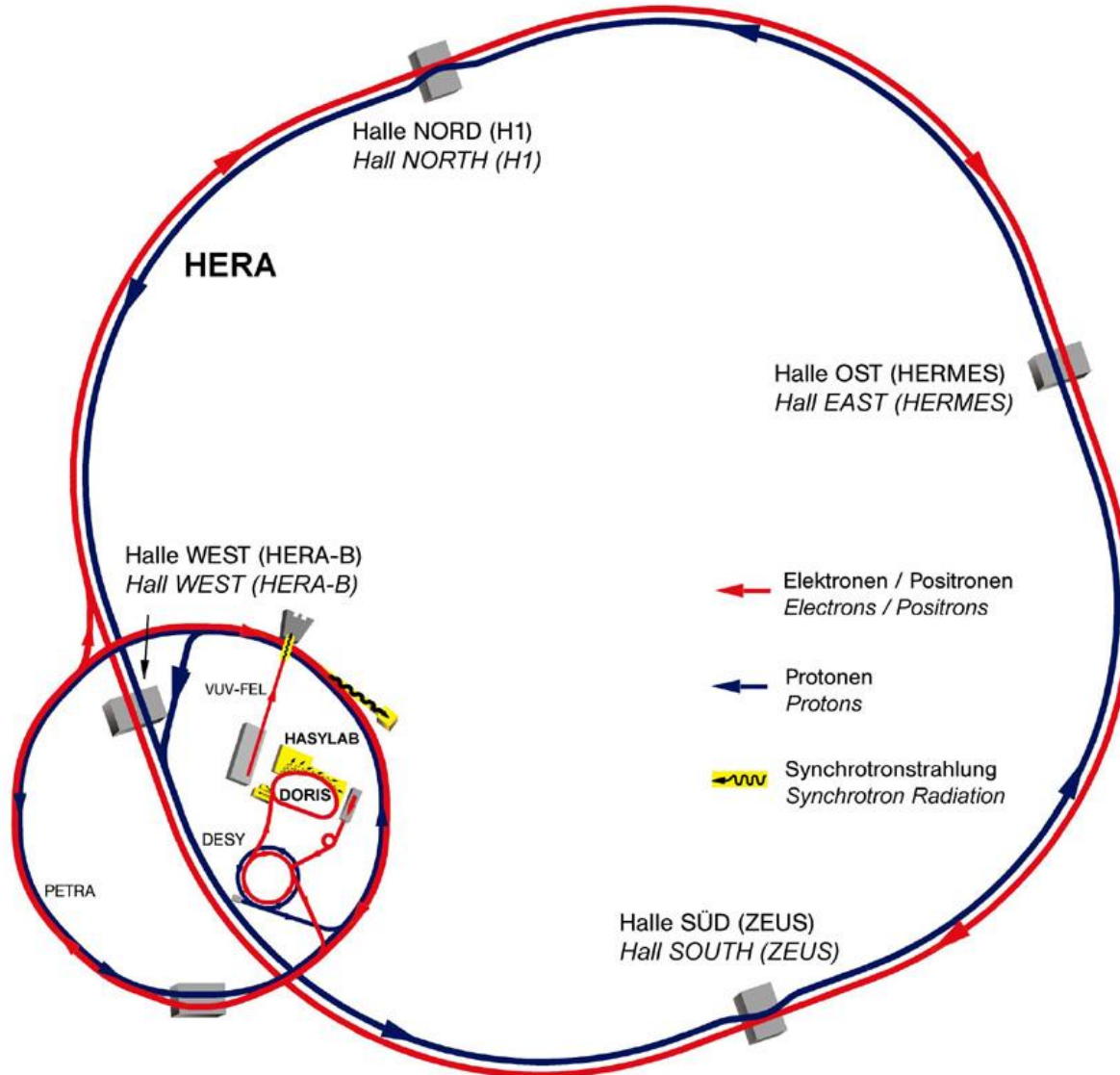
fragmentation



Meson exchange model



# HERA: the only electron-proton collider (so far)

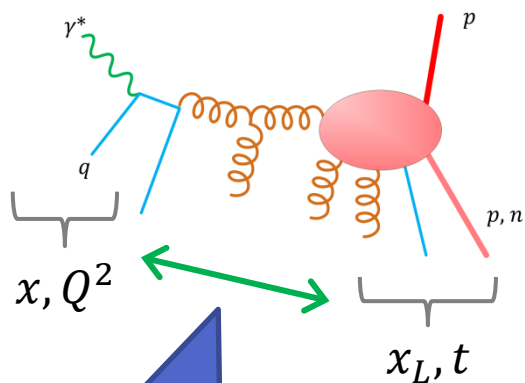


- Circumference: 6.3 km (similar size to the Tevatron at Fermilab)
  - Proton beam: 920 GeV
  - Electron/positron beam: 27.5 GeV
- centre-of-mass energy  $\sqrt{s} = 318 \text{ GeV}$
- Resolve structure down to  $10^{-18} \text{ m}$
- 220 bunch operation (96ns bunch spacing)
  - Operated in 1992-2007

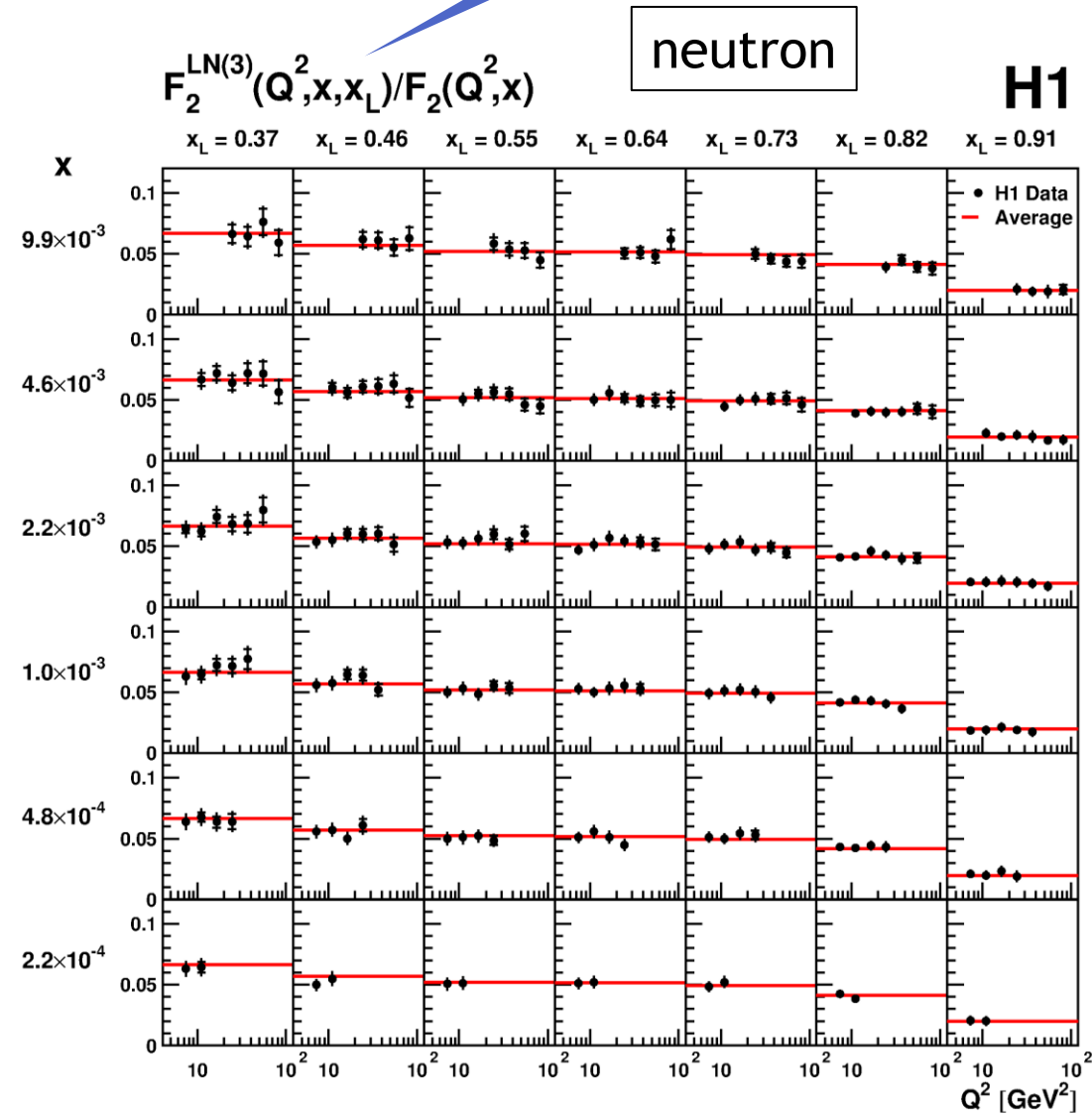
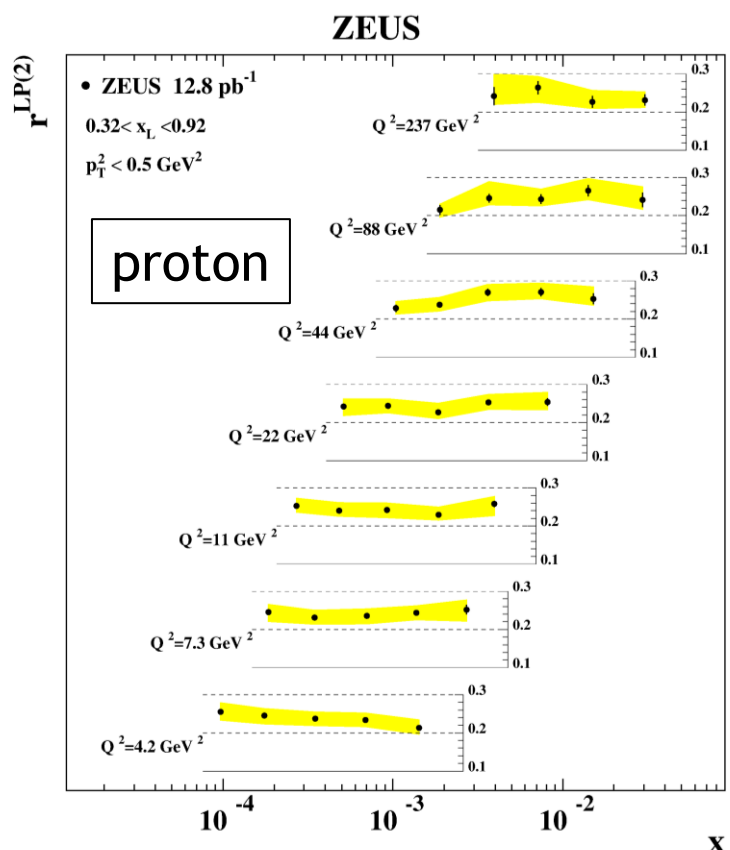
# Does the baryon talk to the virtual photon?

$$x_L \equiv z$$

- The answer is basically no!
  - No strong yield dependence on  $x$ ,  $Q^2$
  - “limiting fragmentation”



The photon and LN vertices are way separated in rapidity

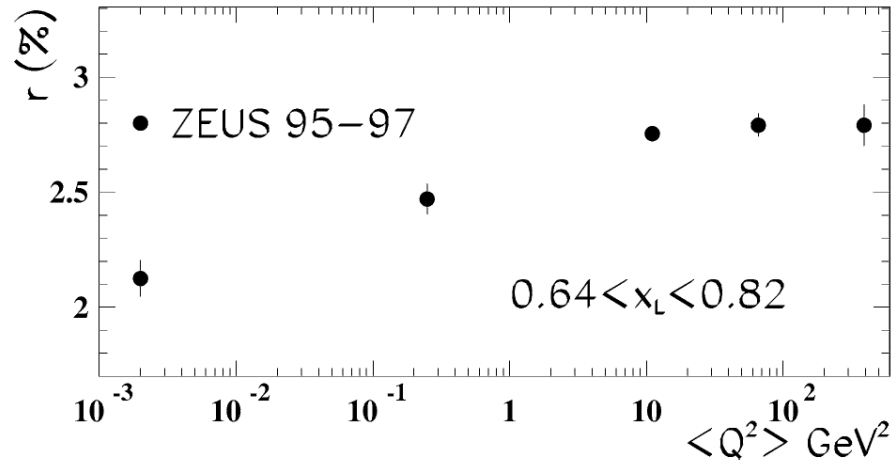


fraction of events with proton is flat in  $x$

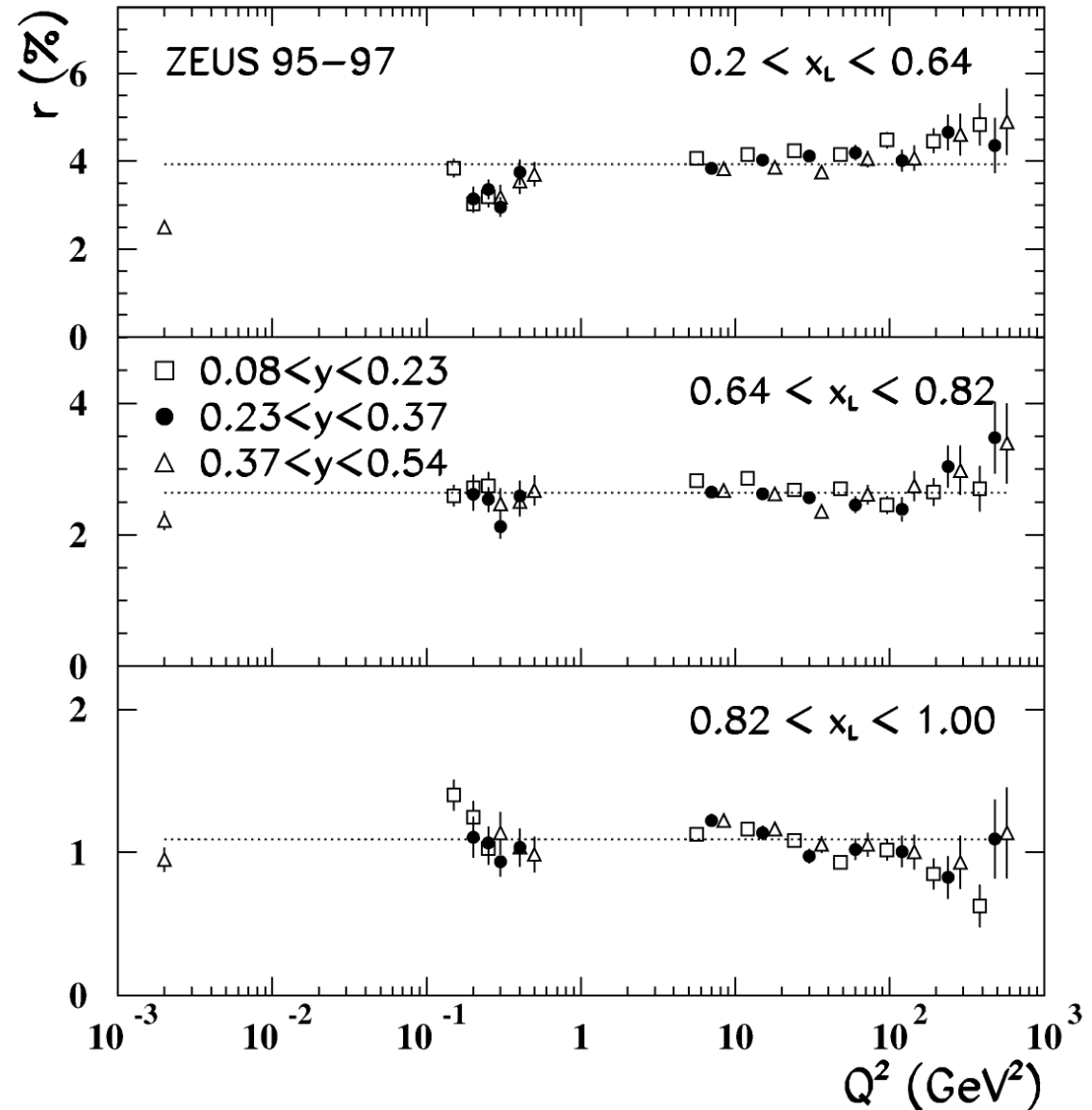
... and flat in  $Q^2$  for given  $(x, z)$  bins

# A bit more in detail: $Q^2$ dependence?

NPB 637(2002) 3-56 (also right figure)



- slight dependence in  $Q^2$ 
  - re-scattering (absorption) for photoproduction events? (photon = hadron)
- different dependence for  $x_L$ 
  - low- $x_L$  : stronger dependence
  - what is the mechanism?

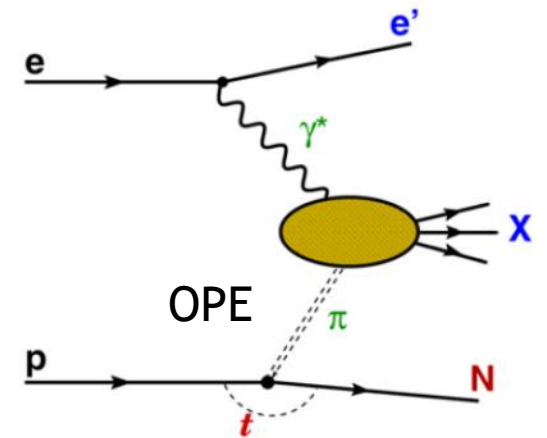
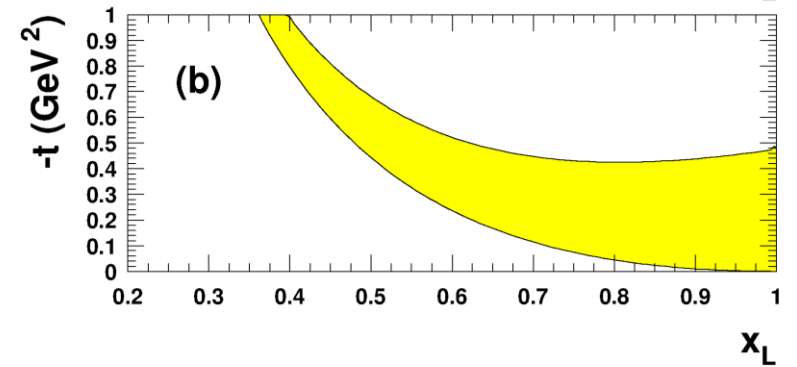
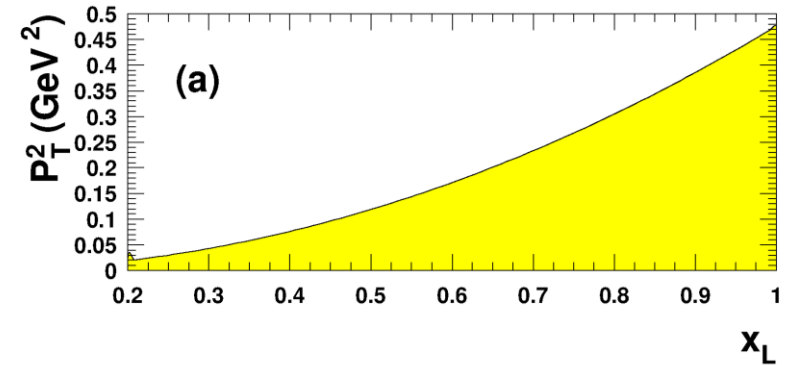
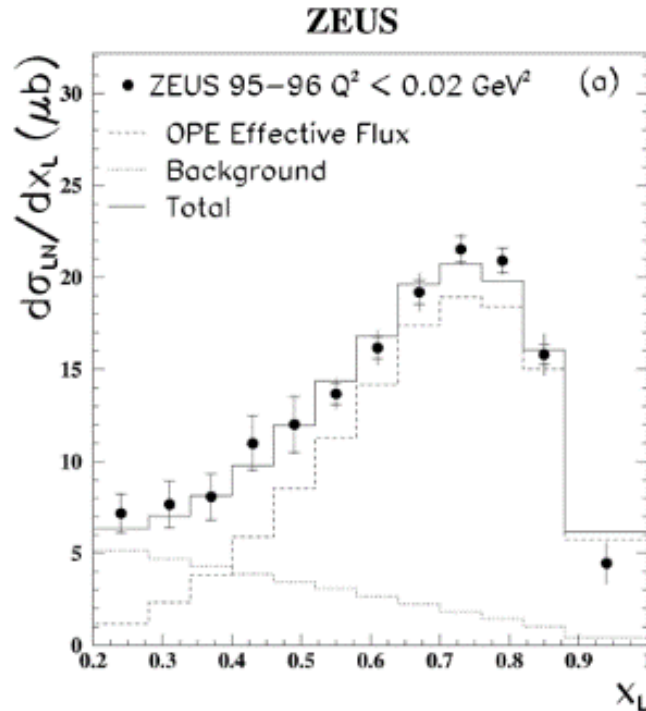
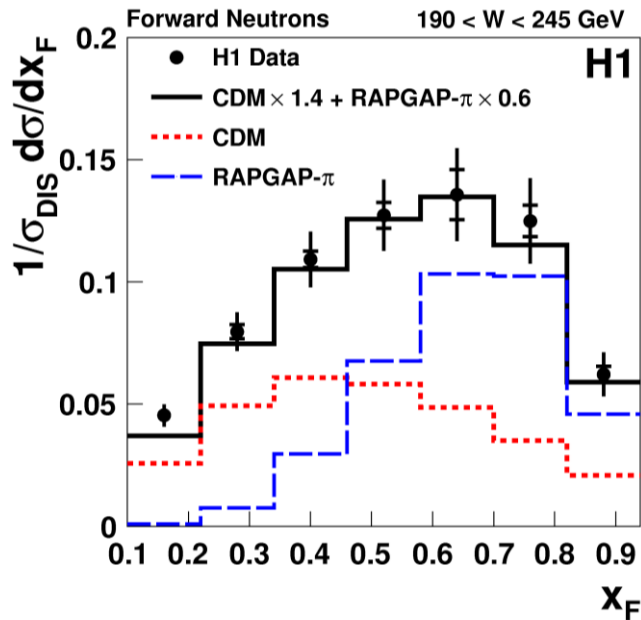


# Longitudinal spectrum of the forward neutron

- Limit by aperture, not corrected
  - yield at high  $z$  is enhanced
- Still a peak, corresponding to pion exchange

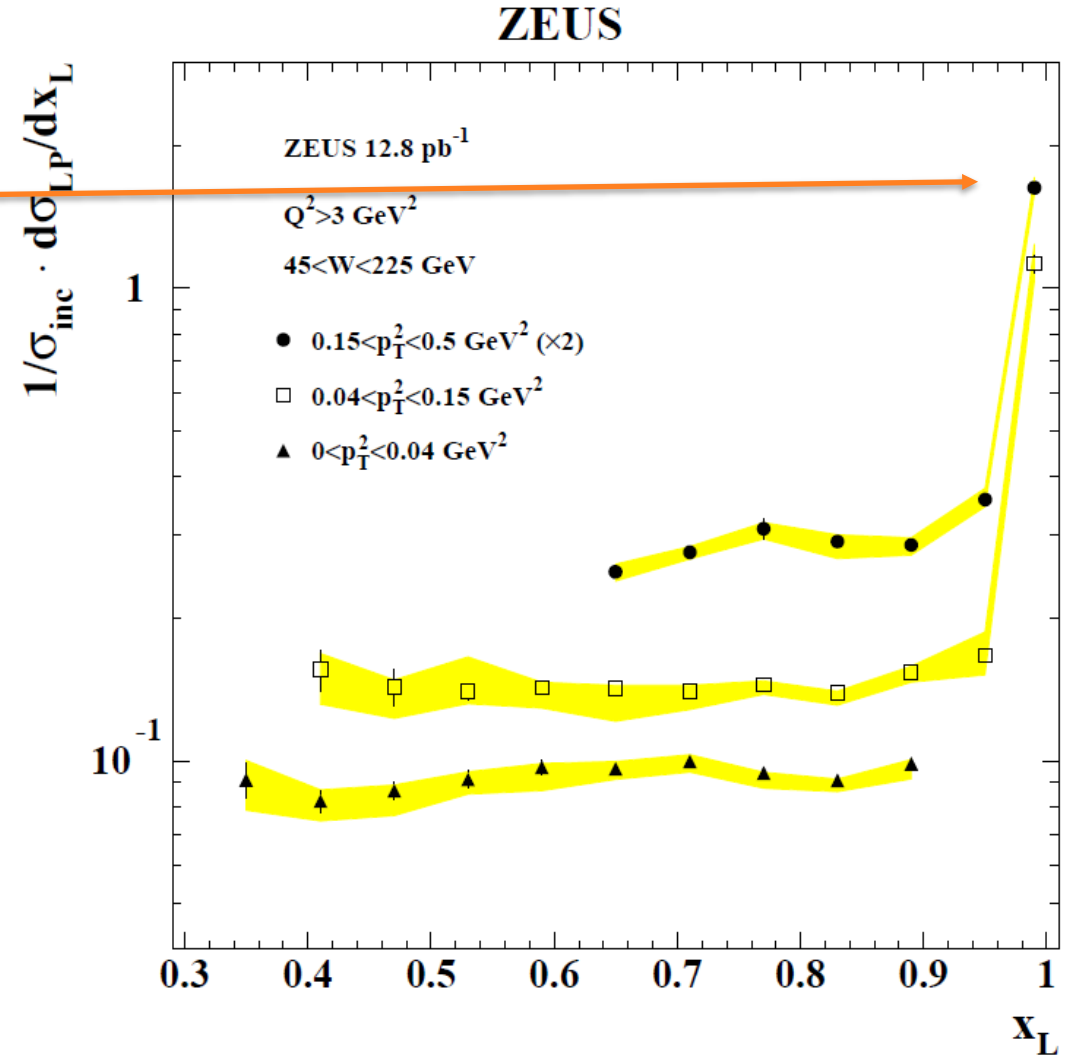
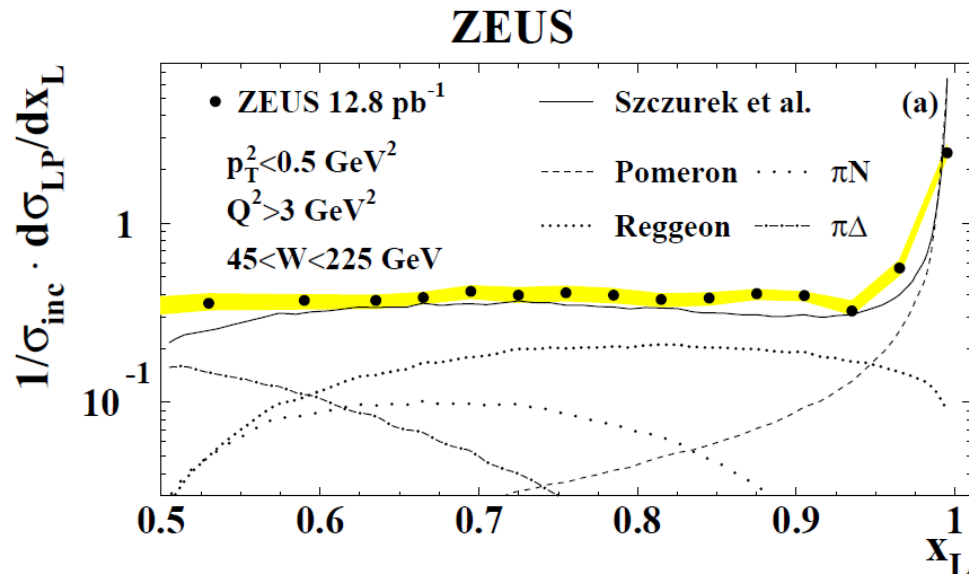
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Nucl.Phys.B637(2002)3



# Longitudinal spectrum of forward proton

- $x_L = p_Z^{LB} / p_{beam}$
- Very flat  
(except for diffractive peak)
- Explained by:
  - Limited fragmentation, or
  - particle exchange model  
(many Regge poles superimposed)



EIC should be able to study much more precisely

# **Meson exchange model and pion structure**

# The formalism of pion exchange

- The cross section =  $\sigma(\gamma^* \pi) \times$  (pion flux): **Regge factorisation**

$$\frac{d^2 \sigma(W^2, Q^2, z, t)}{dz dt} = f_{\pi/p}(z, t) \sigma_{\gamma^* \pi}((1-z)W^2, Q^2)$$

- The  $f_{\pi/p}(z, t)$  takes the following form in most of the models:

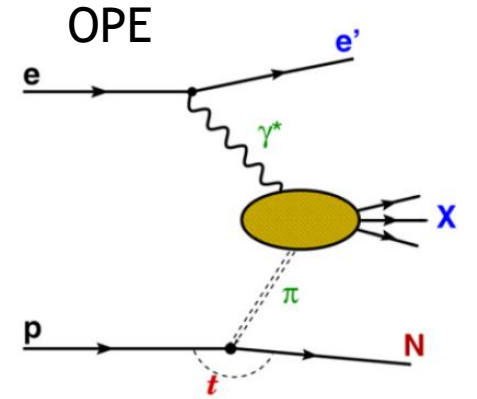
$$f_{\pi/p}(z, t) \propto -\frac{t}{(t - m_\pi^2)^2} (1-z)^{\alpha(t)} F^2(z, t)$$

$\alpha(t)$ : pion trajectory

$F(z, t)$ : introducing cutoff at the low- $t$  end

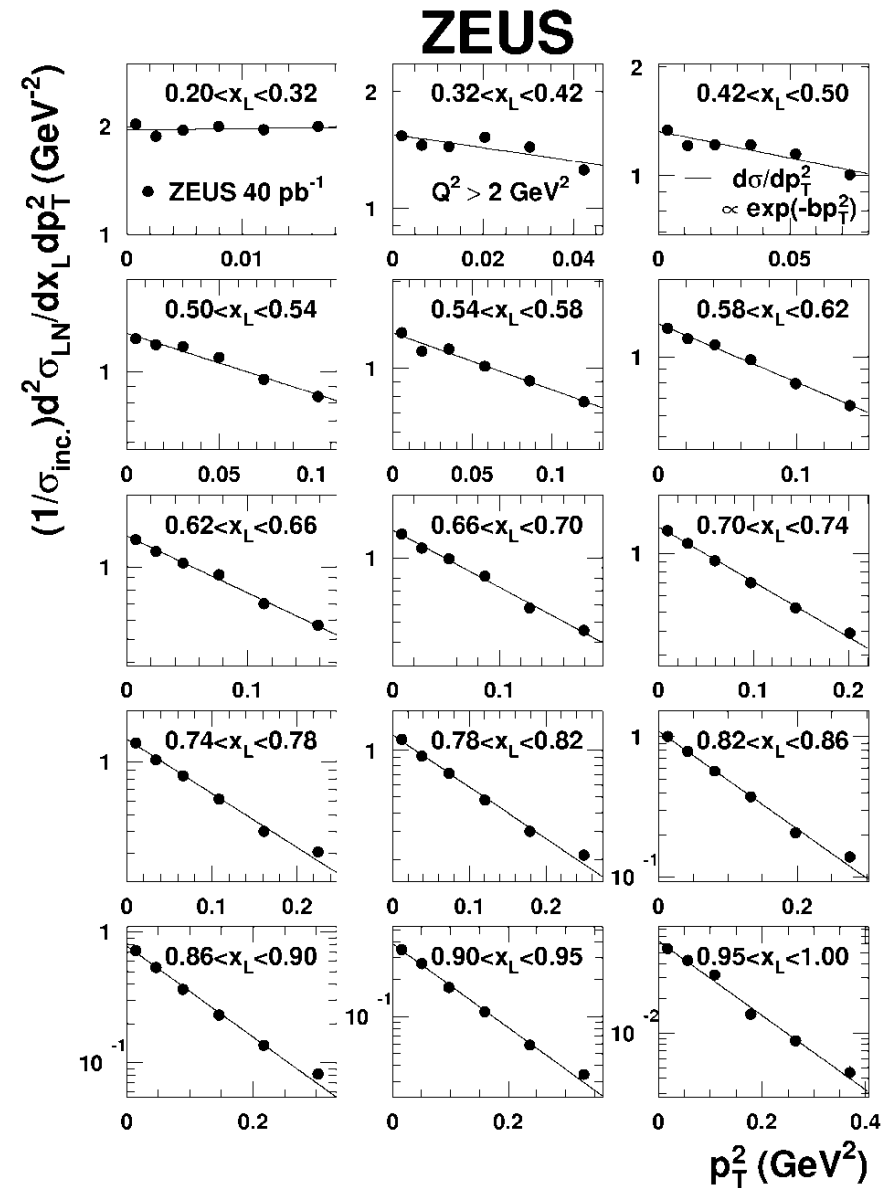
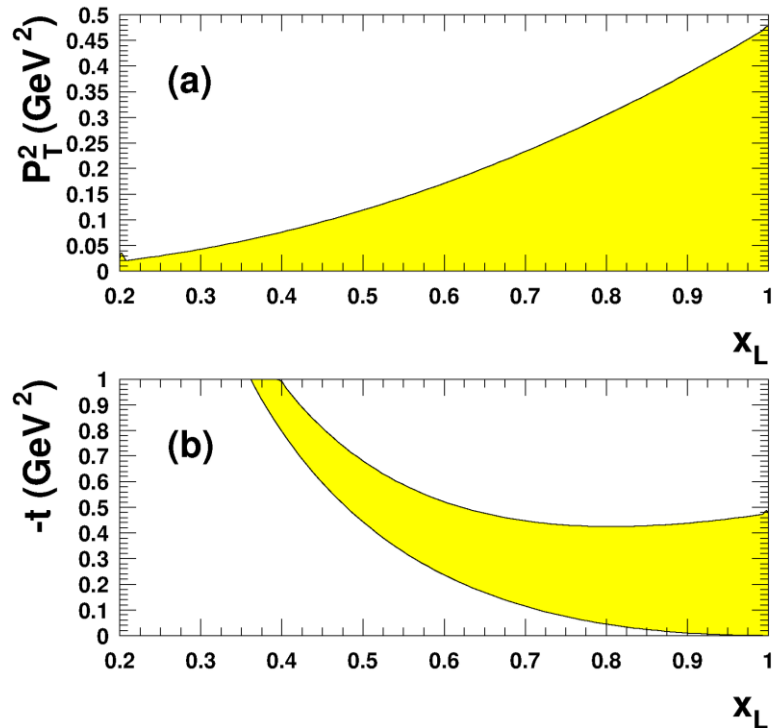
$F = 1$  (Bishari)

$F = \exp \frac{t - m_\pi^2}{\Lambda^2}$  (Holtman and others: generalized form of monopole  $F = \left( \frac{\Lambda^2 - m_\pi^2}{\Lambda^2 - t} \right)$ )



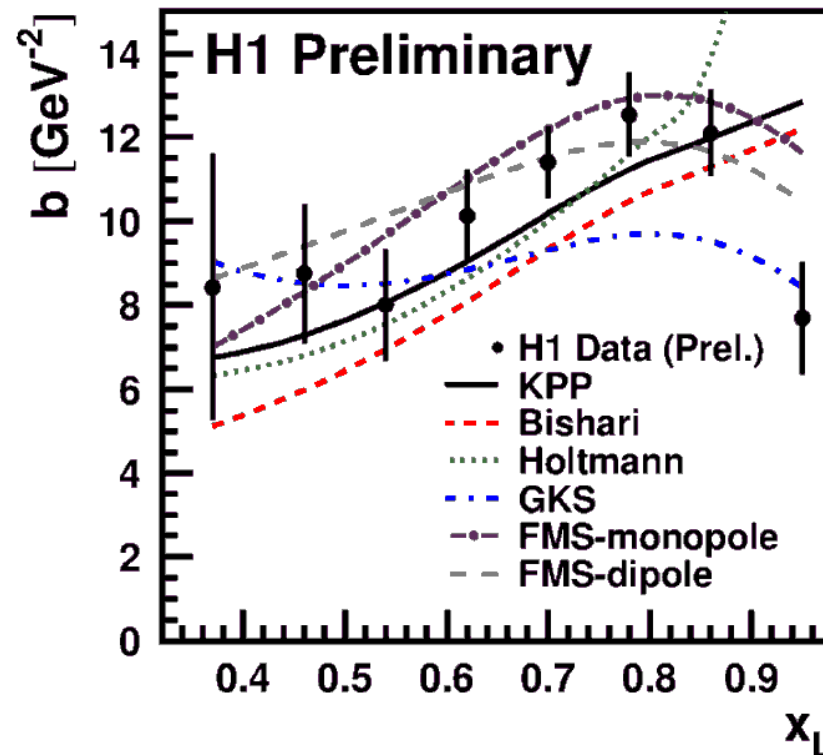
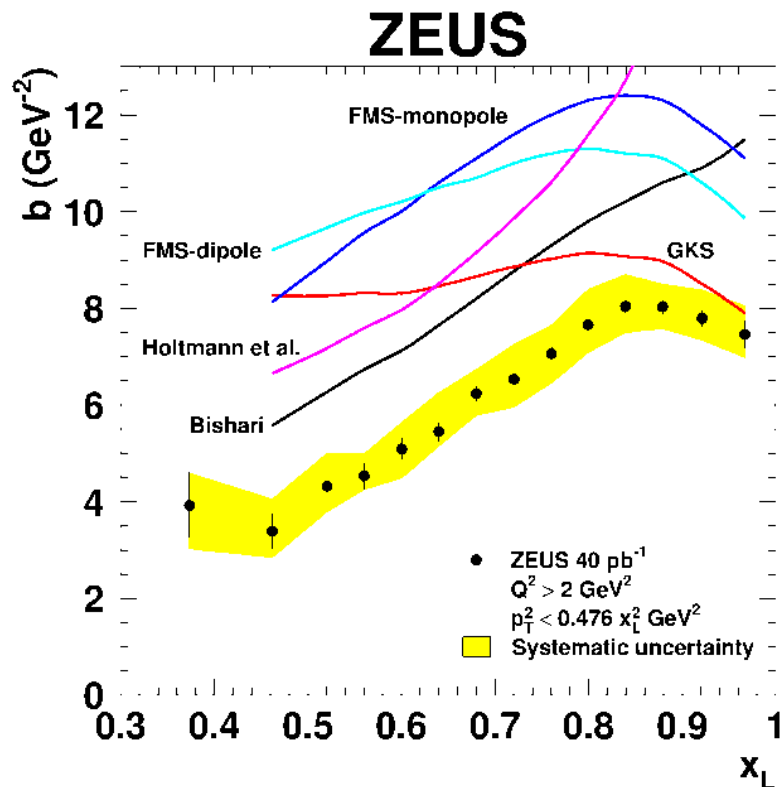
# $p_T$ spectrum of neutrons

- Exponential behaviour in low- $t$  regions
  - Seems to have rapid change in slope as a function of  $x_L = Z$
  - ... even through it is visually enhanced by the acceptance change vs  $x_L$



# Forward neutron: rich structure in b-slope

- Compared to various pion flux (e.g.  $\pi - n$  vertex factor shape)
  - Qualitatively in agreement with various models
  - Need to evaluate sub-leading components near  $z = 1$  and low- $z$ 
    - sub-leading meson exchange and/or non-diffraction (standard fragmentation)

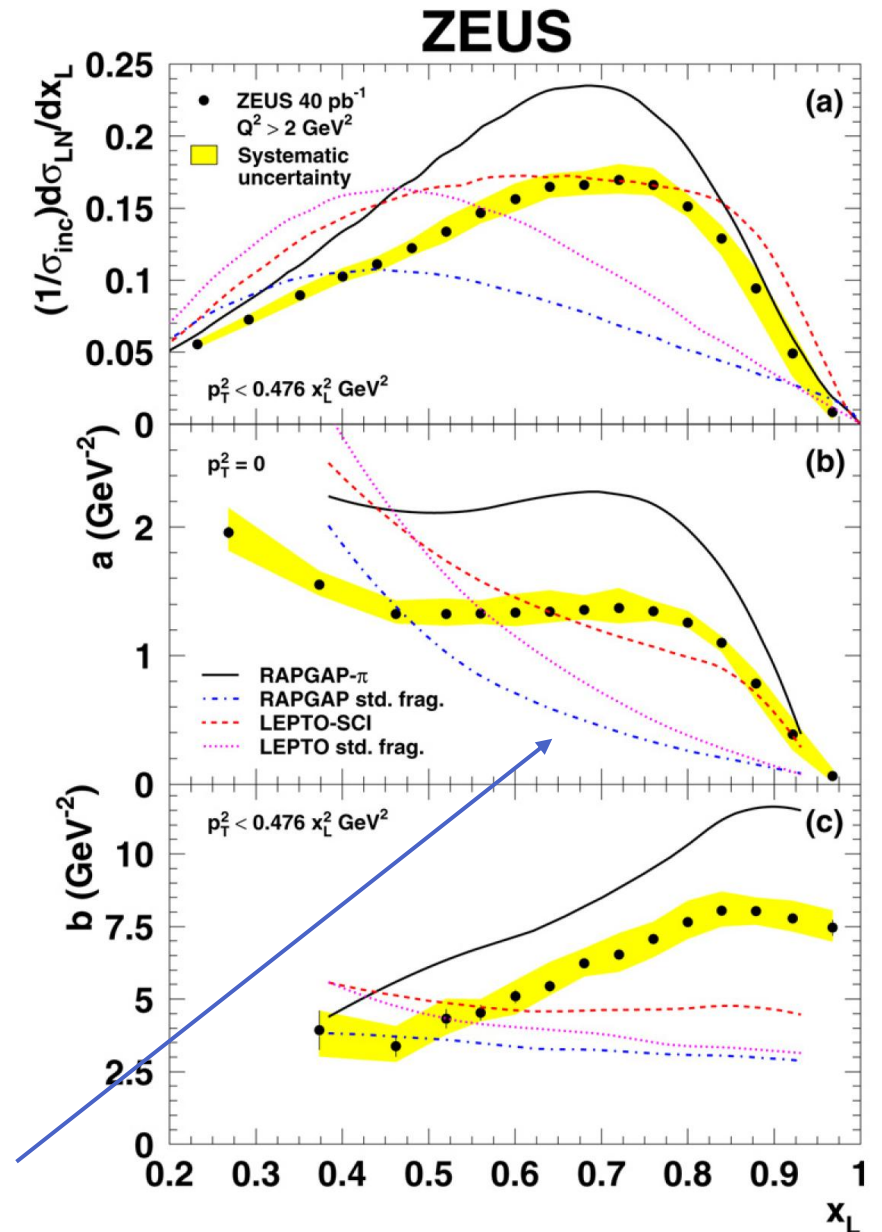


Need precise  
( $p_T, z$ ) spectrum  
at the EIC  
for measuring  
the pion flux

# Comparison with non-pion-exchange models

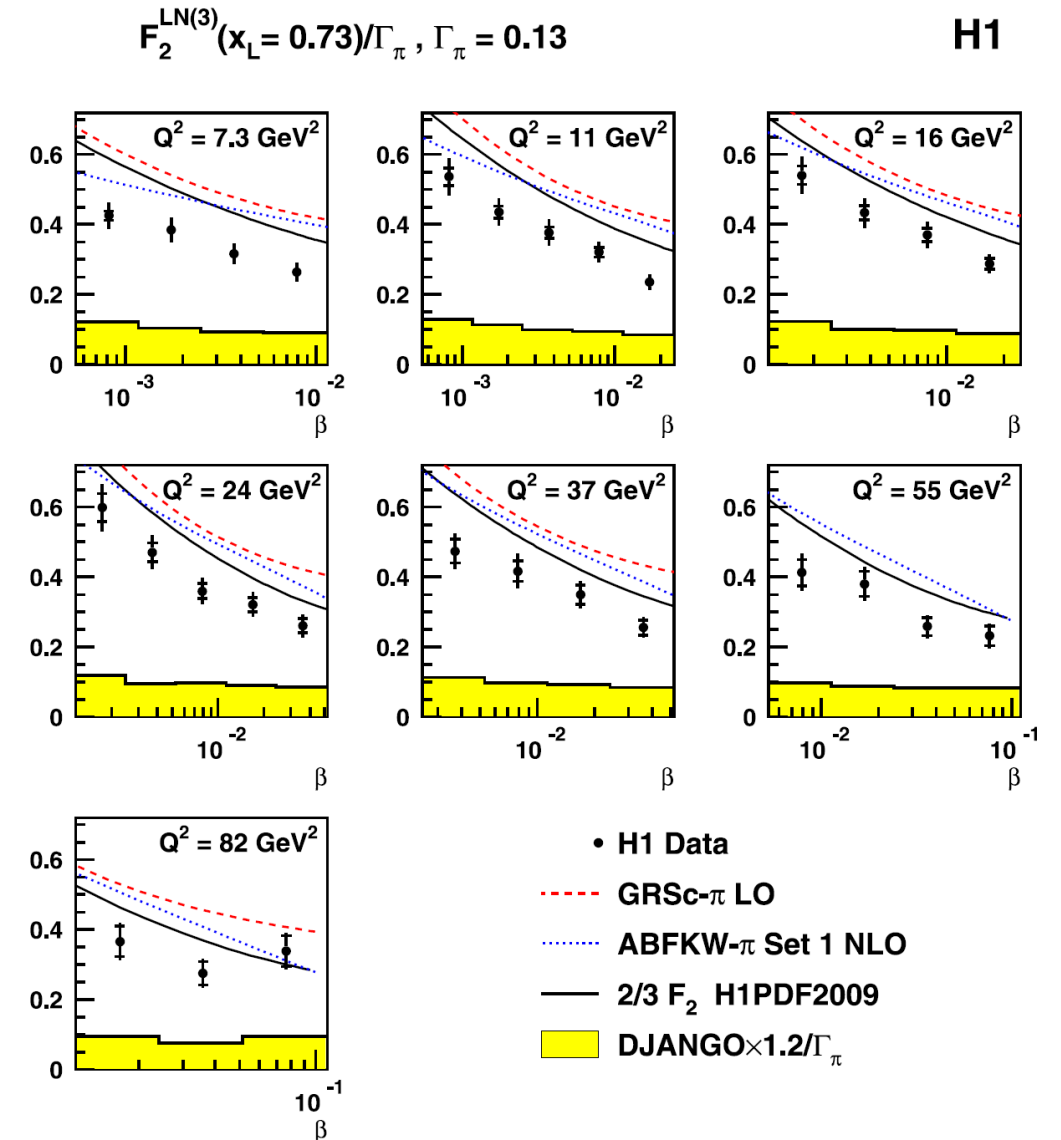
- SCI model: soft-colour interaction, i.e. generic t-channel exchange of soft partons
  - unified model for fragmentation + particle exchange (no double counting)
  - can explain the observed, peak somewhat
  - $p_T$  spectrum too broad
- Models based on fragmentation only
  - "pion" peak missing
- How do we bridge the pion exchange model and the standard fragmentation at middle/low  $z$ ?

Non-diffractive process may be background to OPE  
Need to go to high- $z$  to avoid that



# HERA extraction of the pion structure function

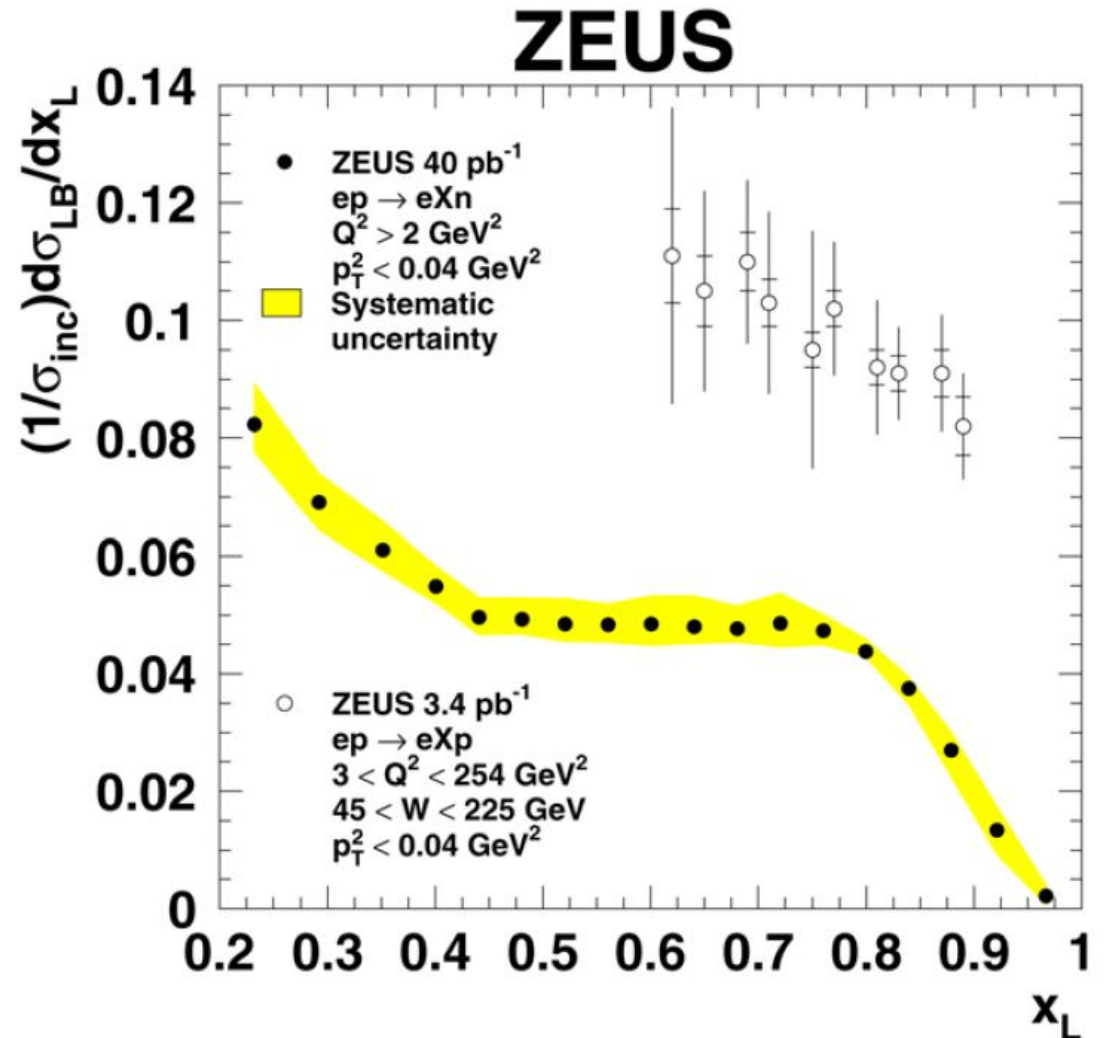
- $x_\pi = \beta \equiv x_{Bj}/(1 - x_L)$   
 $x_L \equiv z$ : neutron momentum fraction
  - the shape of the SF for proton and pion should be the same if  $F_2^p \propto x^{-\lambda}$  with constant  $\lambda$  ( $1 - z$  just rescales  $x$  for  $x_\pi$ )
  - This holds quite well: see 2/3 F2 H1PDF2009, which is the proton structure function!
- The absolute value of the SF is smaller
  - ZEUS similar (even smaller)
  - What is the reason?
- EIC: high- $x_\pi$  SF possible with high statistics



# Why are the pion SFs too small?

- Naïve isovector exchange:  
neutrons are more than proton  
in the final state of hadron-hadron  
collisions
- This was not the case at HERA!  
More protons there for very forward  
range  $p_T^2 < 0.04 \text{ GeV}^2$
- Where are these neutrons?
  - or are there more protons  
than expected?

This must be the first thing to check!



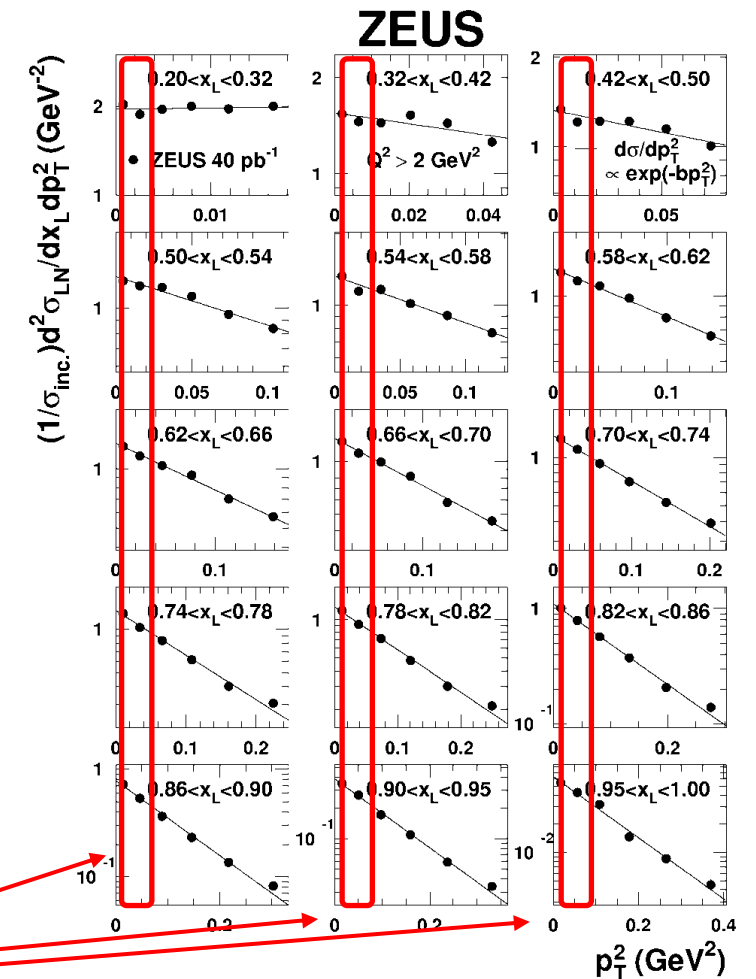
# To be away from the flux uncertainty ...

- The F-term of the pion flux is unity

near  $t = t_{min}$  (e.g.  $F = \exp \frac{t - m_\pi^2}{\Lambda^2}$ )

- Smaller uncertainty in the flux
- Also the exchanged pion is on-shell-like there
- Need measure near the  $t$  threshold = very low  $p_T$
- For that, we need:
  - neutron position resolution
  - very precise determination of zero degree
  - low beam spread

Nuclear Physics B 776 (2007) 1-37

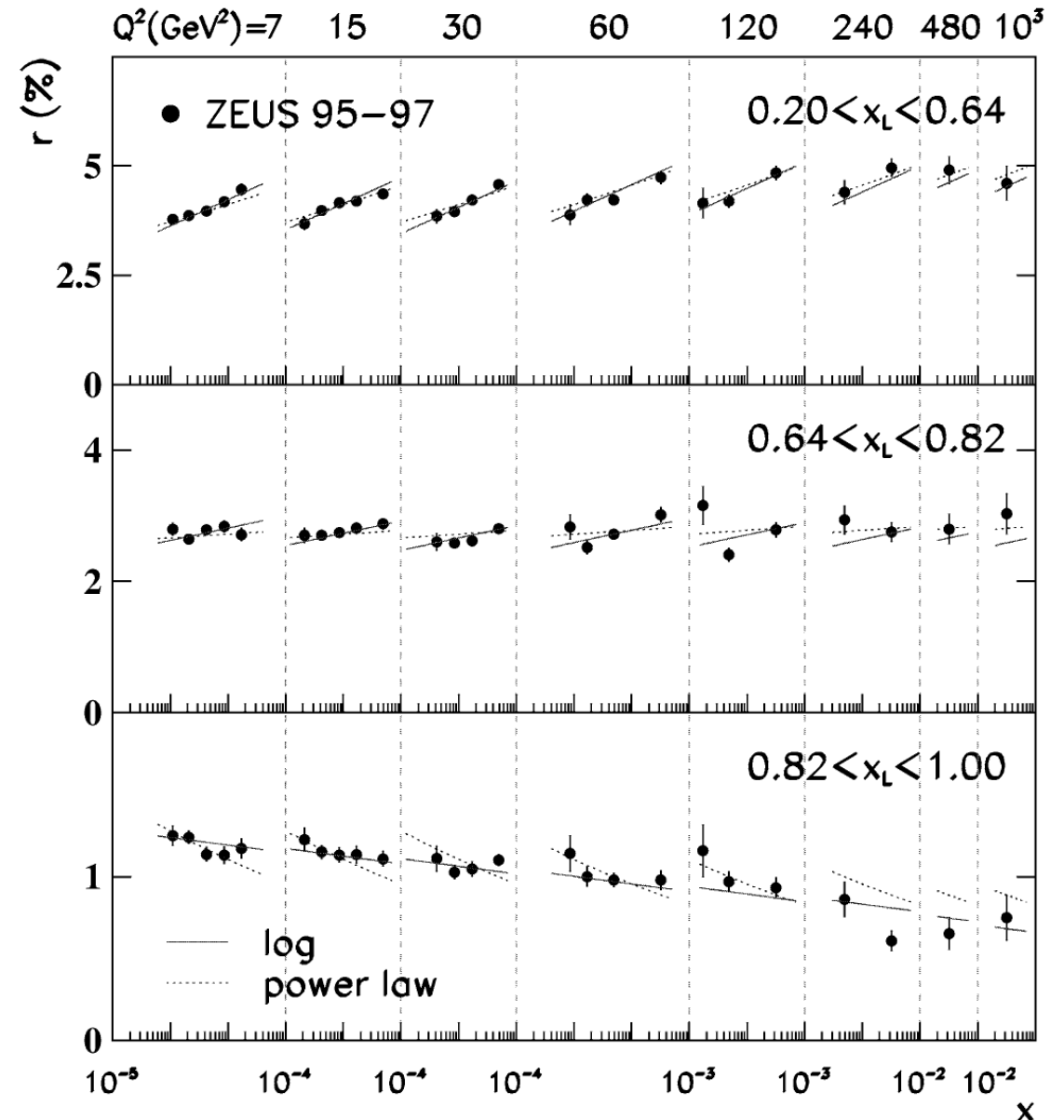


Focusing on these points!

# **Possible factorisation breaking in baryon production**

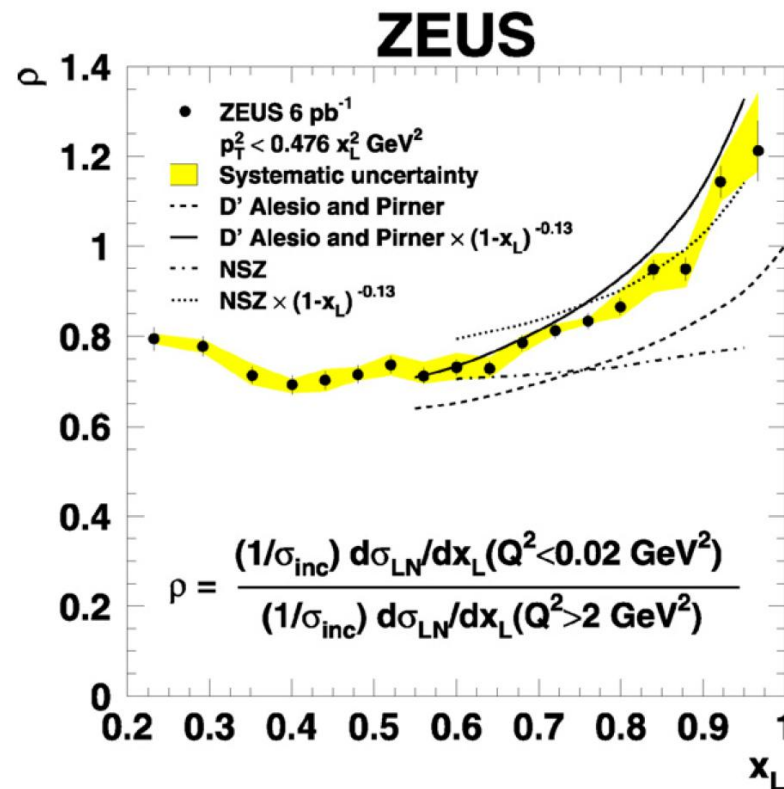
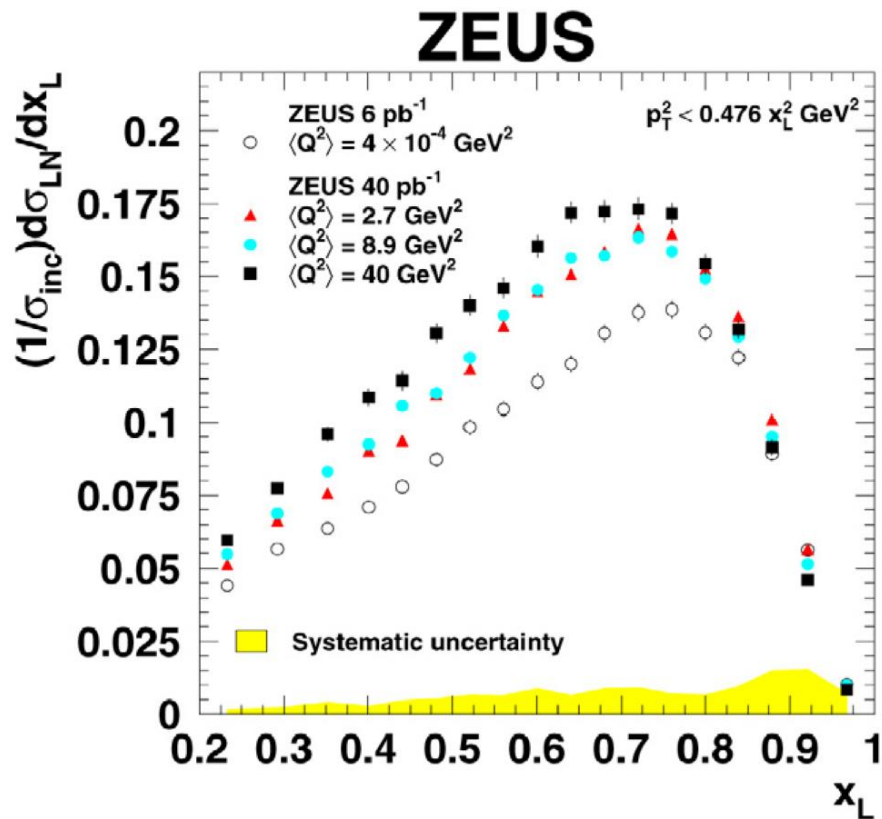
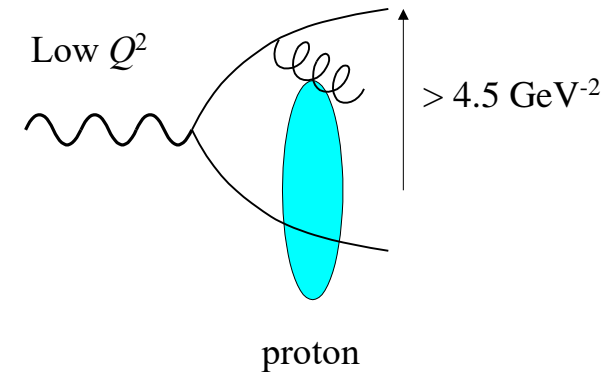
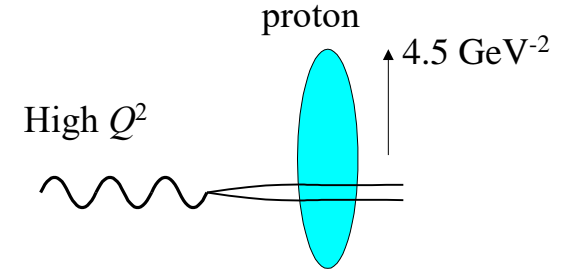
# Fraction of events with forward neutrons: $x$ -dependence?

- Leading neutron events have very high statistics (> 10% of DIS)
  - very detailed kinematic dependence can be studied
  - factorization or not?  
it seems we have logarithmic factorization breaking  
It was not conclusive within DIS



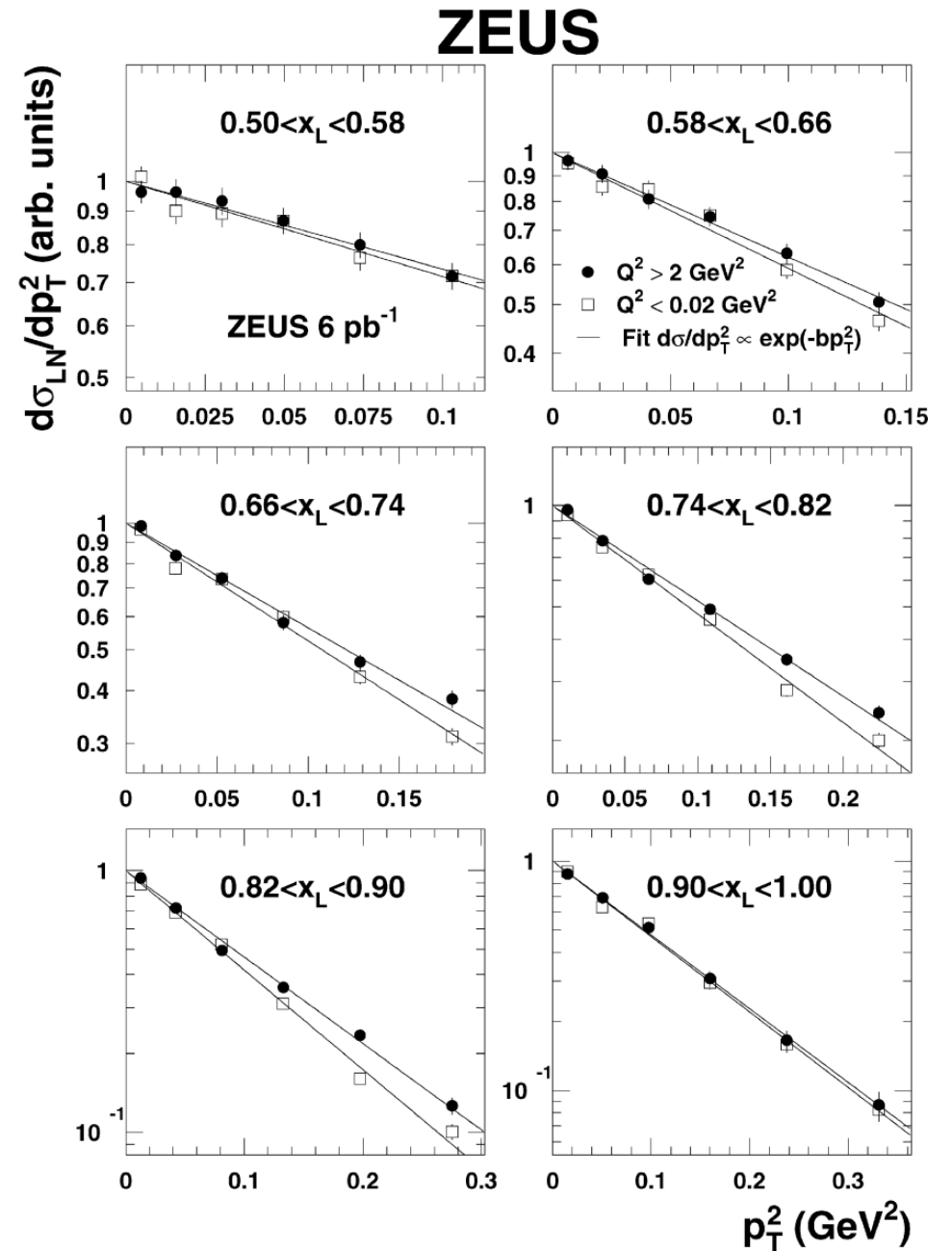
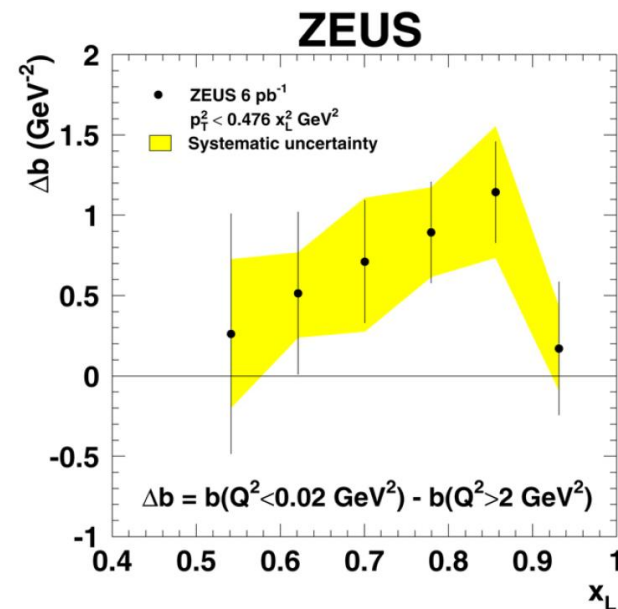
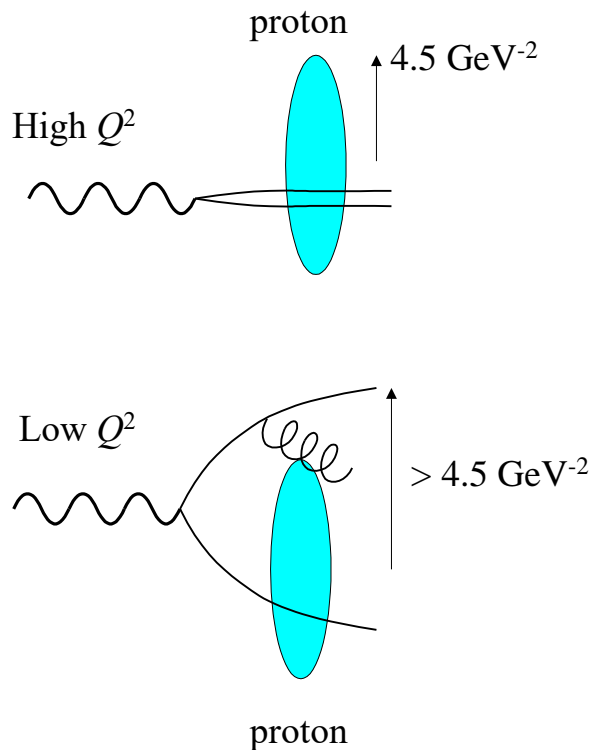
# In $Q^2$ : (photoproduction) / (DIS $Q^2 > 2 \text{ GeV}^2$ )

- Mild suppression in photoproduction = large photon
  - implying 20-30% absorption (re-scattering)
- clear  $x_L = z$  dependence on the ratio photoproduction / DIS



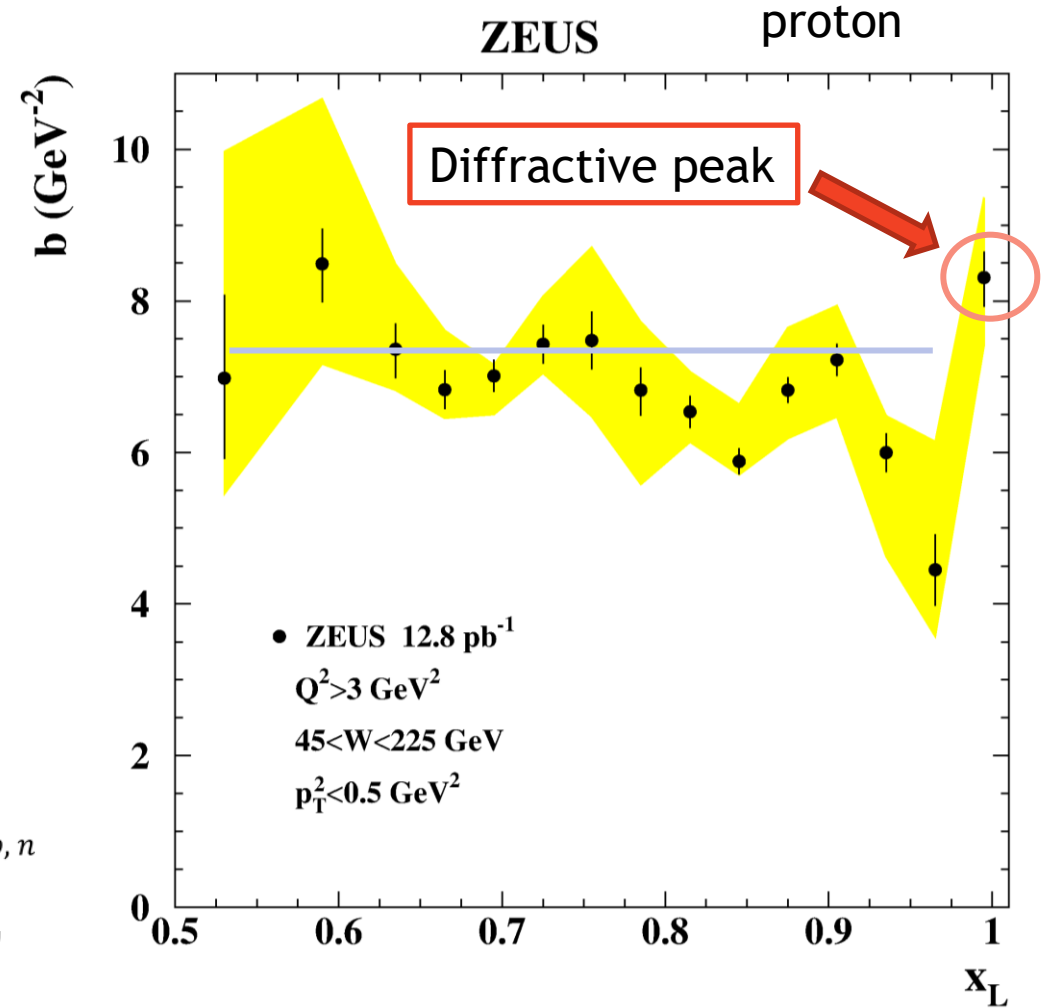
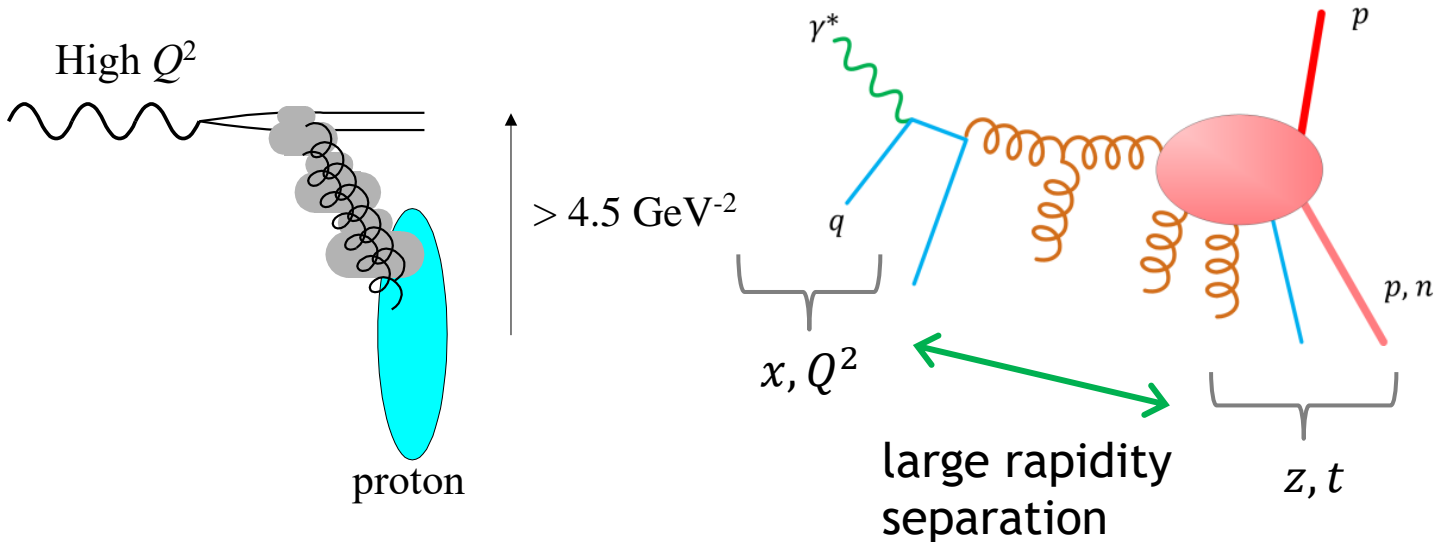
# Slope $Q^2 \simeq 0$ vs $Q^2 > 2 \text{ GeV}^2$

- We see factorisation breaking in the slope too
  - Absorption or rescattering of “small configuration” = large overlap



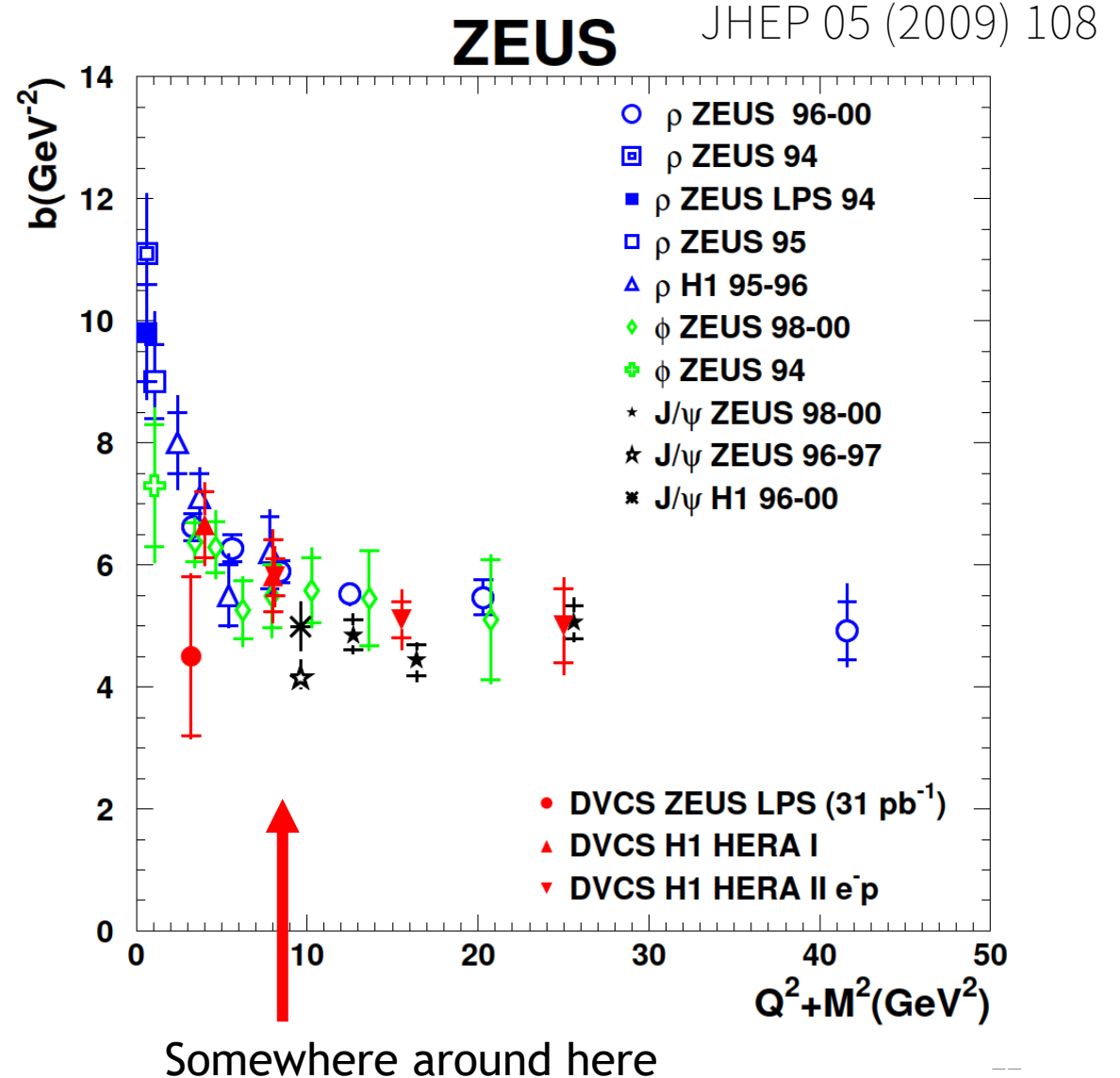
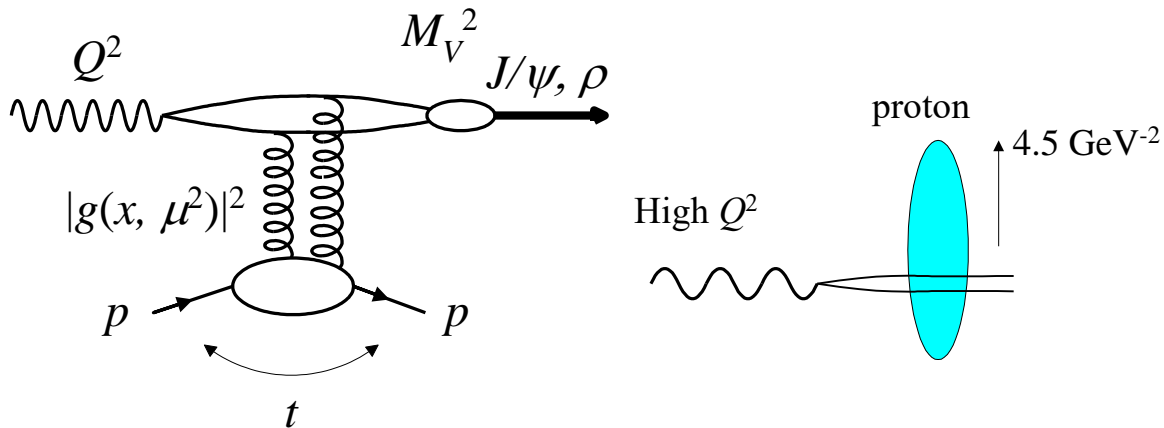
# B-slope of leading baryons at HERA

- Proton: flat at  $b \sim 7 \text{ GeV}^{-2}$  ( $\sigma \propto e^{-bp_T^2}$ )
  - Larger than the proton size
  - The vertex is somewhat soft = peripheral but not too soft either (unlike elastic  $pp$ )
  - Large rapidity separation between the dipole and the forward system at high energies may smear the transverse size



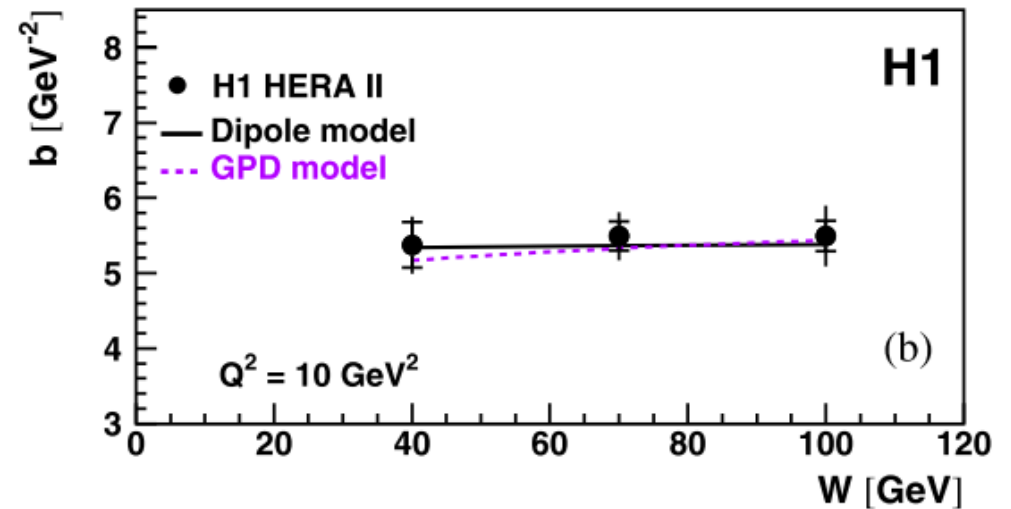
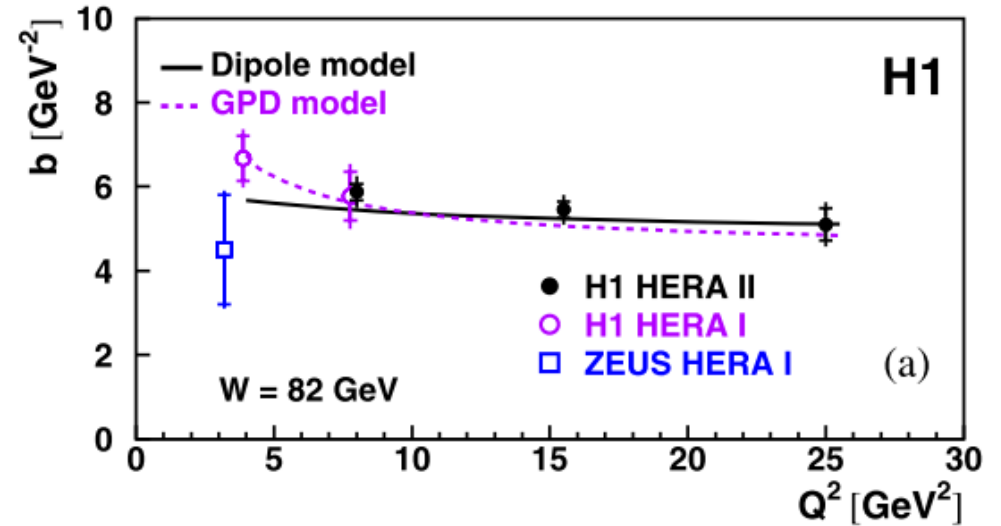
# cf. $b$ -slope for vector meson electroproduction

- Where does the dipole size become negligible?
  - Seems to have universal behaviour as a function of  $Q^2 + M_V^2$
  - $M_V$ : vector meson mass
- Asymptotically approaches to  $b \simeq 5 \text{ GeV}^{-2}$  around  $Q^2 + M_V^2 \simeq 10 \text{ GeV}^2$ 
  - Slightly smaller than the forward baryon  $b$ -slope ( $7 \text{ GeV}^{-2}$ )



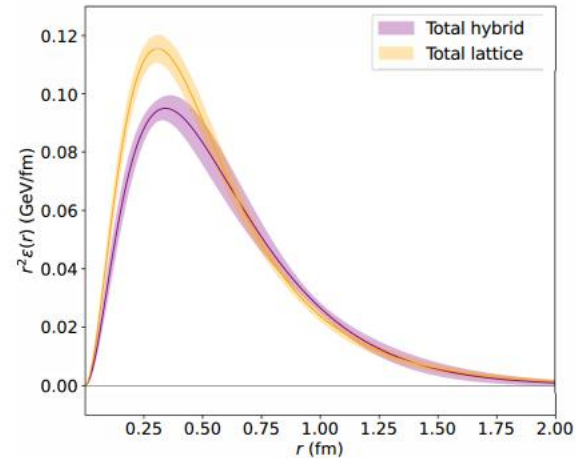
# DVCS $b$ -slope

- GPD model seems to have been describing the increase in  $b$  when the hard scale is absent

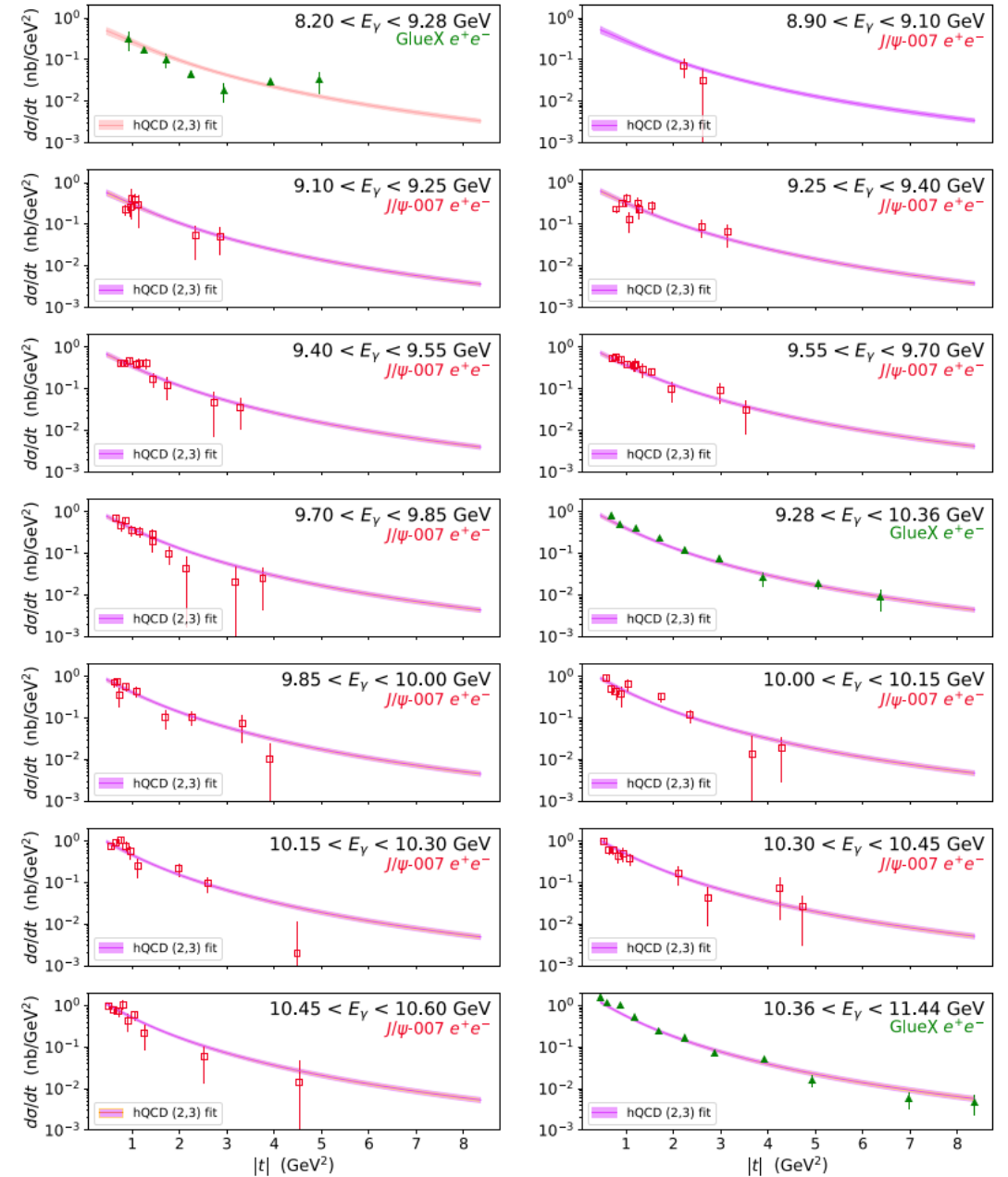


# JLab t-slope for exclusive $J/\psi$ production

- Exclusive photoproduction  $\gamma p \rightarrow J/\psi p$  from  $J/\psi$ -007 and GlueX experiments
  - Compilation in arXiv:2512.15064, by M. Boer et al.
  - Interpreted as energy-momentum tensor, then Fourier-transformed to extract the shape in radius



**Fig. 12** The total mass density where we combined the gluon extraction FFs from Fig. 11 and the quark FFs evaluated in lattice [219].



# **Towards pion structure: day-1 measurements at the EIC**

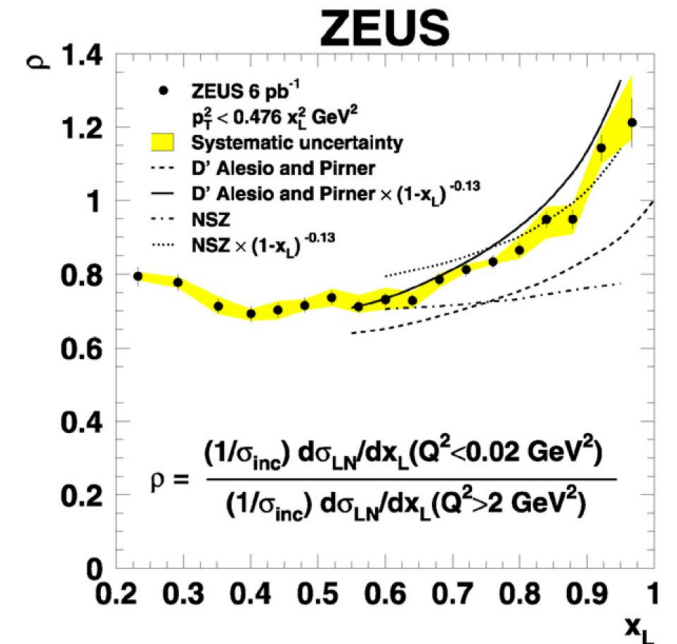
# The day-1 measurement we need to do for neutrons

- $p_T$  spectrum vs  $z$  is absolutely important
  - Very good  $p_T$  (position) resolution detector for neutrons
  - "beam spot" = zero-degree determination
  - need for runs with small beam spread in  $p_X, p_Y$ : **large  $\beta^*$  = low lumi runs**
- Need to measure  $(z, p_T)$  spectrum in sufficiently narrow  $(x, Q^2)$  bins

– to first test the Regge factorisation Ansatz:

$$\frac{d\sigma(W^2, Q^2, z, t)}{dzdt} = f_{\pi/p}(z, t) \sigma_{\gamma^* \pi}((1-z)W^2, Q^2)$$

- At HERA (was  $40 \text{ pb}^{-1}$  data): only  $Q^2 \simeq 0$  vs  $Q^2 > 1 \text{ GeV}^2$ , or fine bins in  $Q^2$  but not in  $(z, t)$ 
  - Need  $\gg 100 \text{ pb}^{-1}$  to map  $(x, Q^2, z, t)$  but may not need  $\gg 1 \text{ fb}^{-1}$
- Need photoproduction data: may need ZDC coincidence for saving photoproduction data if it is not saved by default in the current streaming scheme



# Limiting fragmentation and "Regge" factorization

- **Limiting fragmentation:** the LN variables  $(z, t)$  are independent of  $(x, Q^2)$ 
  - large rapidity separation, (almost) no information propagated

- **Regge factorisation:** the cross section is the product of the (particle flux)  $\times \sigma(\gamma$ -particle)

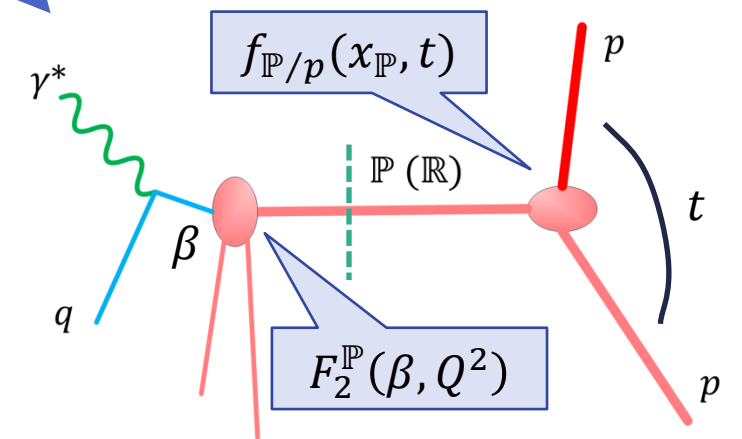
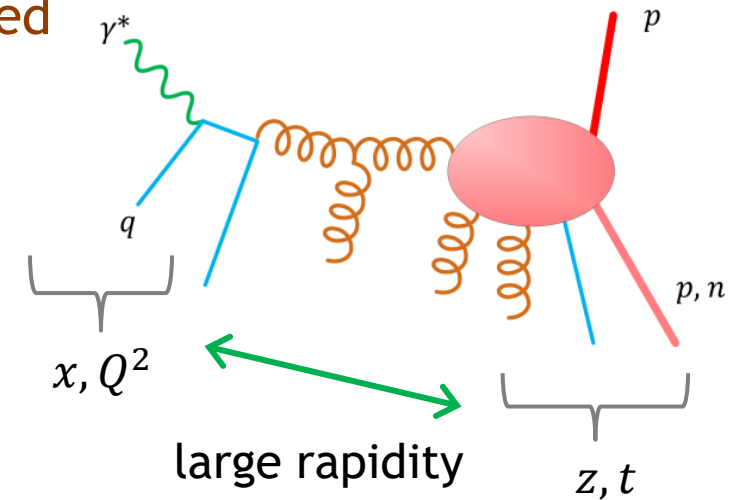
$$\frac{d\sigma(W^2, Q^2, z, t)}{dzdt} = f_{\pi/p}(z, t) \sigma_{\gamma^* \pi}((1-z)W^2, Q^2)$$

- **For both models, factorisation holds between photon vertex  $(x, Q^2)$  and baryon vertex  $(x_L, t)$** 
  - The independence does not distinguish the two models

- Proof of the Regge factorization:

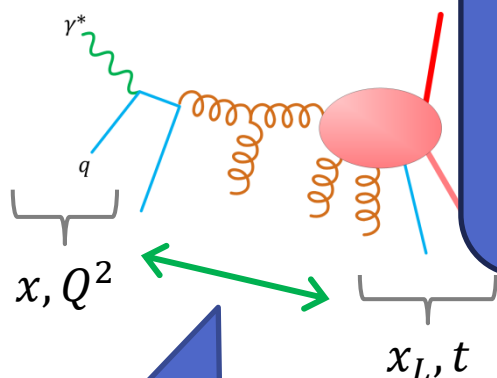
$\sigma_{\gamma^* \pi}$  is independent of the soft vertex variables  $z$  and  $t$

- E.g. pion structure function should be independent of  $x_L$  values: we could measure in a few bins in  $x_L$



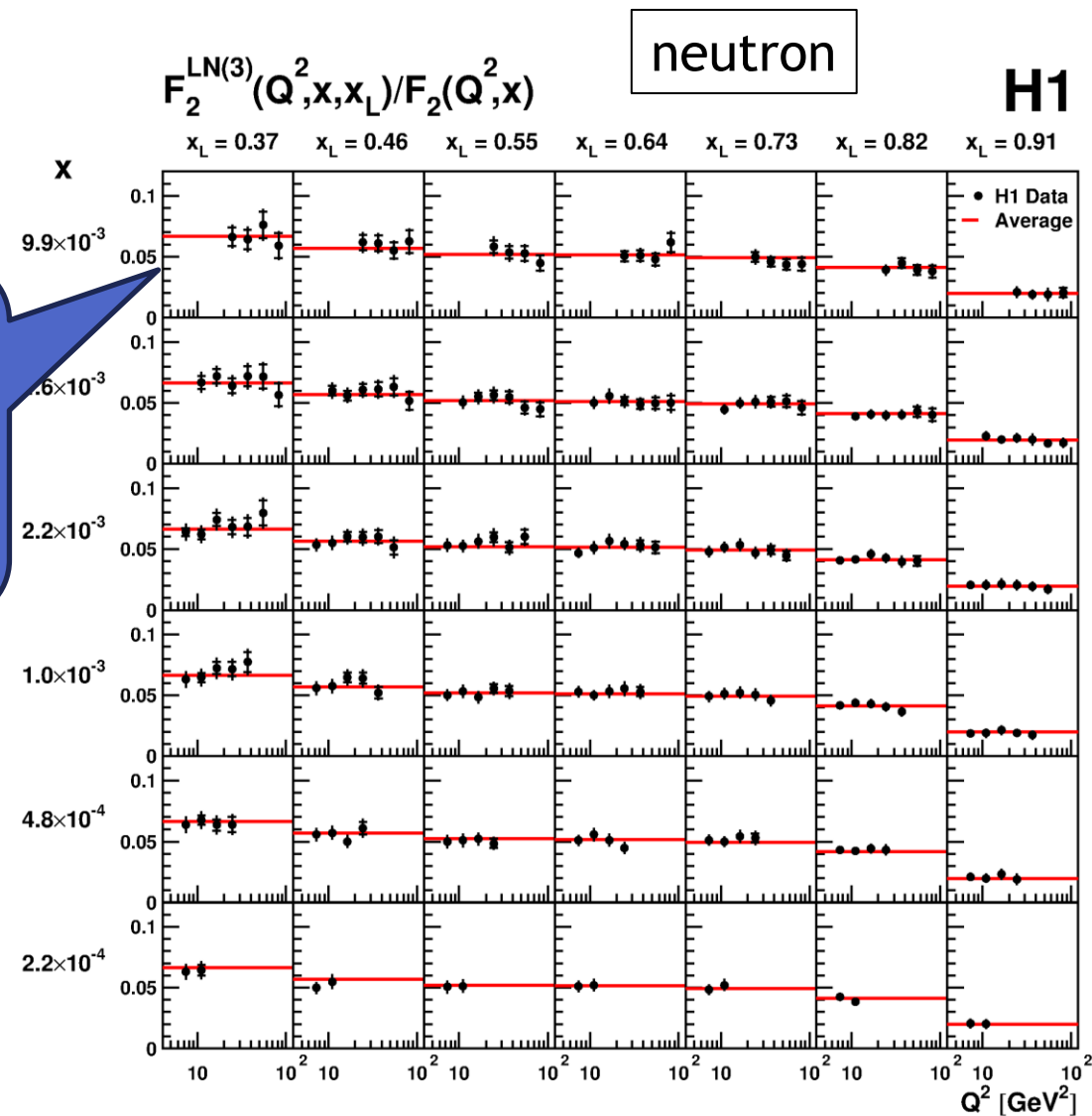
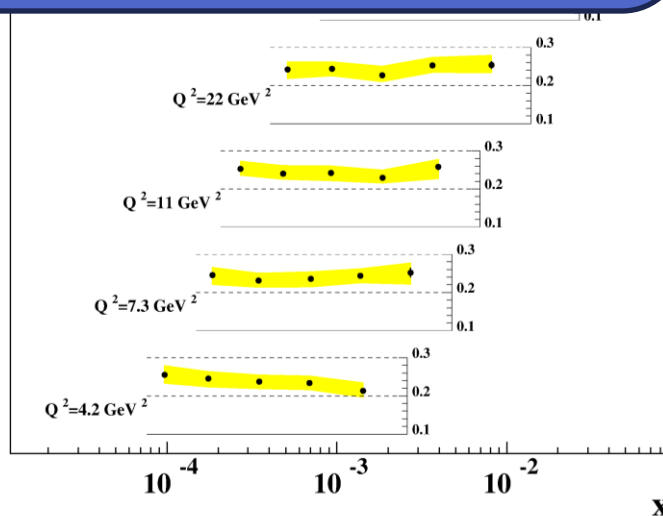
# Does the baryon talk to the virtual photon?

- The answer is basically no!
  - No strong yield dependence on  $x$ ,  $Q^2$
  - “limiting fragmentation”



The trend of the “height of the red lines” vs  $x_L$  (horizontal boxes) should be independent of  $x$  (vertical boxes) and  $Q^2$  ( $x$  axis within the box)

The photon and LN vertices are way separated in rapidity

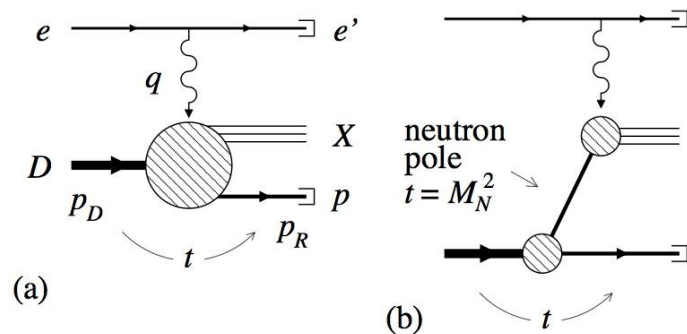


fraction of events with proton is flat in  $x$

... and flat in  $Q^2$  for given  $(x, z)$  bins

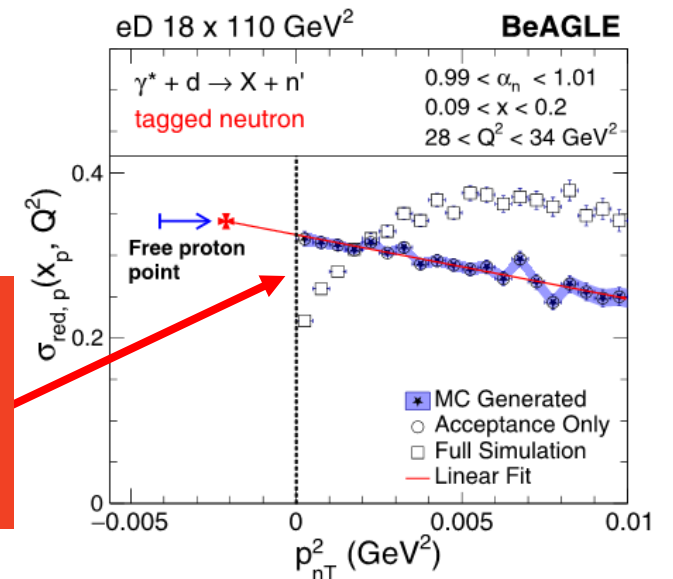
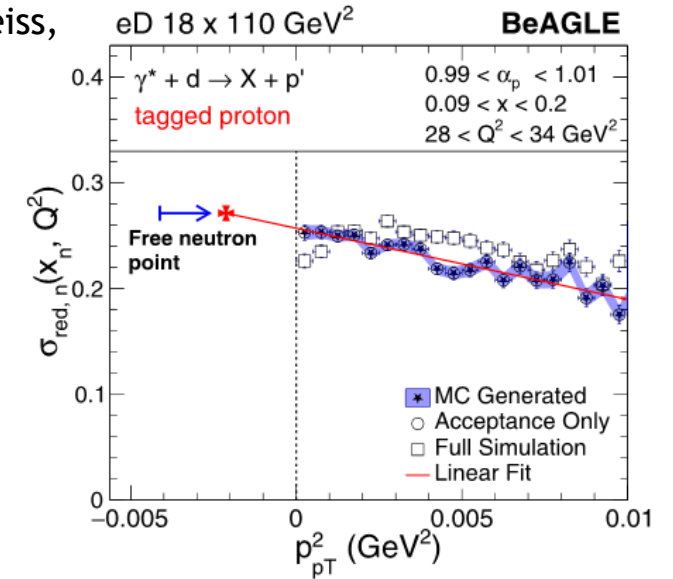
# Example of $t$ reconstruction: $eD$ scattering

- Proton-tagged  $eD$  scattering
  - $e(p + n) \rightarrow en + p$  DIS for neutron
- Neutron-tagged  $eD$  scattering
  - $e(p + n) \rightarrow en + p$  DIS for proton, which can be compared to  $ep$  scattering to understand the  $t$  effect
- Way to understand nuclear (EMC) effect
  - Or short-range correlation (SRC), by comparing ( $ep + en$ ) with untagged  $eD$  DIS



large smearing in measuring the recoil  $t$  due to finite resolution: severer for  $n$  than  $p$

A. Jentsch, Z. Tu, C. Weiss, PRC 104, 065205 (2021)



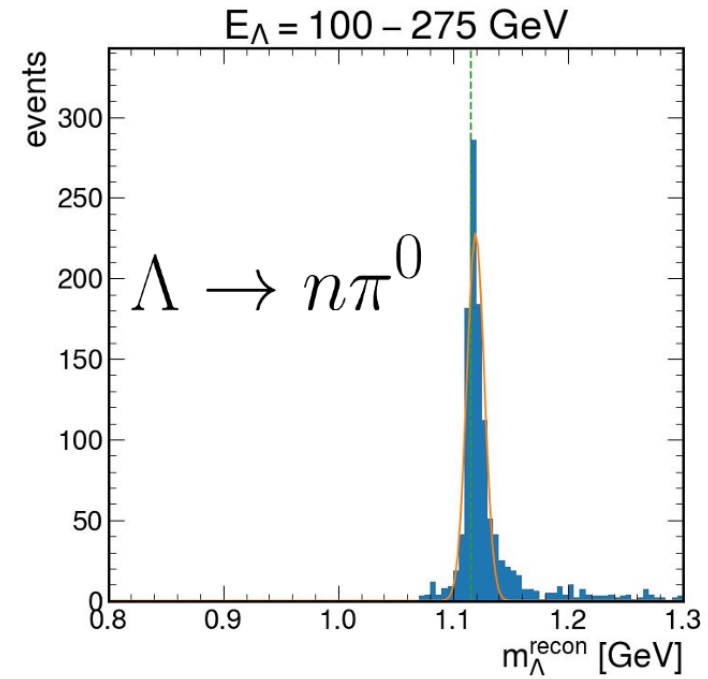
# Summary

- Old experimental results for forward baryon production are reviewed
- The neutron production at HERA supports the presence of the Sullivan process, allowing us to extract the meson structure
- A few remarks, however:
  - We need good transverse momentum resolution if we like to evade the meson flux uncertainty
  - We seem to have had suppression in forward neutron production  
We may have some mechanism to “absorb” the forward baryon even for high- $Q^2$  DIS
- ZDC position and incident angle resolutions are important
  - Also for reconstructing the  $\Lambda$  decay vertex through  $\gamma$  direction: we also need photon angular resolution

# **Backup: kaon and pion structure**

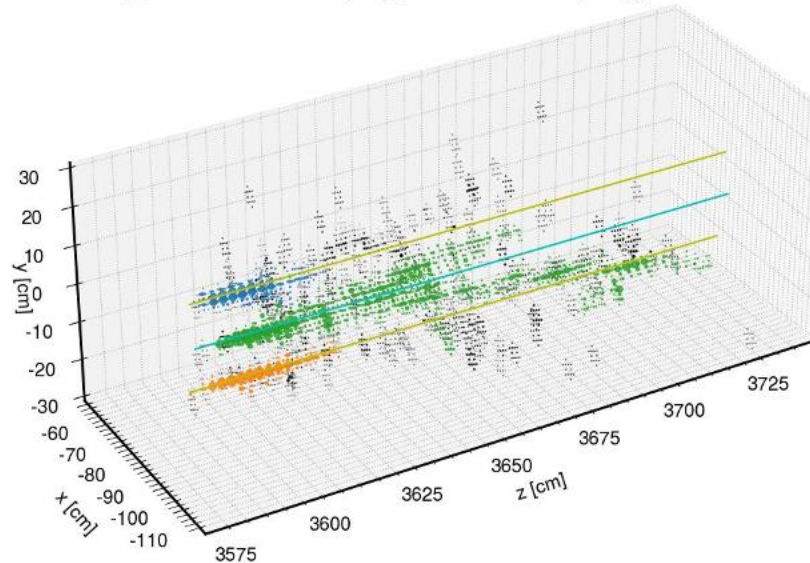
# Kaon structure: measurement

- using  $\gamma^* p \rightarrow (\gamma^* K^+) + \Lambda^0$ ,  $\Lambda^0 \rightarrow n\pi^0$  (35.9%)
  - all decay products go to ZDC if:
    - $\pi^0$  is energetic ( $\gtrsim 10$  GeV) or decay close to ZDC
  - precise determination of the **direction of 2  $\gamma$ 's** from  $\pi^0$  allows to reconstruct  $\Lambda$  vertex, hence **more precise  $m_\Lambda$**  to reject background

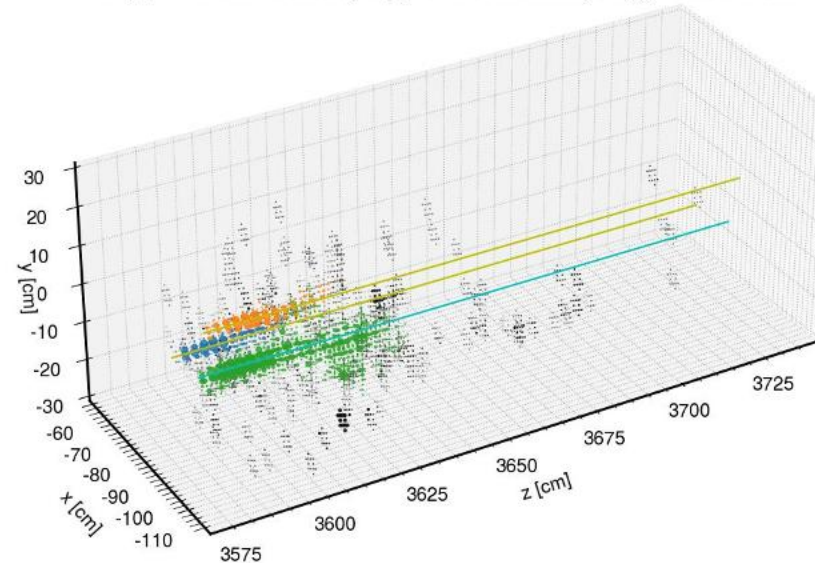


Credit: Sebouh Paul

$E_\Lambda = 100$  GeV,  $\theta_\Lambda = 2.1$  mrad,  $z_{\text{vtx}} = 15.0$  m



$E_\Lambda = 100$  GeV,  $\theta_\Lambda = 2.4$  mrad,  $z_{\text{vtx}} = 17.3$  m

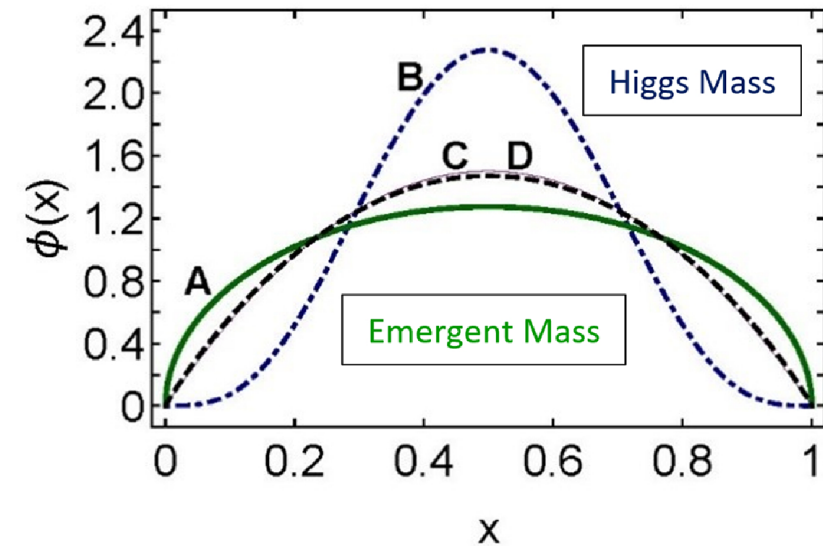


M. Arratia  
ePIC meeting, July 2024

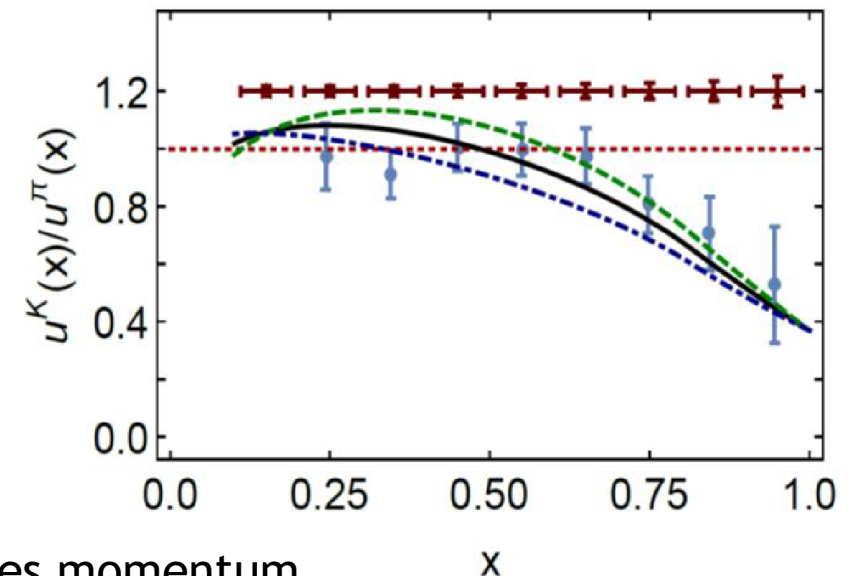
# Why Kaon structure?

- Kaon mass is more dominated by the Higgs mass of strange quarks than that for pions
  - Fraction of mass from kinetic energy of partons is small
- The Fermi motion may be smaller
  - the valence quark momentum is less spread
- Important ingredient of the origin of the mass
- The difference from pion is large at  $x_{\pi/K} \rightarrow 1$

Ratio of parton densities  $K/\pi$  slightly enhanced at middle  $x$ , but less pronounced if glue takes momentum

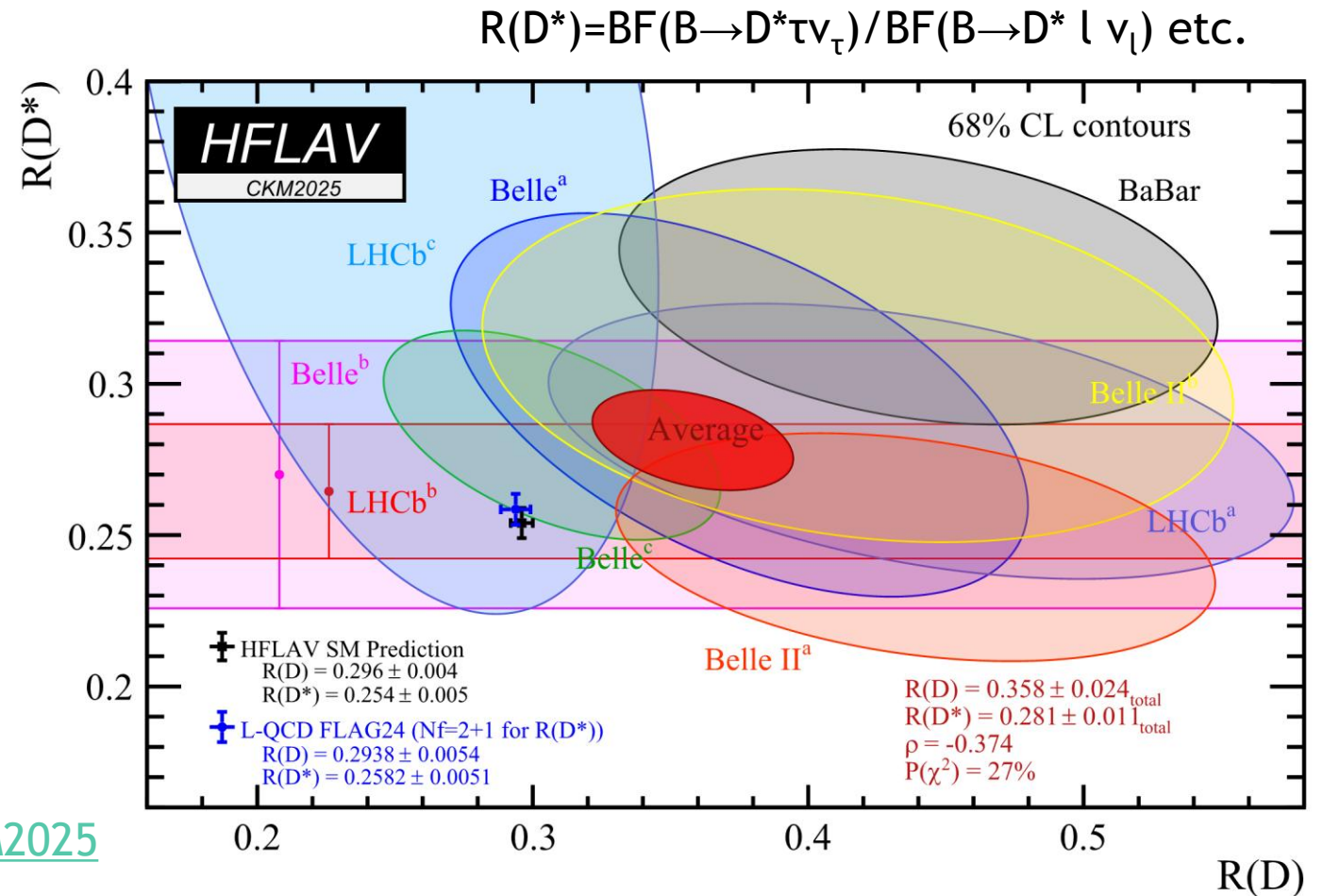


Parton distribution amplitude, A: pion, B:  $\eta_c$ , D: "heavy pion":  $m_{u,d}$  replaced by  $m_s$



# Meson structure and HEP

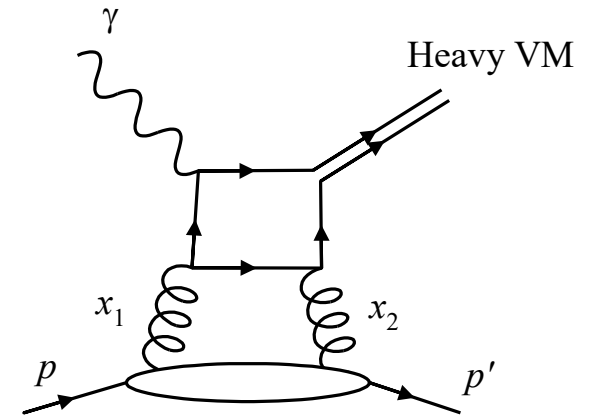
- Heavy quark meson structure needs to be known better
  - there may be hidden uncertainty in the flavor anomaly
- HEP theorists are counting for lattice QCD
  - experimental data on hadron structure with different quark mass may give insight



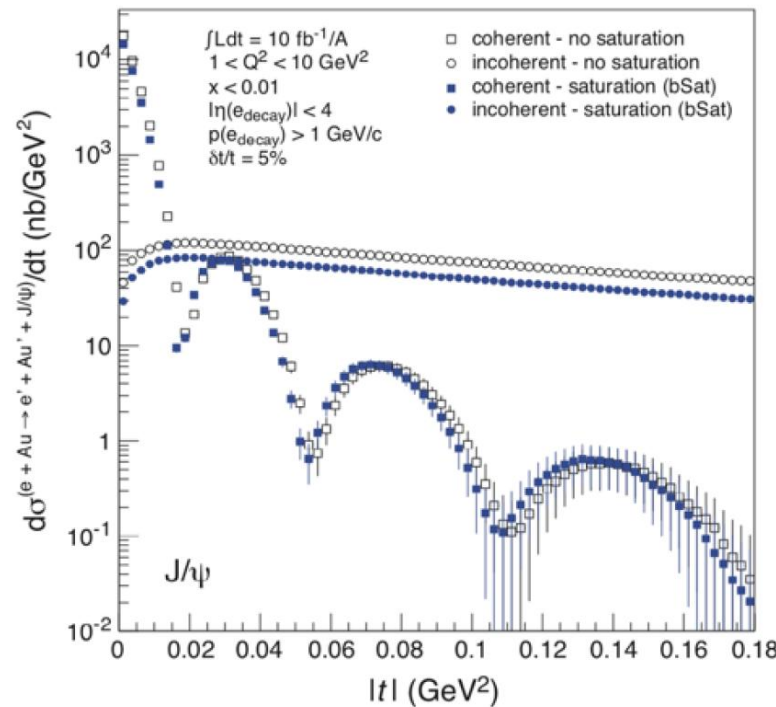
From [HFAG for CKM2025](#)

# Photons from de-excitation

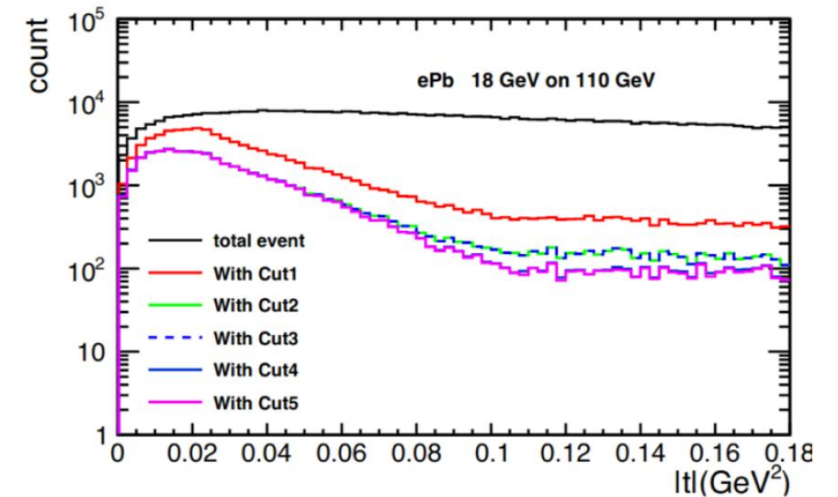
- vetoing nuclear excitation events in heavy  $eA$  collisions by tagging  $\sim 300$  MeV photons from de-excitation
- sensitivity to saturation effect through quasi-elastic  $J/\psi$  production
  - $t$ -distribution of vector meson reveals the transverse parton profile of the ion: central part may not be visible if it is saturated
  - $t$  distribution becomes narrower



reduction of the incoherent contribution through no spectator neutron (red), no additional photon (green) and no proton (purple) requirements

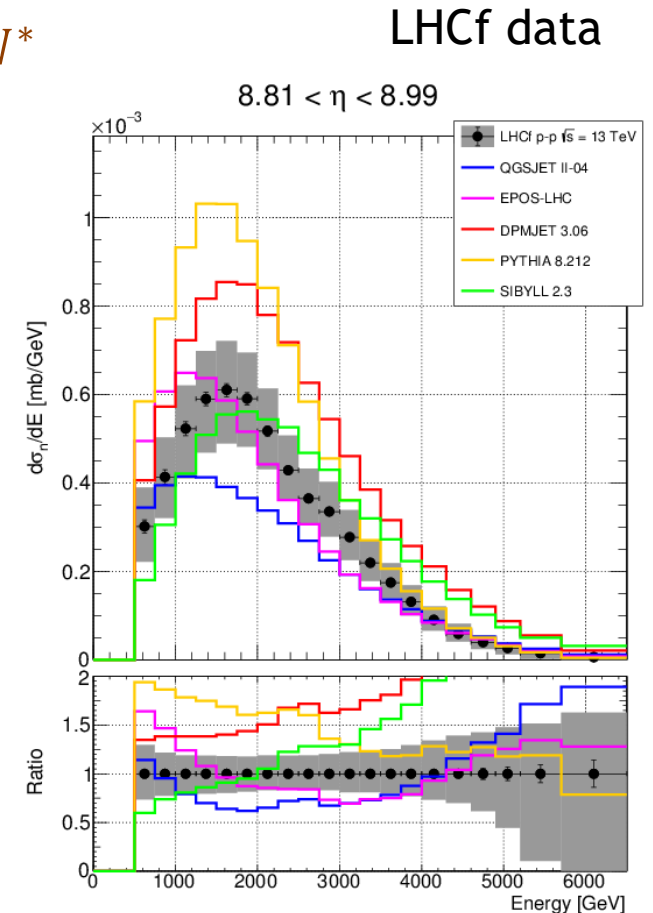
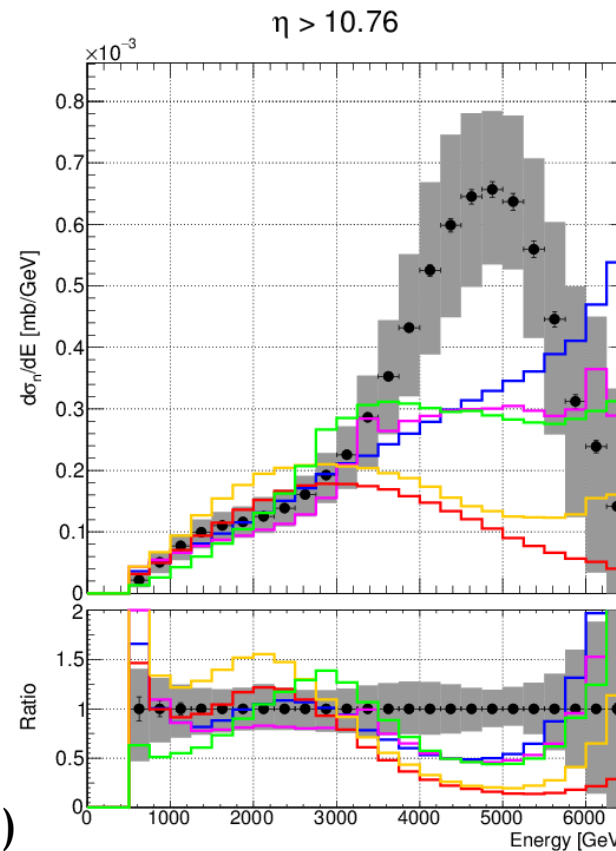
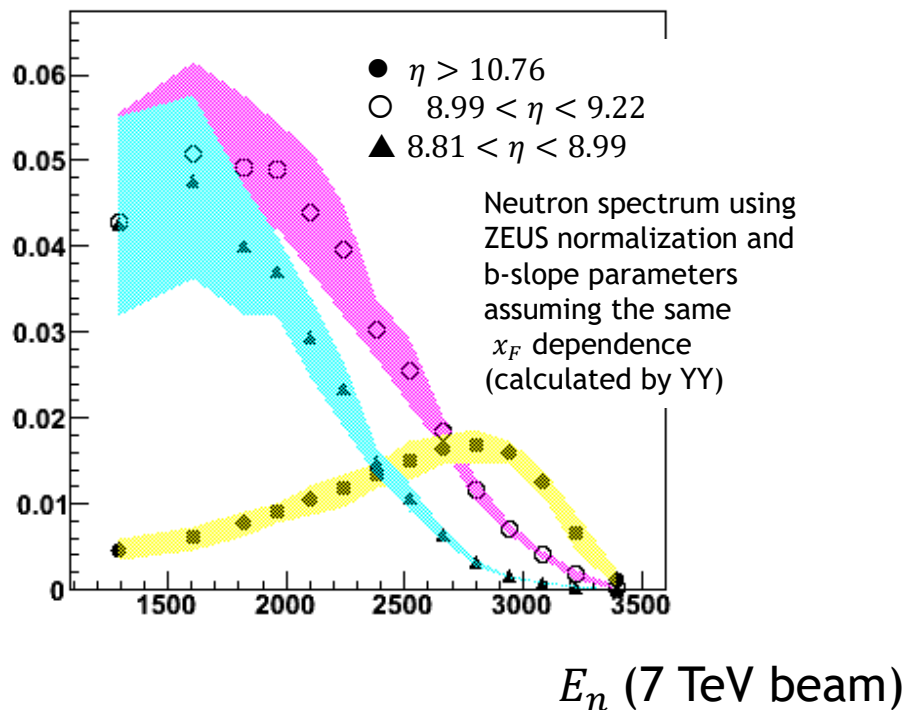


e+Au YR Fig. 7.83



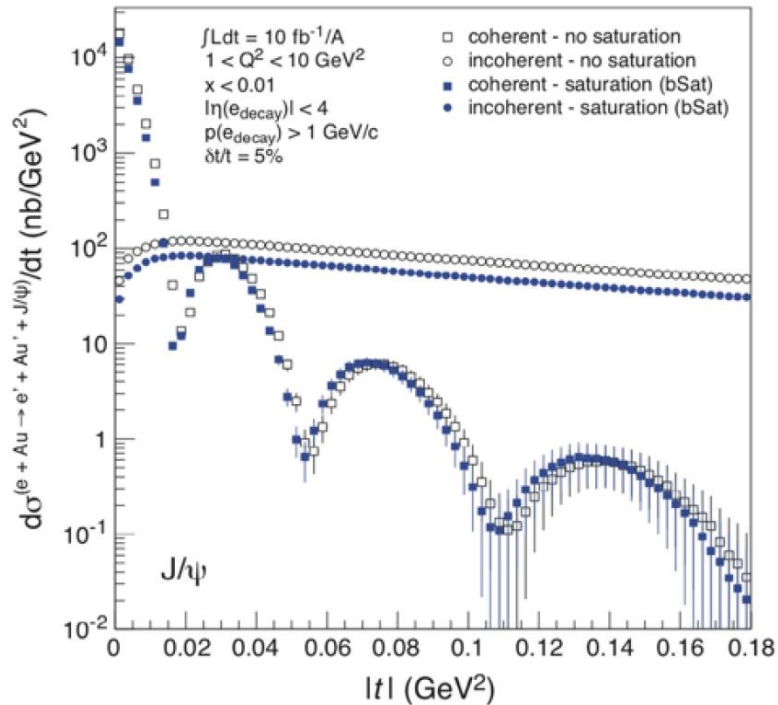
# Neutron puzzle (2): $pp$ vs $ep$

- Limited fragmentation  $\Rightarrow$  the same spectra
- LHCf data similar to  $ep$ , but models suggest harder spectrum at  $x_F \sim 1$ 
  - due to projectile fragmentation?  $pp \rightarrow N^* + Y, N^* \rightarrow n + (\text{hadrons})$
  - Corresponding to proton dissociation for  $ep$  DIS:  $\gamma^* p \rightarrow XN^*$   
LRG-tagged neutron?

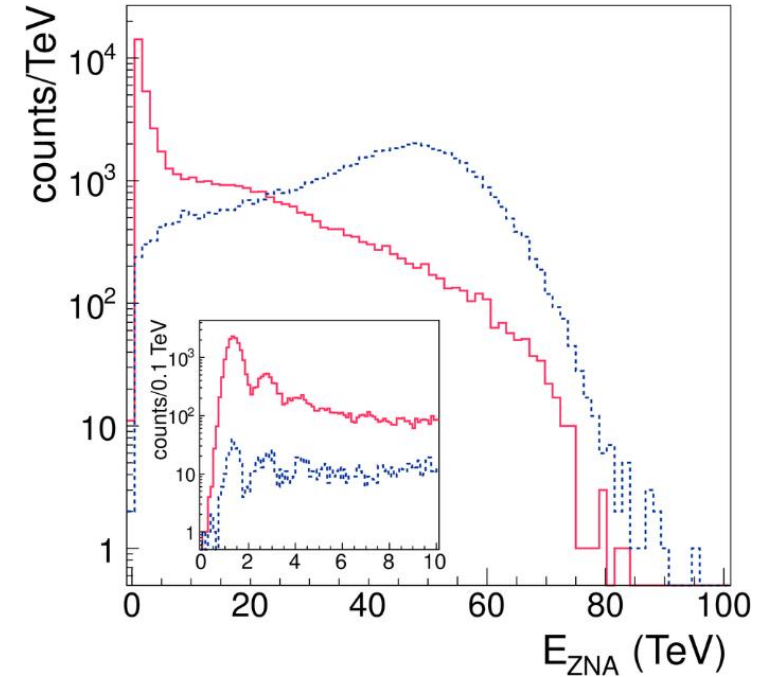
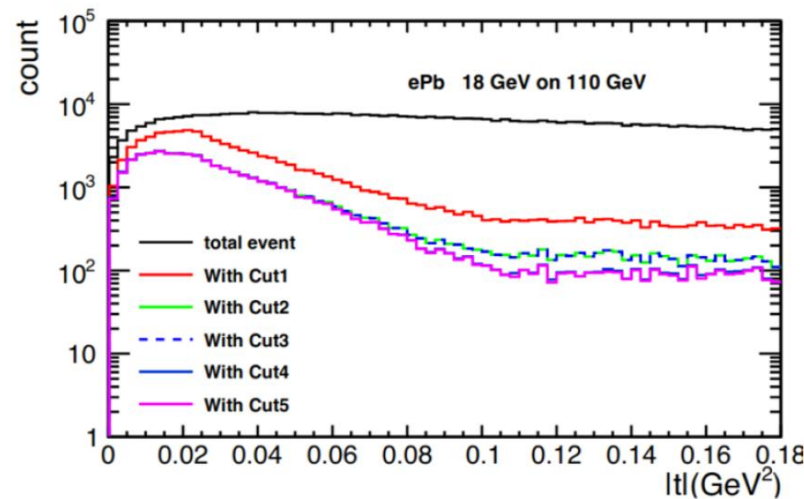


# EIC Yellow Report physics using ZDC +B0 (1)

- Spectator tagging
  - for inelastic eA collisions:  
heavy ions, eD / e<sup>3</sup>He :  $n \times E_{beam}$  energy in ZDC
  - vetoing nuclear excitation events  
by tagging ~300 MeV photons from de-excitation



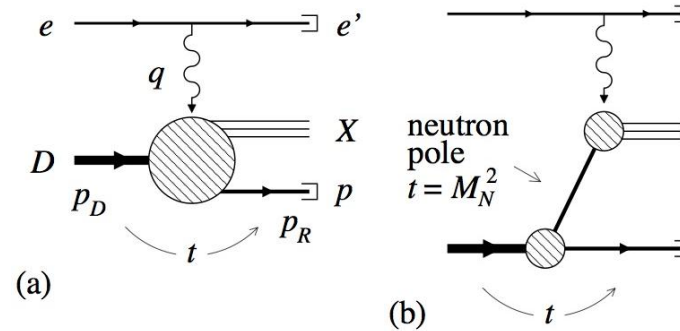
e+Au YR Fig. 7.83



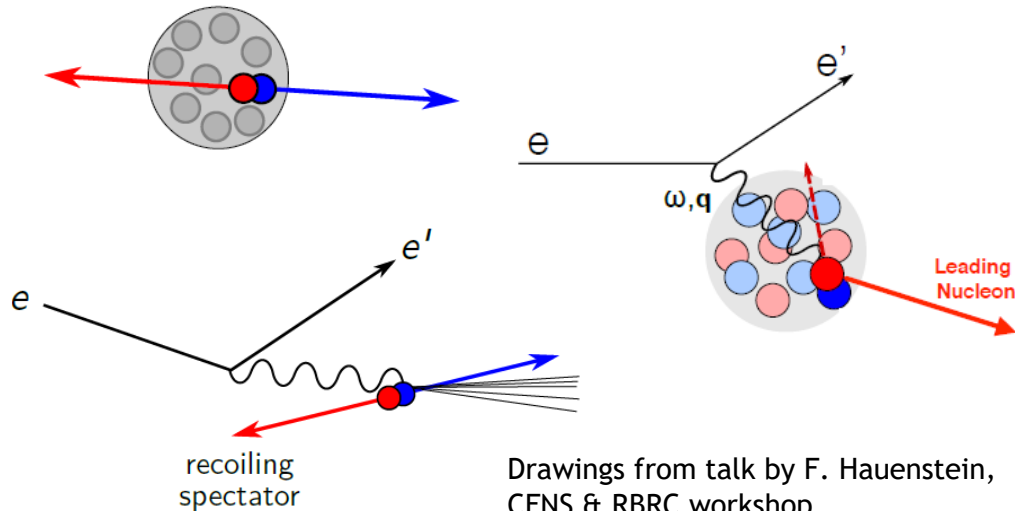
ALICE ZDC (A-side)  
with and without  
activities in plug area  
2.76 TeV run

# Proton/neutron tagged eD/eA DIS

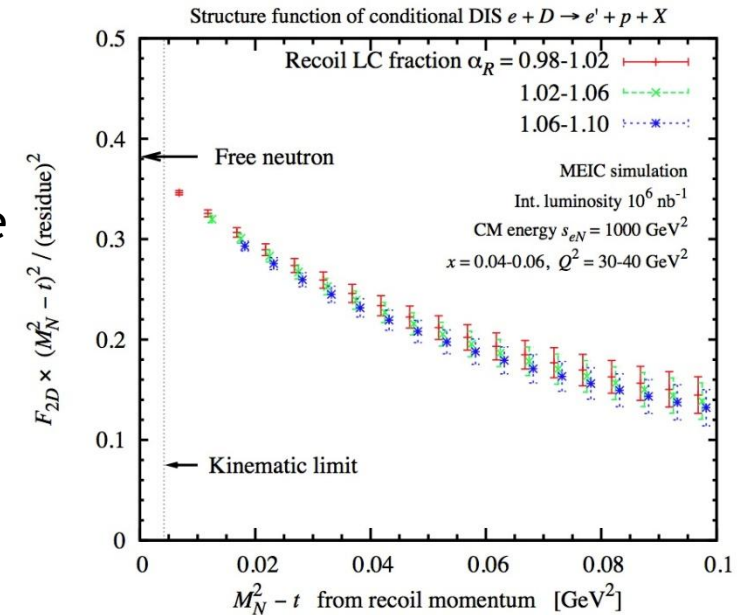
- Proton-tagged  $eD$  and  $eA$  scattering
  - $e(p + n) \rightarrow en + p$  DIS for neutron!
  - Way to understand nuclear (EMC) effect or short-range correlation (SRC) by comparing small and large system



- Neutron-tagged ( $ep + n$ ): proton structure with  $t$ 
  - Cross-check with  $en$  runs:  $t = 0$  reference given. Need to interpolate to  $t = 0$

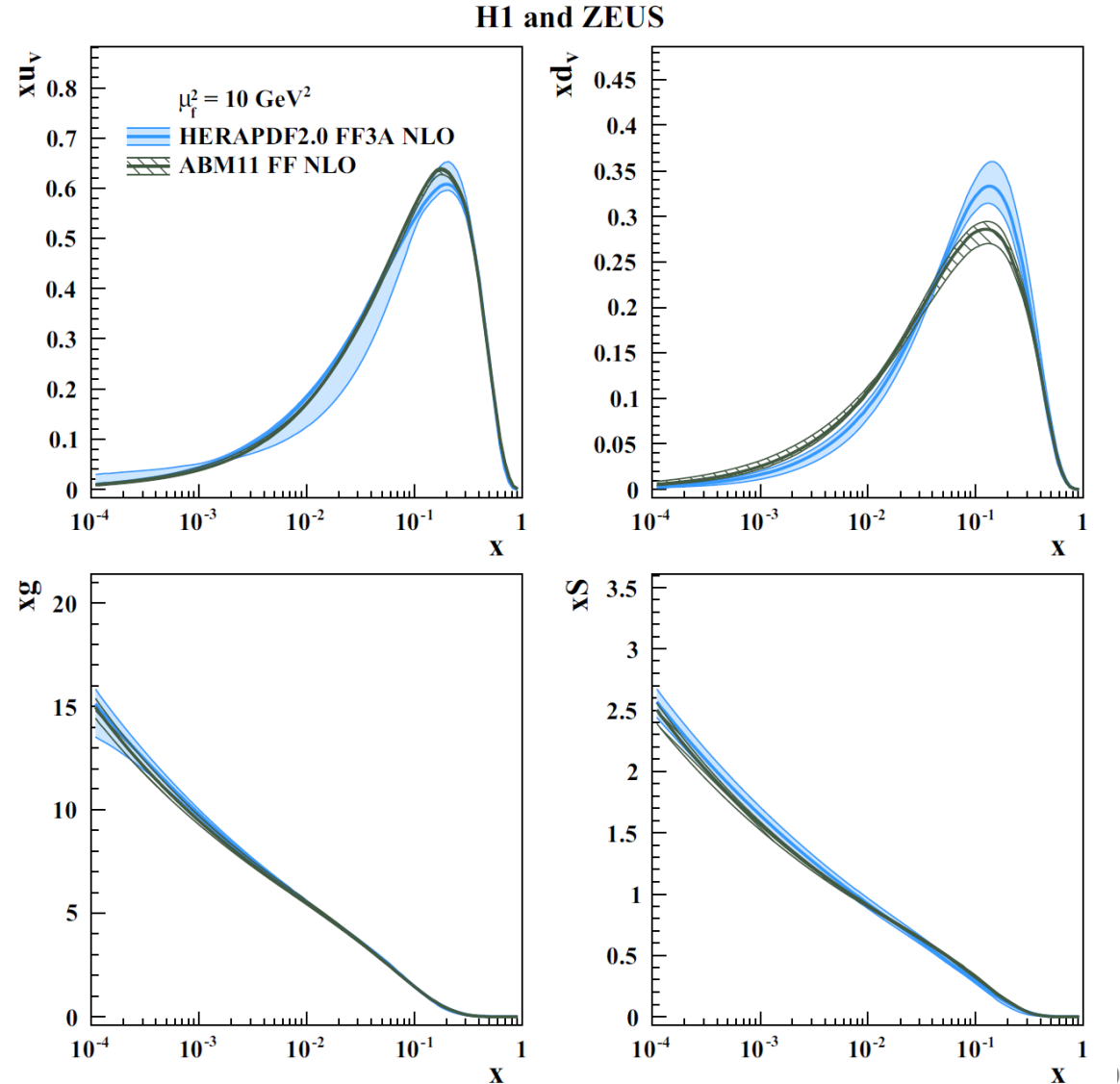


Drawings from talk by F. Hauenstein, CFNS & RBRC workshop <https://indico.bnl.gov/event/6568/>



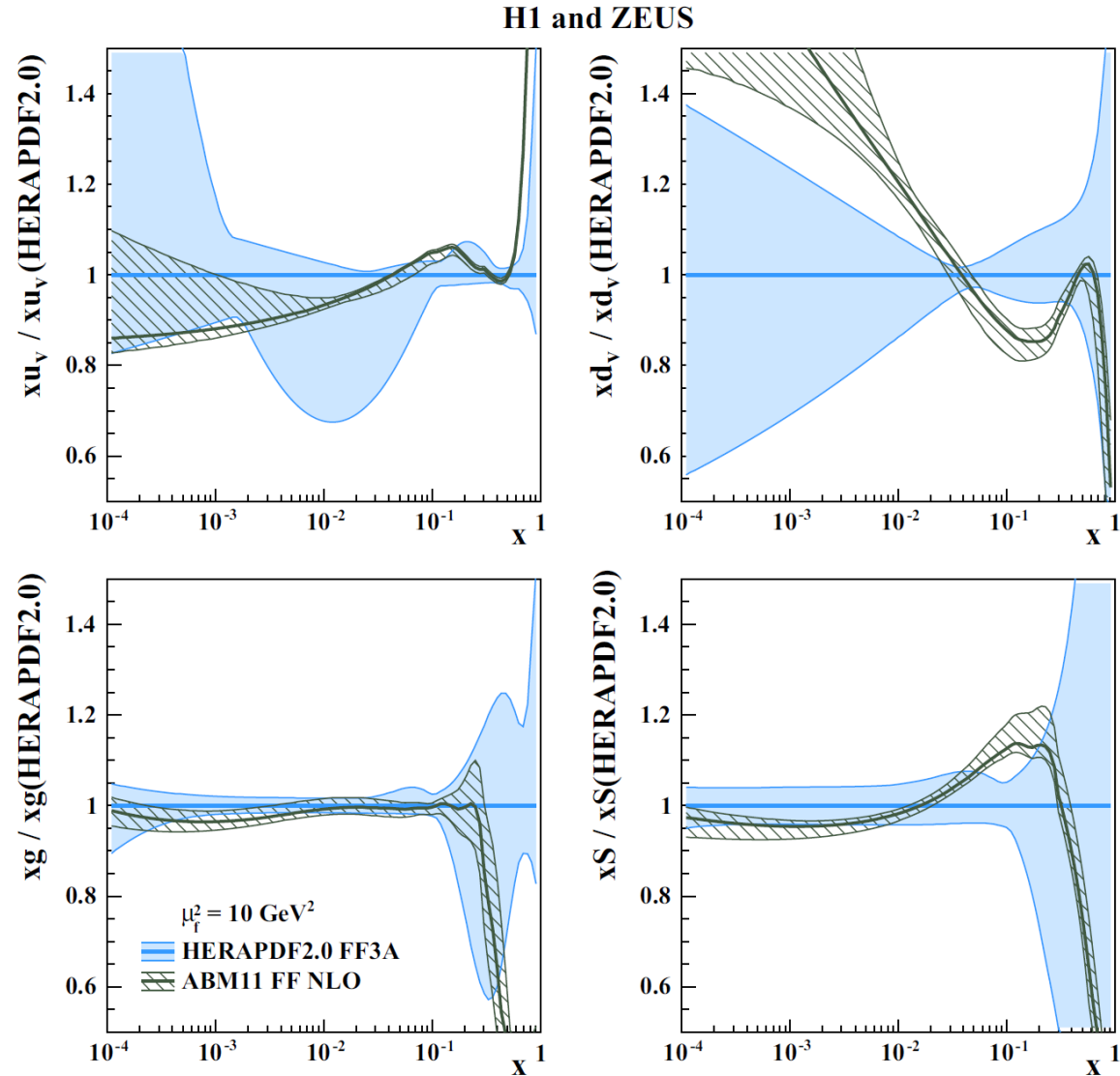
# Understanding EMC effect for HEP

- PDF fits show some mild tension between HERA and fixed-target data
- Example: HERAPDF 2.0 vs ABM11
  - HERAPDF 2.0: HERA data only
  - ABM11 (PRD 86, 054009): including
    - BCDMS, NMC, SLAC
    - Drell-Yan from FNAL
    - Dimuon from  $\nu N$
- PDF at high- $x$  will be one of the major systematics for high-mass BSM state search ( $M > 1$  TeV) at the LHC



# PDF comparison: HERA-only and with fixed target data

- This might simply be systematic of DIS experiments
  - Repeating the measurement with much better detector and environment at the EIC
- Or a real nuclear effect?
  - eD and e<sup>3</sup>He data should help



# Consideration for 300MeV photons

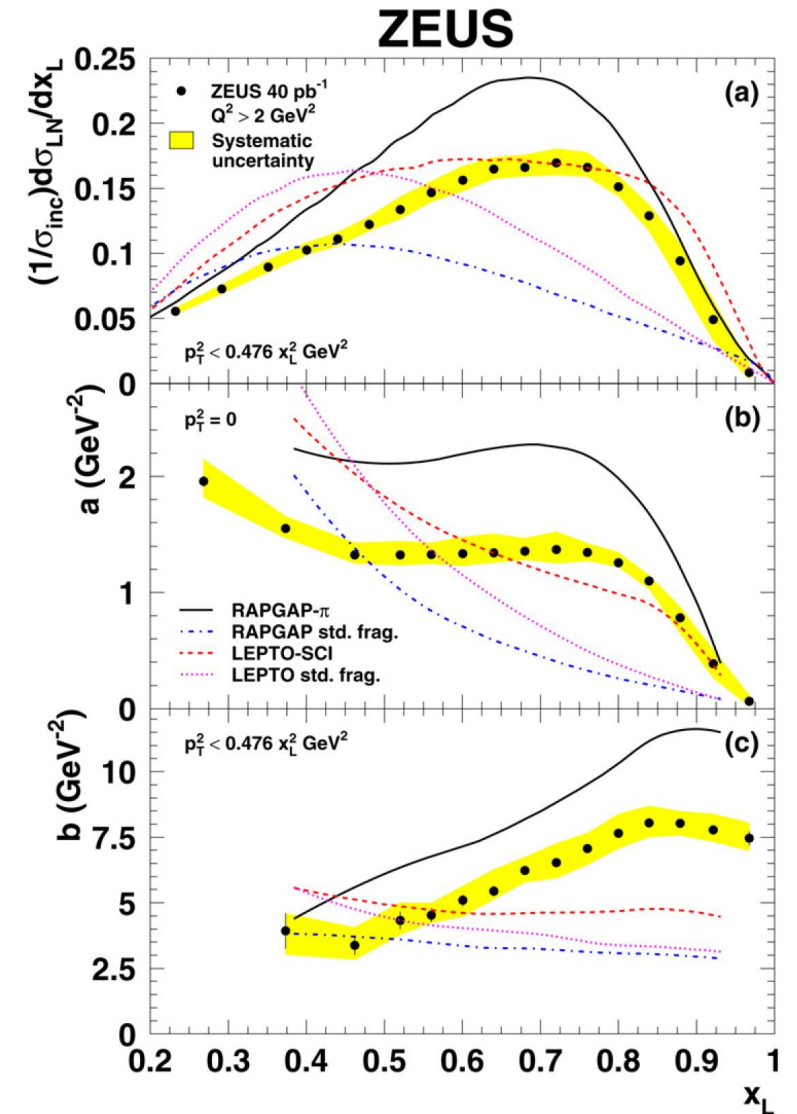
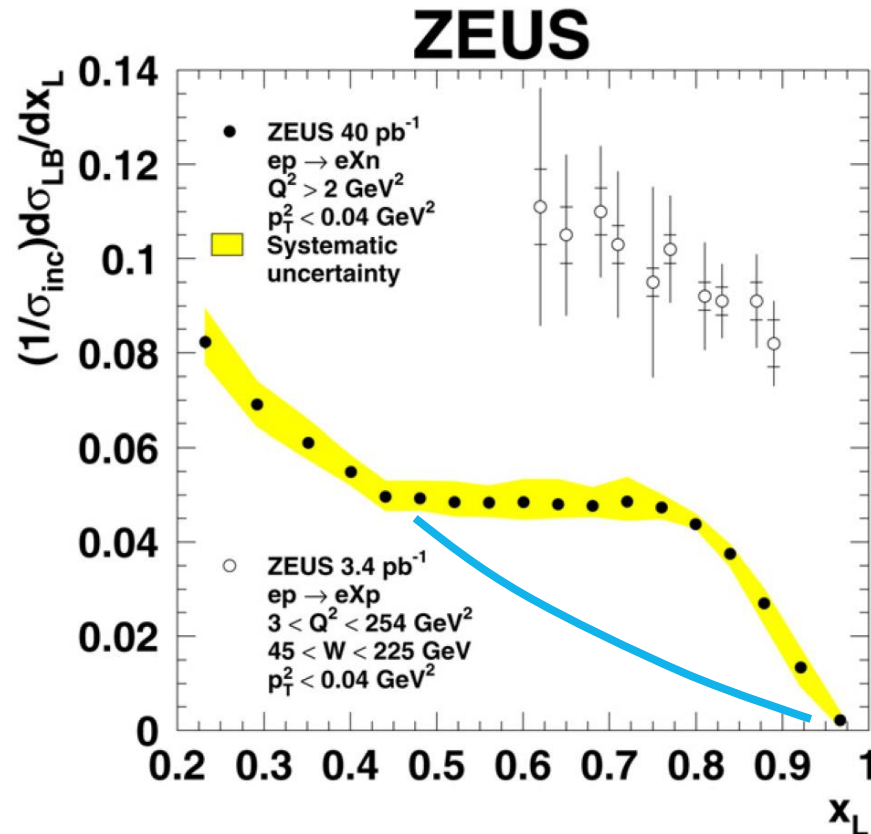
- Any initial state radiation from heavy ion?
  - maybe serious for Au, Pb etc.  
This is irreducible
- Pile-up of stray particles in the ZDC area?
  - charged pions/kaons
  - fast neutrons of  $O(1\text{GeV})$
- Need charged particle veto (tracking) in front of ZDC
- May need timing to remove most of the background, especially neutrons
  - better to use fast crystals for EM ZDC

# Vertex factorization: summary

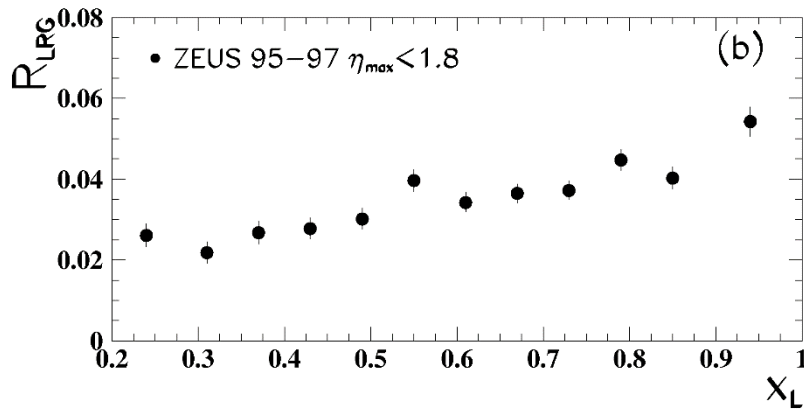
- Factorisation holds approximately
  - hint of mild absorption (20-30%)
  - very little  $(x, Q^2)$  dependence on leading baryon production probability

- Fragmentation model: one-pion (+sub-leading) peak missing

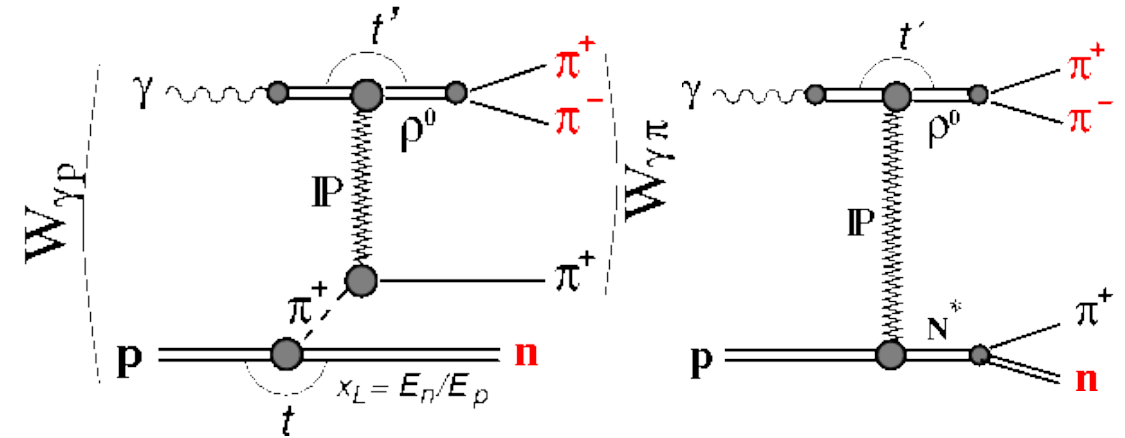
- no clear sign of meson poles in the proton spect however



# Events with a neutron at HERA and rapidity gap



← inclusive DIS (7%)

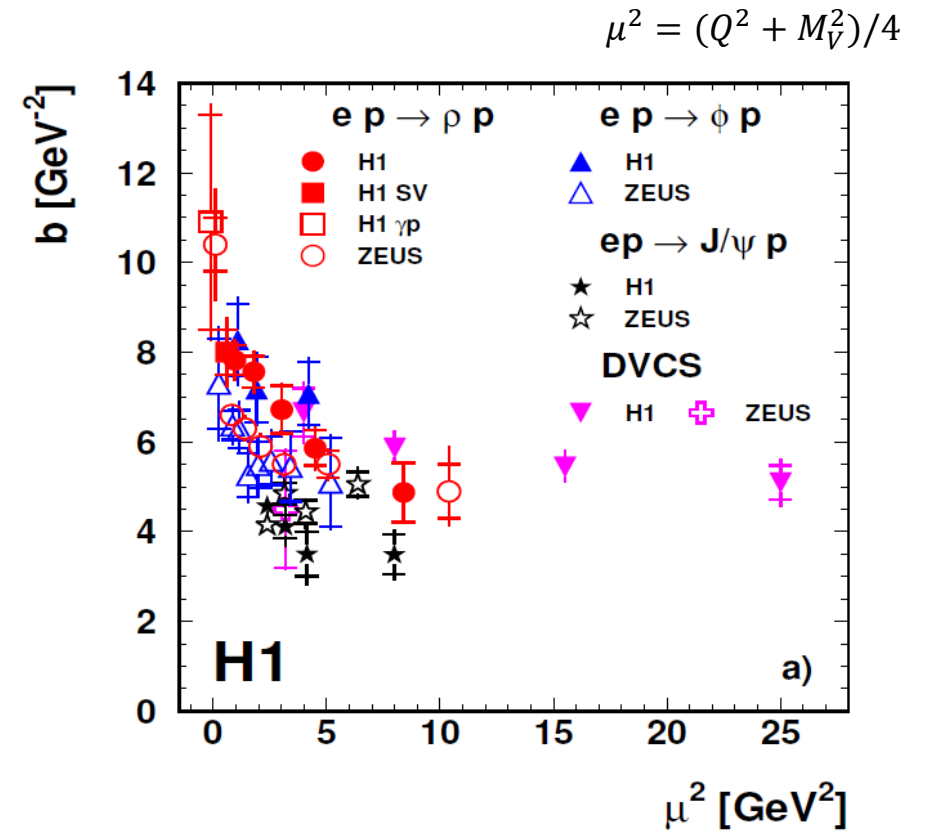
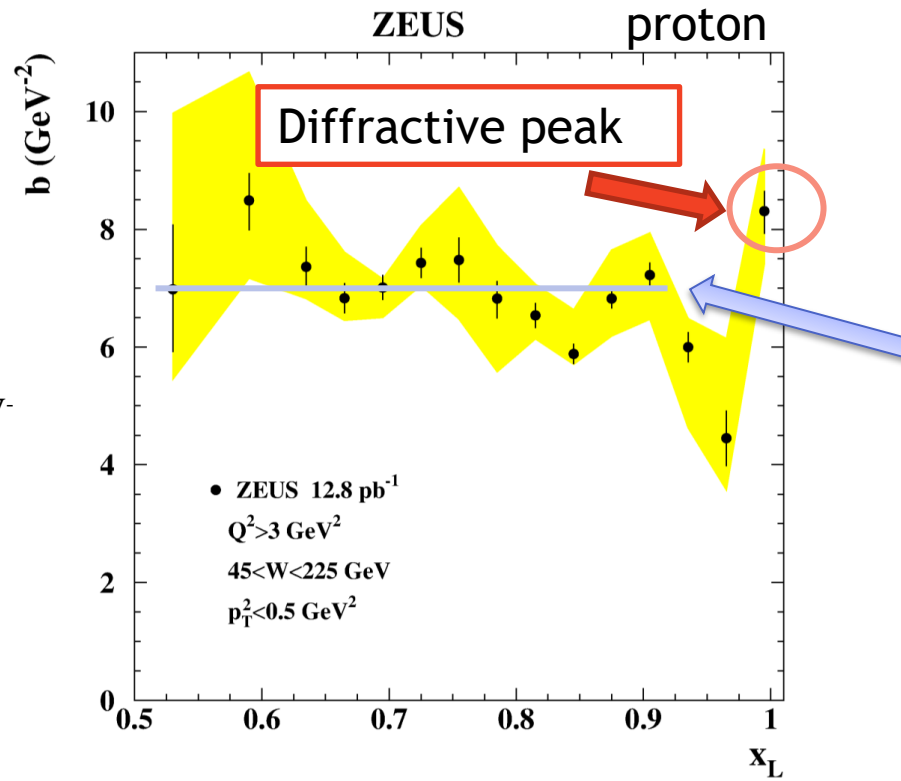
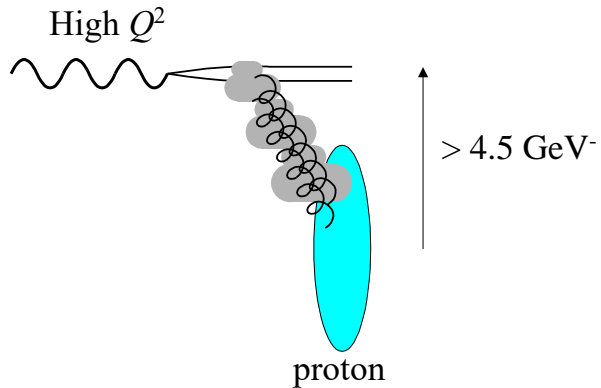


- Fraction of LRG (Large Rapidity Gap) events in neutron-tagged events: 40 - 70% of inclusive DIS
- proton elastic/diffractive: no neutron, hence decrease the above ratio
- Neutrons may come from:
  - $\gamma^* \pi$  diffractive scattering: LRG + neutron
  - proton diffractive dissociation: may yield LRG + neutron through  $N^*$ ,  $\Delta$  etc
  - fraction of LRG increases with  $z$  : possible sign of above two processes

Possible background source to OPE

# $p_T$ dependence for forward protons

- $p_T$  dependence: again almost flat for proton
  - $b \sim 7 \text{ GeV}^{-2}$  ( $\sigma \propto e^{-bp_T^2}$ ), constant
  - Slightly larger than proton size
    - Somewhat peripheral? Semi-soft, not directly probing proton



cf. quasi-elastic vector-meson production  
Strong  $Q^2$  and  $M_{VM}$  dependence

# Limiting fragmentation and "Regge" factorization

- **Limiting fragmentation:** the LN variables  $(z, t)$  are independent of  $(x, Q^2)$ 
  - large rapidity separation, (almost) no information propagated

- **Regge factorisation:** the cross section is the product of the (particle flux)  $\times \sigma(\gamma$ -particle)

$$\frac{d\sigma(W^2, Q^2, z, t)}{dzdt} = f_{\pi/p}(z, t) \sigma_{\gamma^* \pi}((1-z)W^2, Q^2)$$

- **For both models, factorisation holds** between photon vertex  $(x, Q^2)$  and baryon vertex  $(x_L, t)$ 
  - The independence does not distinguish the two models

- Proof of the Regge factorization:

$\sigma_{\gamma^* \pi}$  is independent of the soft vertex variables  $z$  and  $t$

- The factorization may be broken if a **perturbative partonic state** is exchanged between two vertices

- Precise measurement at EIC may be able to

