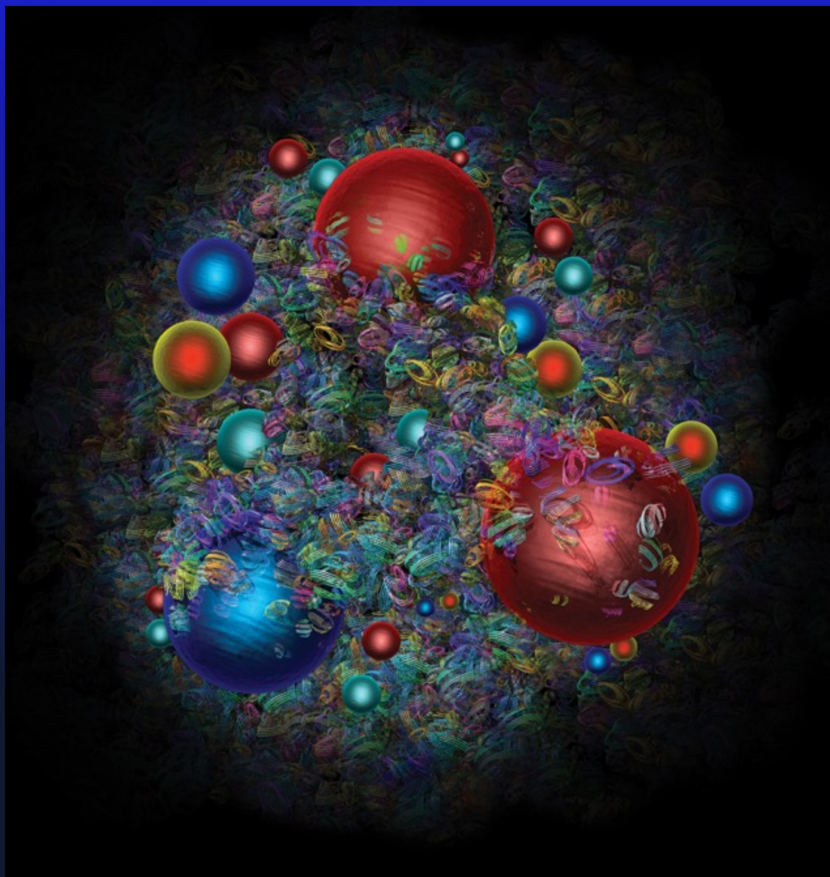
The Electron-Ion Collider: the ultimate electron microscope



Gordon Baym
University of Illinois
Urbana, Illinois

TIDC 2nd EIC Meeting
Institute of Physics, Academia Sinica
January 3, 2023

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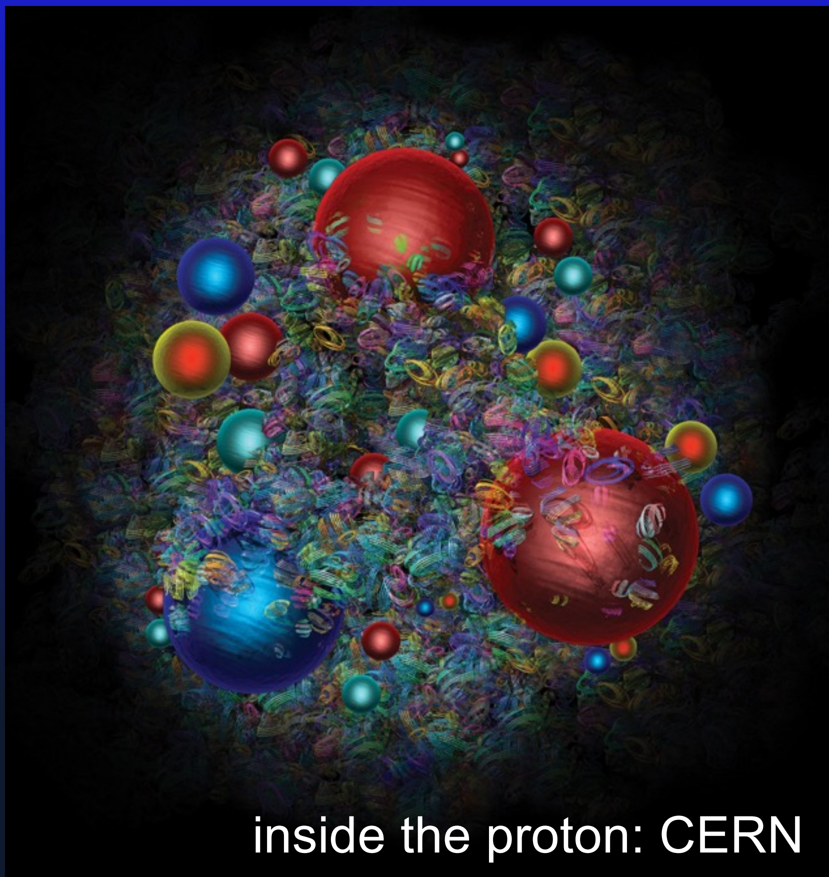

inside the proton: CERN

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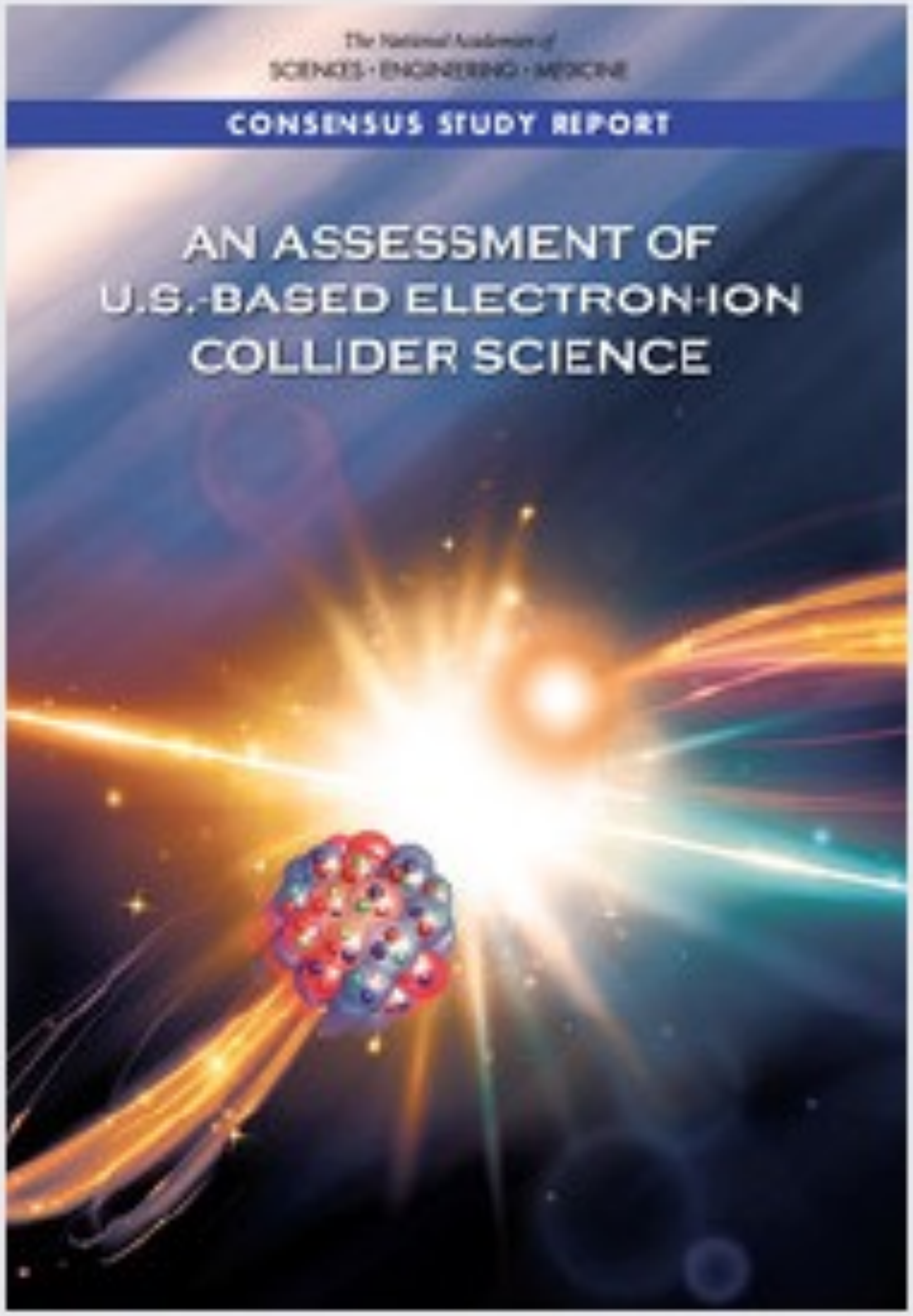
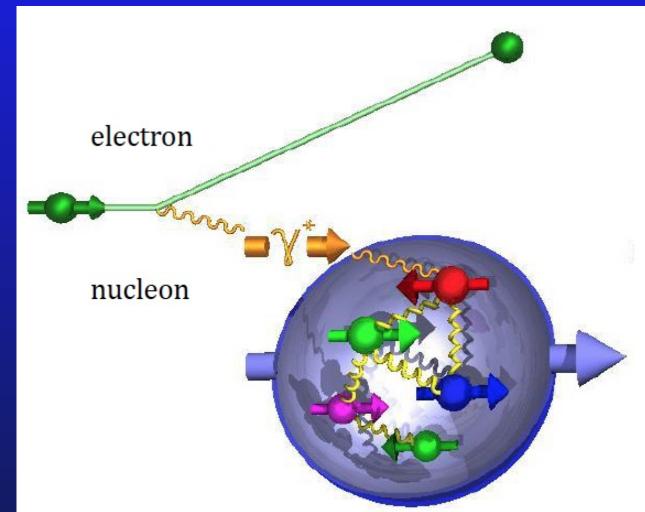


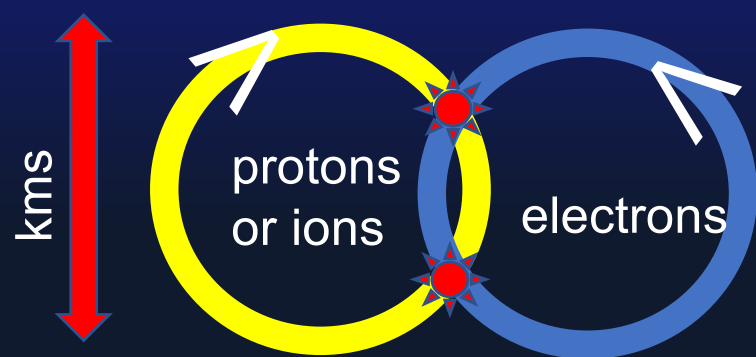
National Academy of Sciences – National Research Council report, July 2018, on the science of an EIC

Download pdf from
https://www.nap.edu/catalog/25171

The Electron-Ion Collider

A very big accelerator -- colliding beams of electrons with beams of protons or heavier ions (atomic nuclei).
A giant electron microscope for peering at the quarks and gluons deep inside the nucleon, as well as atomic nuclei. QCD machine.



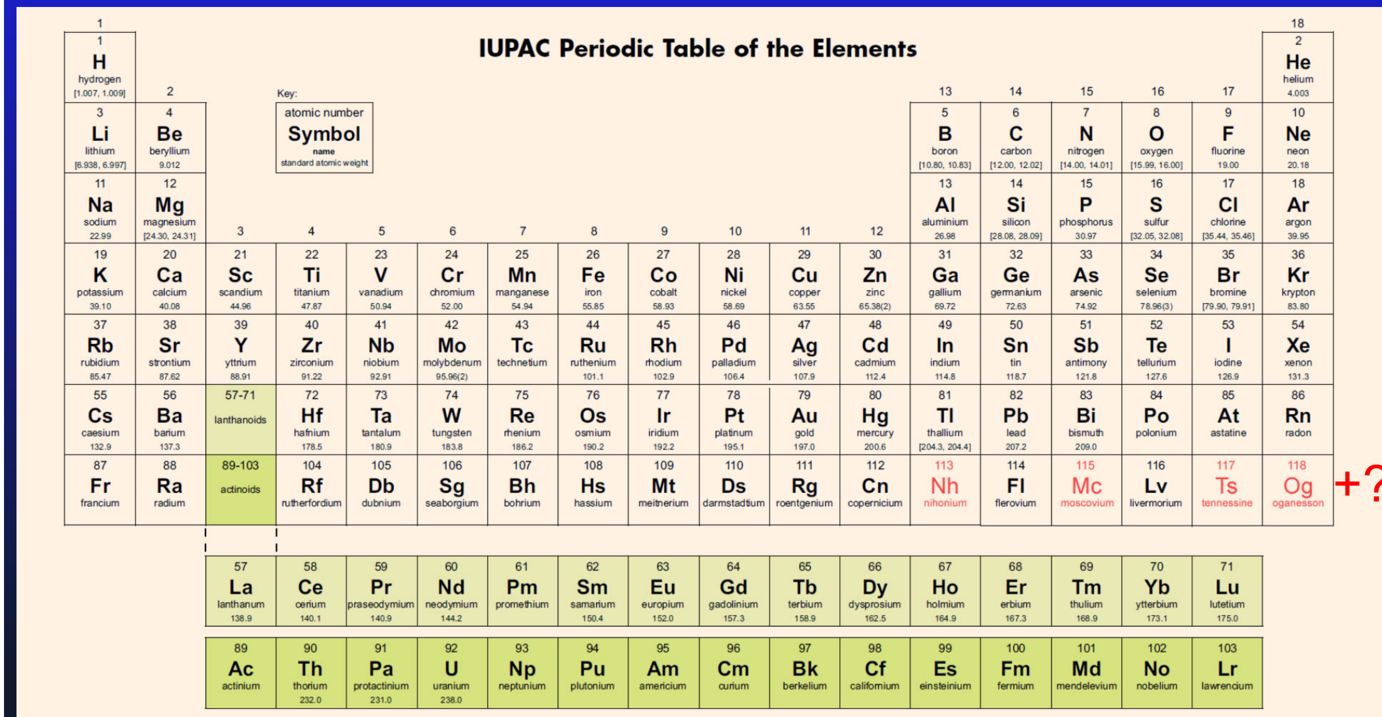


Electron microscope Invented 1931

ca. 1940

Electron-ion center of mass energy:
$\sqrt{s}$ ~ 28~140 GeV.
High luminosity (event rate) and spin polarized beams!

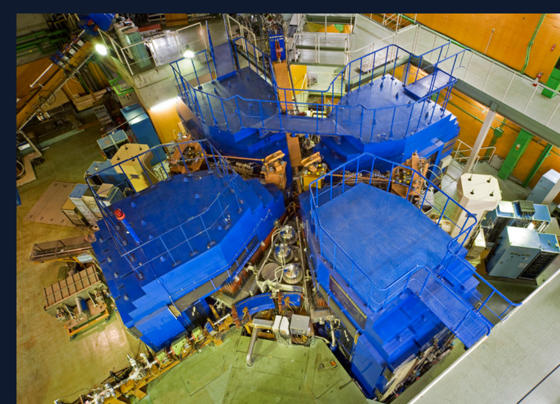
Atomic nuclei: building blocks of the everyday world:



The periodic table of chemical elements is over 150 years old.

Are there further elements out there? Any of them stable?

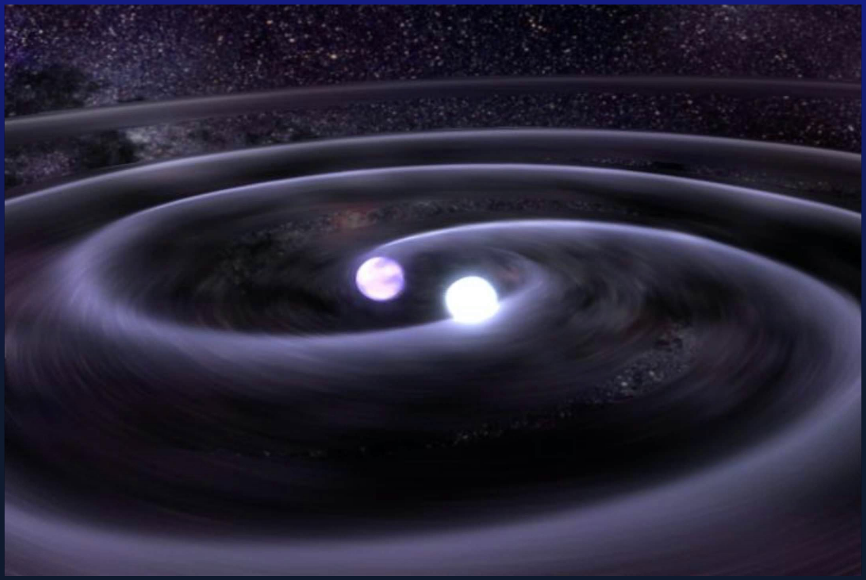
How many isotopes does each element have?
Answers from rare isotope accelerators (e.g., RIBF at RIKEN in Japan and FRIB at MSU) studying nuclei far from stability.

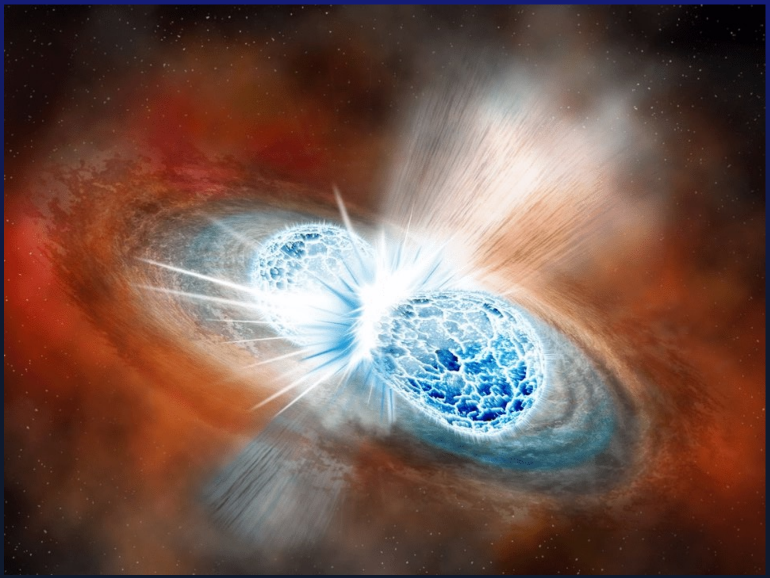


RIKEN

F = facility, RIB = rare isotope beam

How to make the heavier elements? Answers, remarkably, from multimessenger studies of binary neutron star mergers:
Merger GW170817 observed on 17 Aug. 2017 by LIGO and Virgo (gravitational radiation), FERMI (gamma ray telescope) + some 70 other electromagnetic observatories.





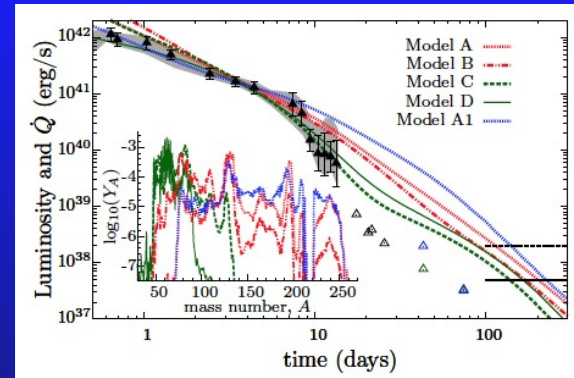
Two neutron stars merging, emitting gravitational radiation and post-merger formiing:

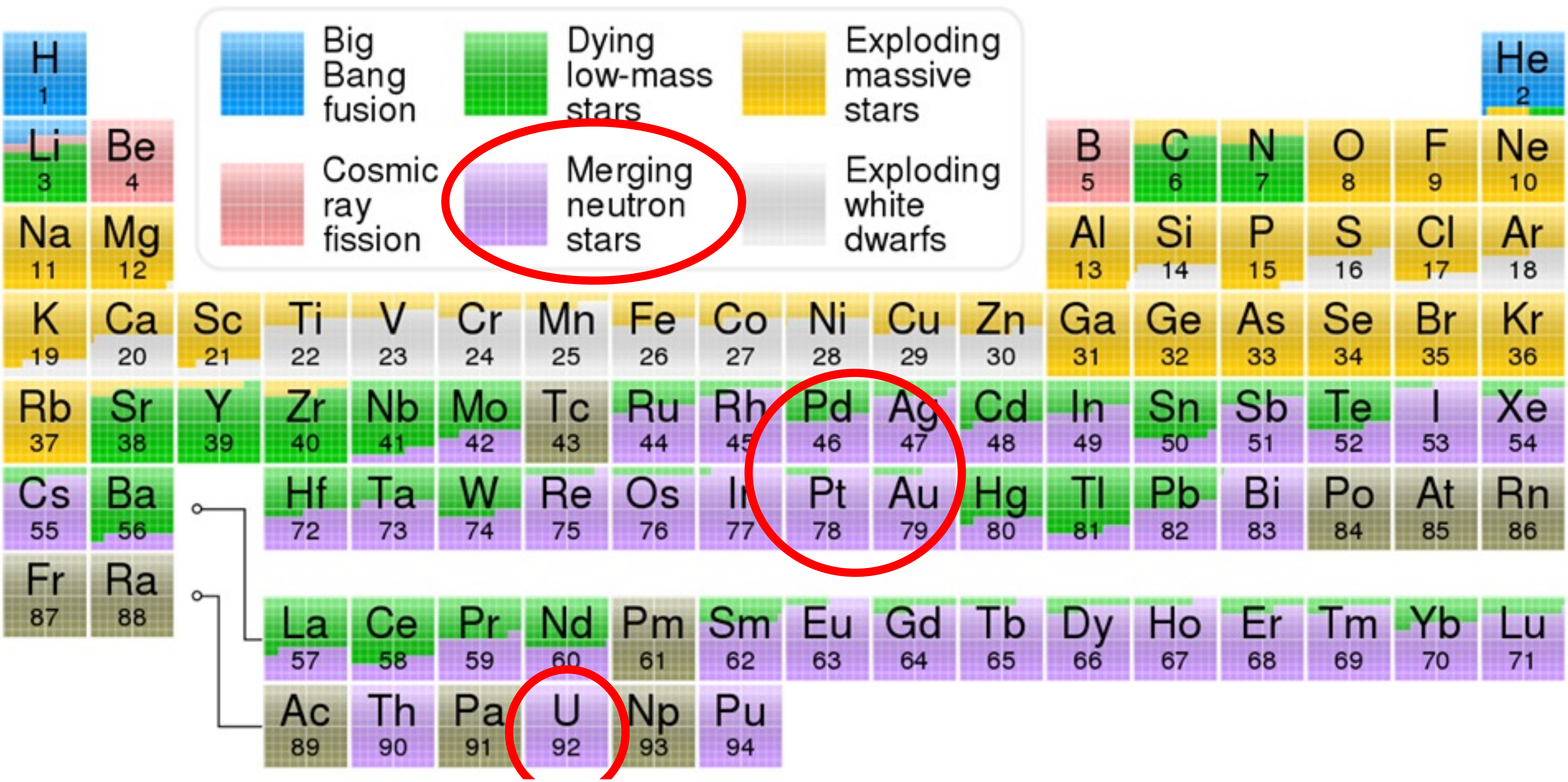
Kilonova: neutron-rich site of r-process

Binary neutron star mergers likely site of heavy element production

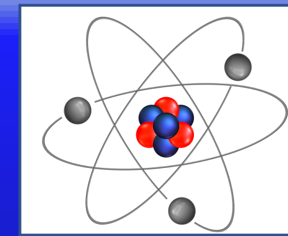
r-process in *kilonova* =>
earth-scale masses of
Au, Ag, Pt, U... and Sr

Kilonova
N= 50, 82, 126





Atoms are made of electrons and nuclei.
Similarly neutrons, protons, and nuclei are made of quarks and gluons. But how? How to explain nucleon masses, spin, magnetic moments, etc. in terms of quarks and gluons?



Quarks = fractionally charged spin-1/2 fermions, baryon no. = 1/3, with internal SU(3) color degree of freedom.

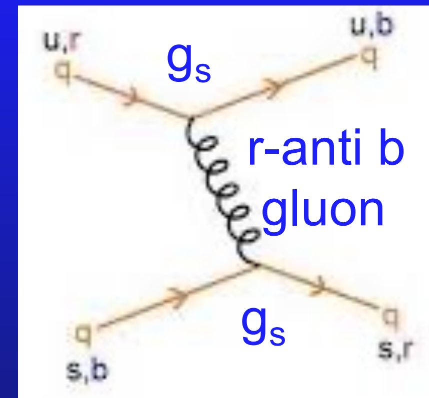
Flavor	Charge/|e|	Mass(MeV)
u	2/3	~2
d	-1/3	~5
s	-1/3	~ 94
c	2/3	~1280
b	-1/3	~4200
t	2/3	~175,000

proton = u + u + d
neutron = u + d + d
$\pi^+ = u + \bar{d}$, etc.

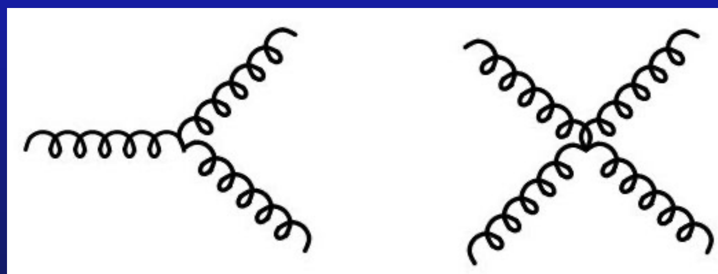
Form of baryons in the early universe at t < 1 microsec (T > 100 MeV) and in the deep interiors of heavier neutron stars.

Strong interactions – quantum chromodynamics

Quarks interact by exchanging gluons – massless vector bosons (like photon) with spin 1, and coming in 8 colors. Gluons also interact with each other!!



QCD

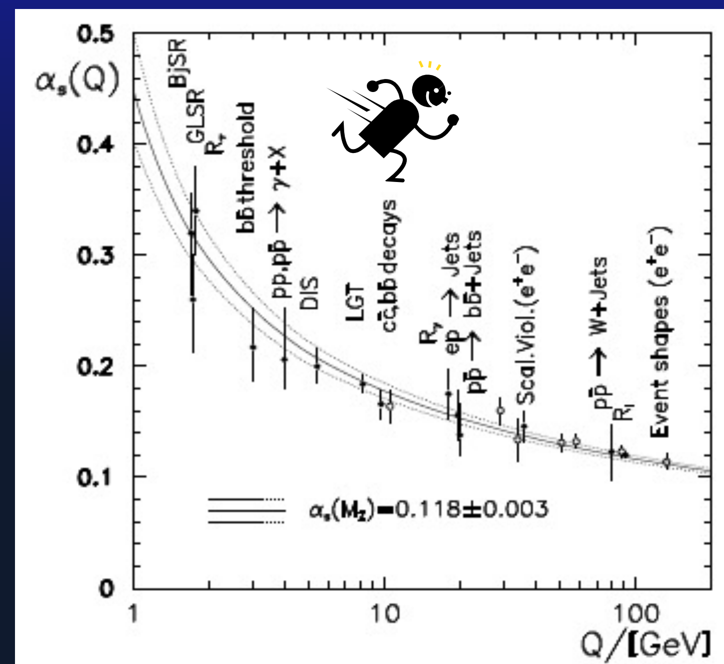


$$\alpha_s(\mu) = \frac{g_s^2}{4\pi} = \frac{6\pi}{(33 - 2N_f)\ln(\mu/\Lambda_{QCD})}$$

μ = energy scale Λ_{QCD} ~ 340 MeV

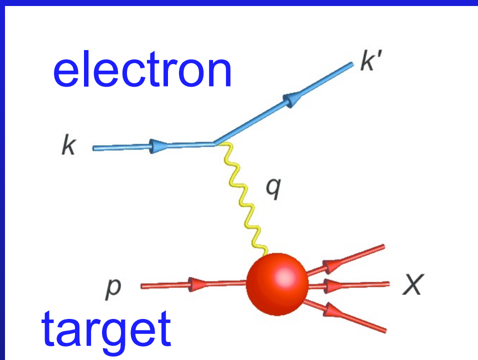
Asymptotic freedom as $\mu \to \infty$

(Even at Grand Unified (GUT) scale, 10^{15} GeV, g_s ~ ½; cf. electrodynamics: $e^2/4\pi = 1/137 \Rightarrow e \sim 1/3$)

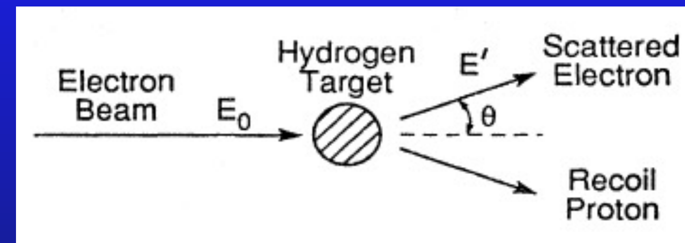


Running coupling constant

Electron scattering on nucleons and nuclei: origins



Electrons (without internal structure) are precise probe of the complex structure of nucleons and nuclei.



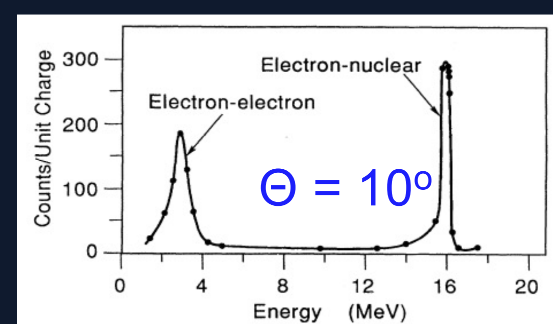
First scattering of electrons on nuclei, Illinois Betatron 1951

Al Hanson,

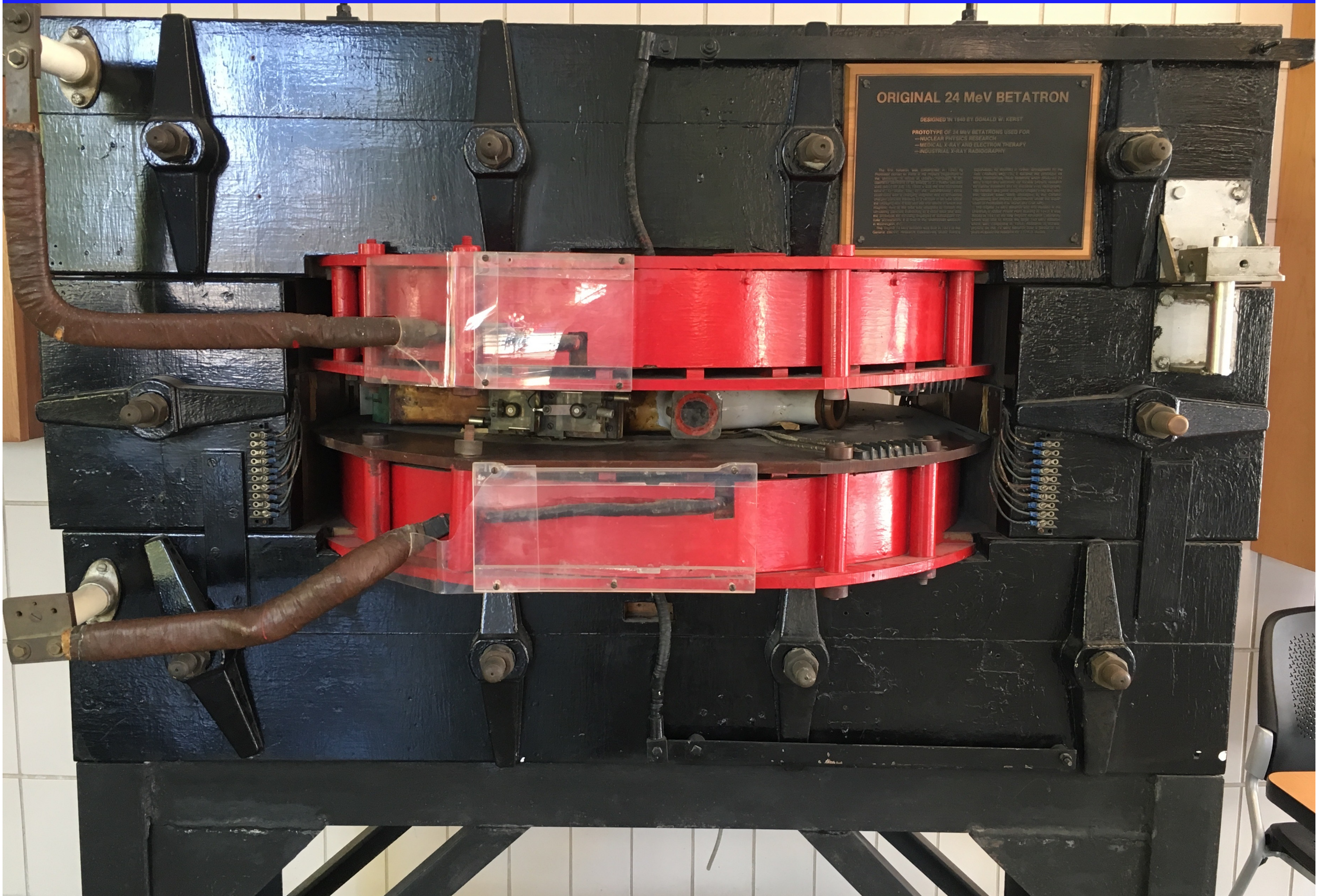
Ernie Lyman

and Merrill B. Scott

15.7 MeV electron beam on Be and Au foils



Donald Kerst, U of I, with the Betatron, ca. 1941


ORIGINAL 24 MeV BETATRON
DESIGNED IN 1940 BY DONALD W. KERST
PROTOTYPE OF 24 MeV BETATRONS USED FOR
—NUCLEAR PHYSICS RESEARCH
—MEDICAL X-RAY AND ELECTRON THERAPY
—INDUSTRIAL X-RAY RADIOGRAPHY

First scattering of electrons on protons and alpha particles

R. W. McAllister and Robert Hofstadter, 1956, with 187 MeV electrons at Stanford High Energy Physics Lab (HEPL) on H_2 and He gaseous targets, and then on polyacetylene (CH_2).

Found first hint of internal structure of the proton from the angular distribution:

proton radius $r_{proton} \sim 0.7$ fm (1fm = 10^{-13}cm)

But is the inside of the proton a continuous "pudding" or are its constituents point particles?

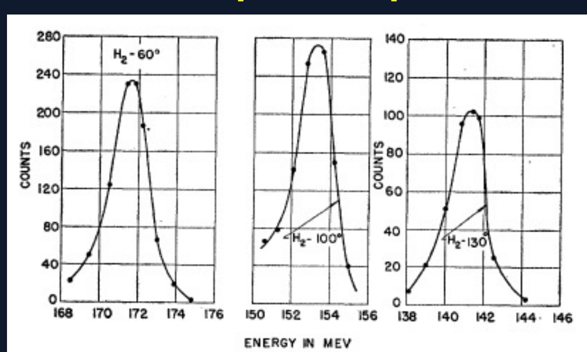
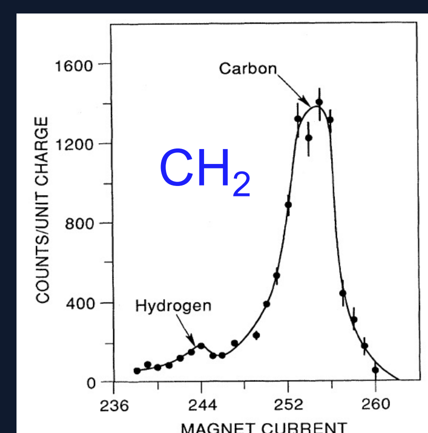
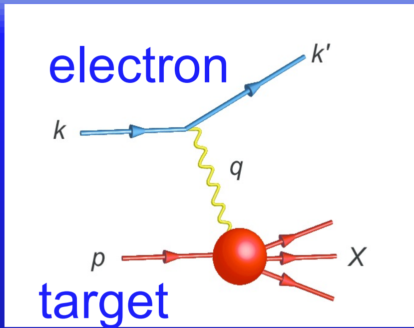


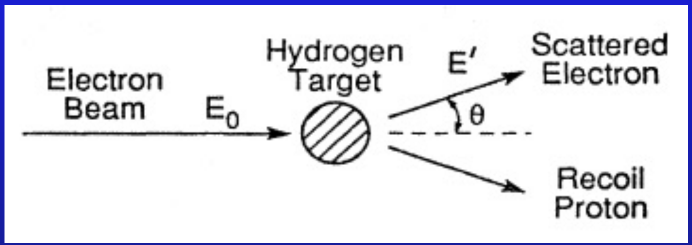
FIG. 3. Typical elastic profiles obtained with hydrogen gas at 185 Mev.

Elastic scattering on H_2 at various Θ





Kinematic variables in electron scattering from nucleons and nuclei



Electron scattering angle: θ

Electron energy transfer in lab: $\nu = E_e - E_e'$

3-momentum transfer from electron: $\mathbf{q}$

4-momentum transfer squared: $Q^2 = \mathbf{q}^2 - \nu^2$
Larger Q^2 = higher (transverse) resolution

Elastic scattering on target of mass m:
Energy conservation $(m^2 + q^2)^{1/2} = m + \nu$ => $\nu = Q^2/2m$

Bjorken scaling variable: $x = Q^2/(2m_{proton}\,\nu)$
"shutter speed" ~ 1/x

The variables Q^2 and x define the landscape of electron scattering

Discovery of quarks 1967-69

Friedman, Kendall, and Taylor do deep-inelastic scattering (DIS) destroying proton target – at the SLAC (Stanford Linear Accelerator Center) 2 mile long, 20 GeV electron accelerator.

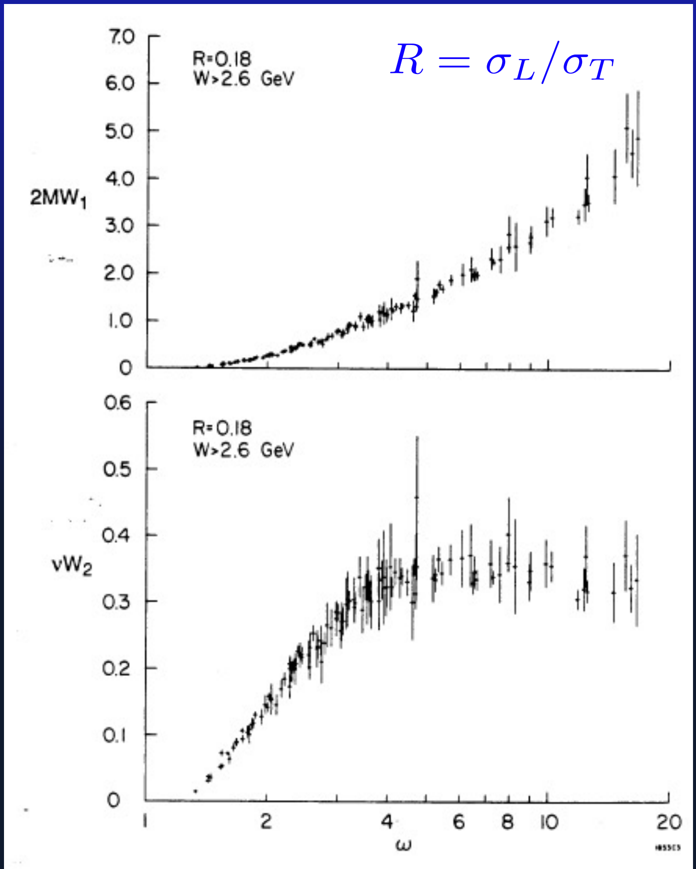
Measure electron angular cross section:

$$\sigma(E,E',\theta)=\frac{4e^4E'^2}{Q^4}\left\{W_2(\nu,Q^2)\cos^2\frac{\theta}{2}+2W_1(\nu,Q^2)\sin^2\frac{\theta}{2}\right\}$$

For scattering from point particles inside proton, W_1 and νW_2 depend only on Bjorken scaling variable

$$x=\frac{Q^2}{2m_{proton}\nu}$$

Observe dependence on x only; shows that proton is made of point particles.



Bjorken scaling indicating quark structure. ω = 1/x
(W = mass of recoiling target)

Proposal of quarks as mathematical model 1964

Volume 8, number 3 PHYSICS LETTERS 1 February 1964

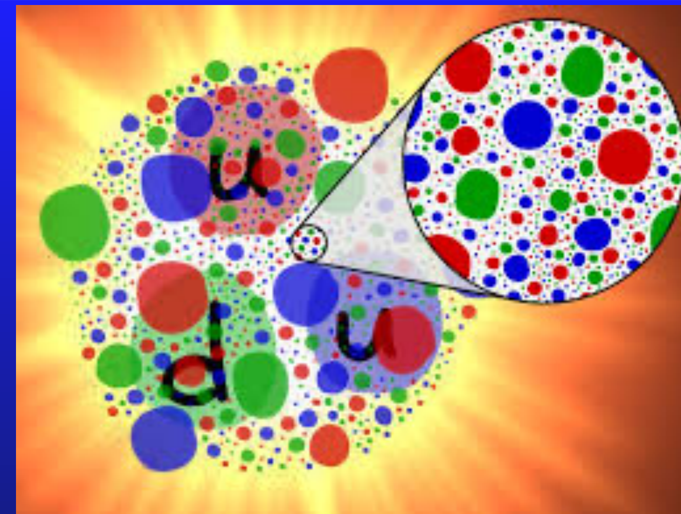
A SCHEMATIC MODEL OF BARYONS AND MESONS *

M. GELL-MANN
California Institute of Technology, Pasadena, California

Ordinary matter near the earth's surface would be contaminated by stable quarks as a result of high energy cosmic ray events throughout the earth's history, but the contamination is estimated to be so small that it would never have been detected. A search for stable quarks of charge $-\frac{1}{3}$ or $+\frac{2}{3}$ and/or stable di-quarks of charge $-\frac{2}{3}$ or $+\frac{1}{3}$ or $+\frac{4}{3}$ at the highest energy accelerators would help to reassure us of the non-existence of real quarks.

Parton model (1969)

R.P. Feynman and J.D. Bjorken:



Understand electron scattering in terms of *partons* -- quasiparticles

Given parton carries momentum p (in beam direction) -- a fraction x of the total target proton momentum p_{proton} .   Same x

$$\frac{p_{parton}}{p_{proton}} = x = \frac{Q^2}{2m_{proton}\nu}$$

energy conservation in scattering on parton $\nu = \frac{Q^2}{2m_{parton}}$

$$m_{parton} = x\, m_{proton}$$

Measure x in "infinite momentum" frame, i.e., with the proton moving at (nearly) the speed of light.

Point partons identified with quarks, antiquark, and gluons -- all governed by QCD, with asymptotic freedom. Brodsky and Farrar

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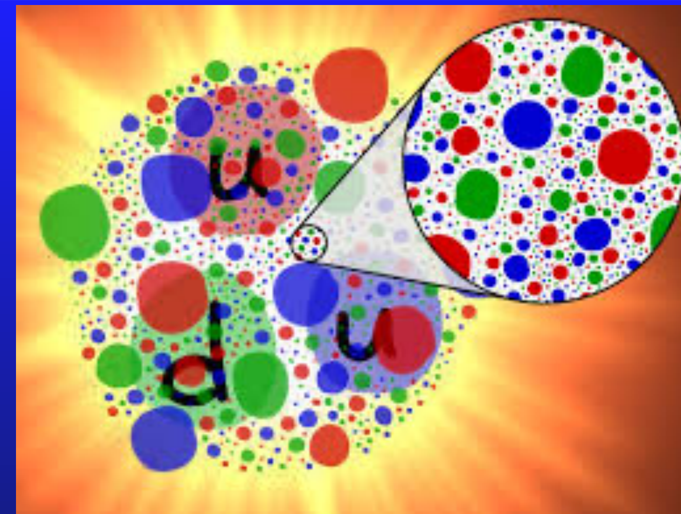

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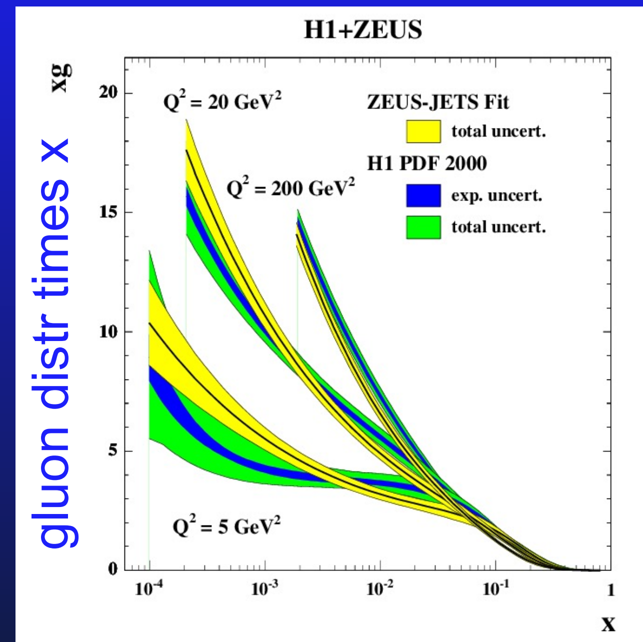
1991-2007 e scattering at HERA
(Hadron-Electron Ring Accelerator, Hamburg)

High-energy collisions of 27.5 GeV electron and positron beams (polarizable) with 920 GeV proton beams (unpolarized).

HERMES (spin) fixed-target experiment.

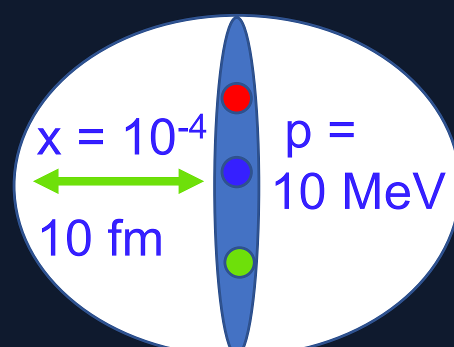
HERA => great abundance of low x gluons within nucleon.



How can electrons probe electrically neutral gluons?

$g \rightarrow q + \bar{q}$ causing an effective electric dipole moment

View proton in frame in which proton is “slower” than ∞ momentum. Heisenberg => low x “wee” partons stick out.

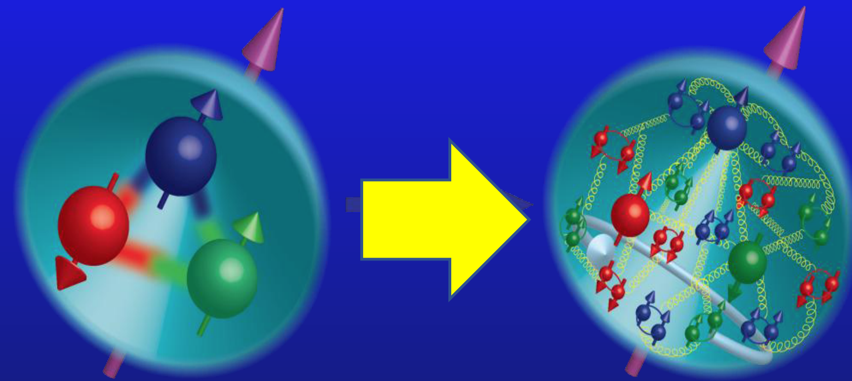


Lorentz contracted

100 GeV proton

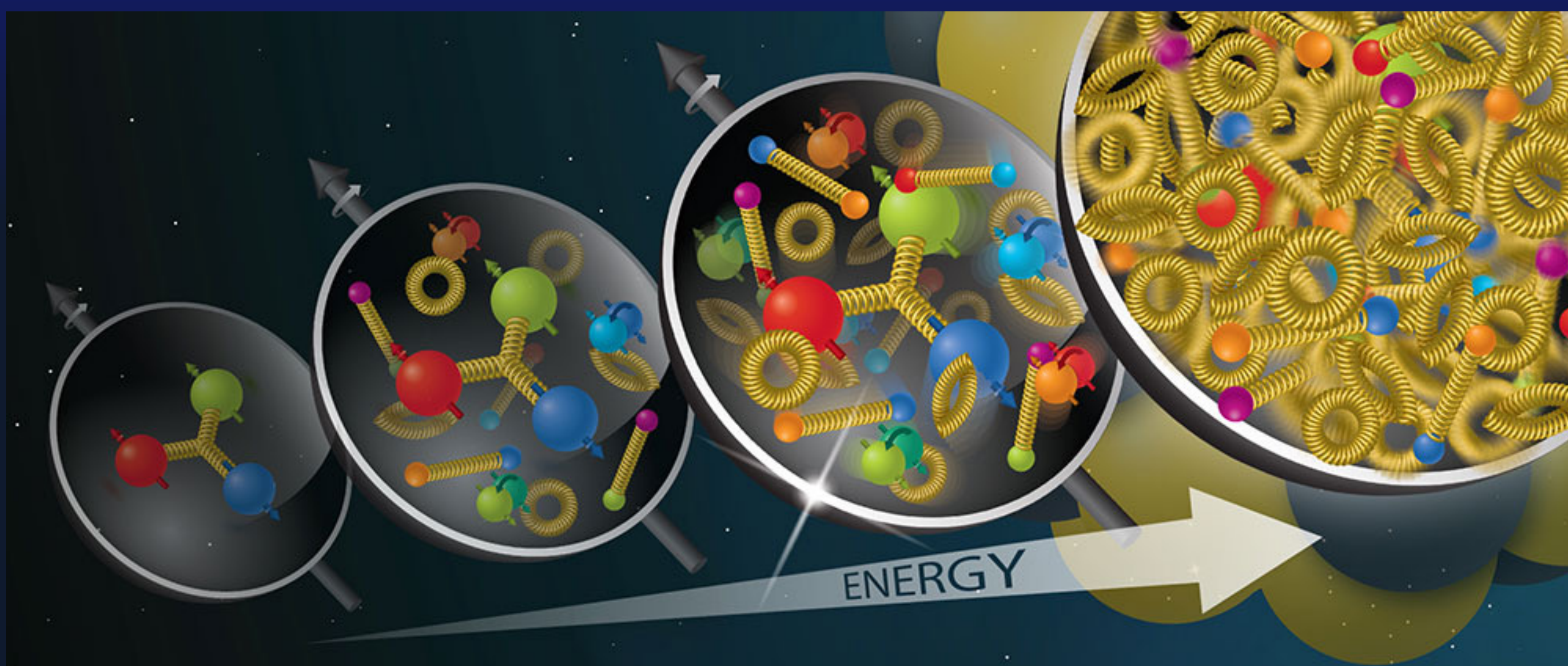
Physics goals of an EIC

Our picture of the nucleon has evolved considerably from the first simple picture as a "bag" of three valence quarks.
Have sea of quark-antiquark pairs (u, d, s) as well as cloud of gluons, all interacting!



1970s cartoon — Now

Z-E Meziani



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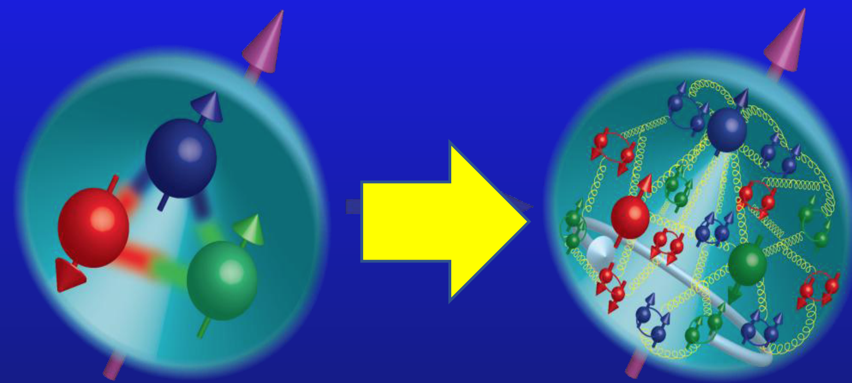

1970s cartoon Now

Z-E Meziani

Where are all the quarks and gluons -- in space and in momentum space? What is the many-body physics of all these interacting degrees of freedom.

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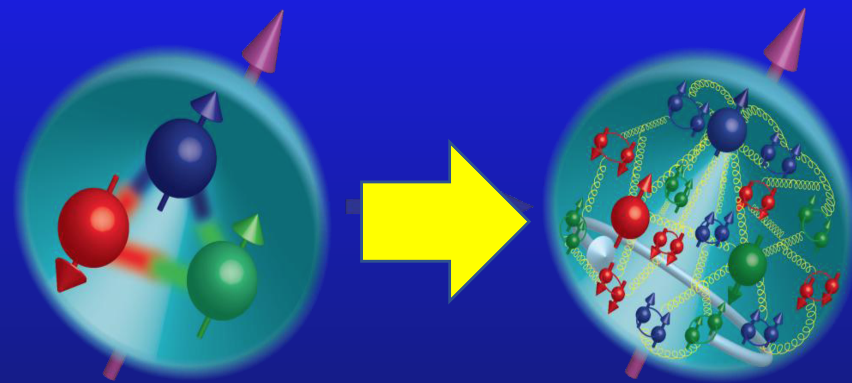

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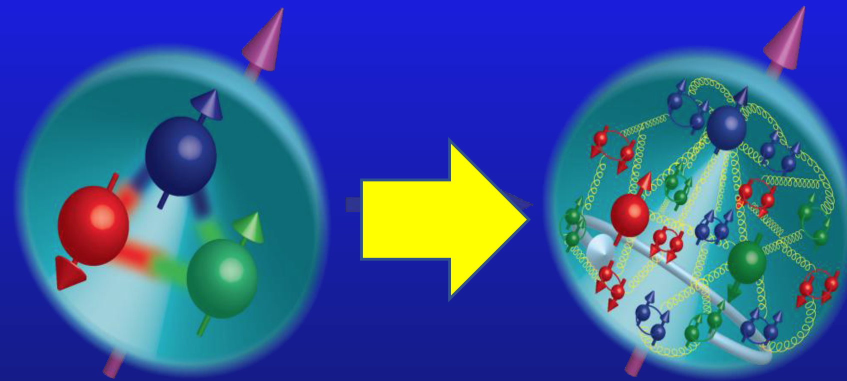


1970s cartoon Now

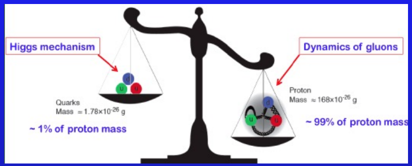
Z-E Meziani

Where are all the quarks and gluons -- in space and in momentum space? What is the many-body physics of all these interacting degrees of freedom.

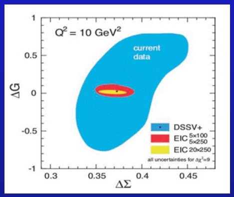
Sketch of worm vs. modern electron microscope picture (deep ocean worm)

Basic science questions for an EIC

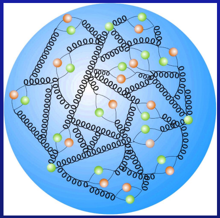
How does the nucleon get its mass?



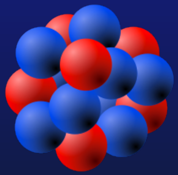
How does the spin of the nucleon arise from its elementary quark and gluon constituents?



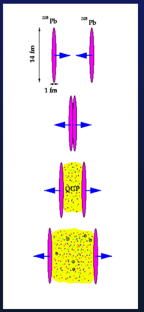
What are the emergent properties of dense systems of gluons?



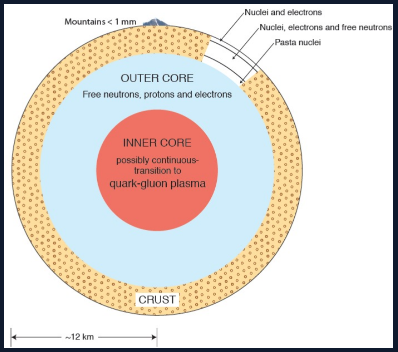
What is the internal structure of nuclei? How do nuclei differ from being a simple collection of nucleons?



What is the initial state in ultrarelativistic heavy-ion collisions?



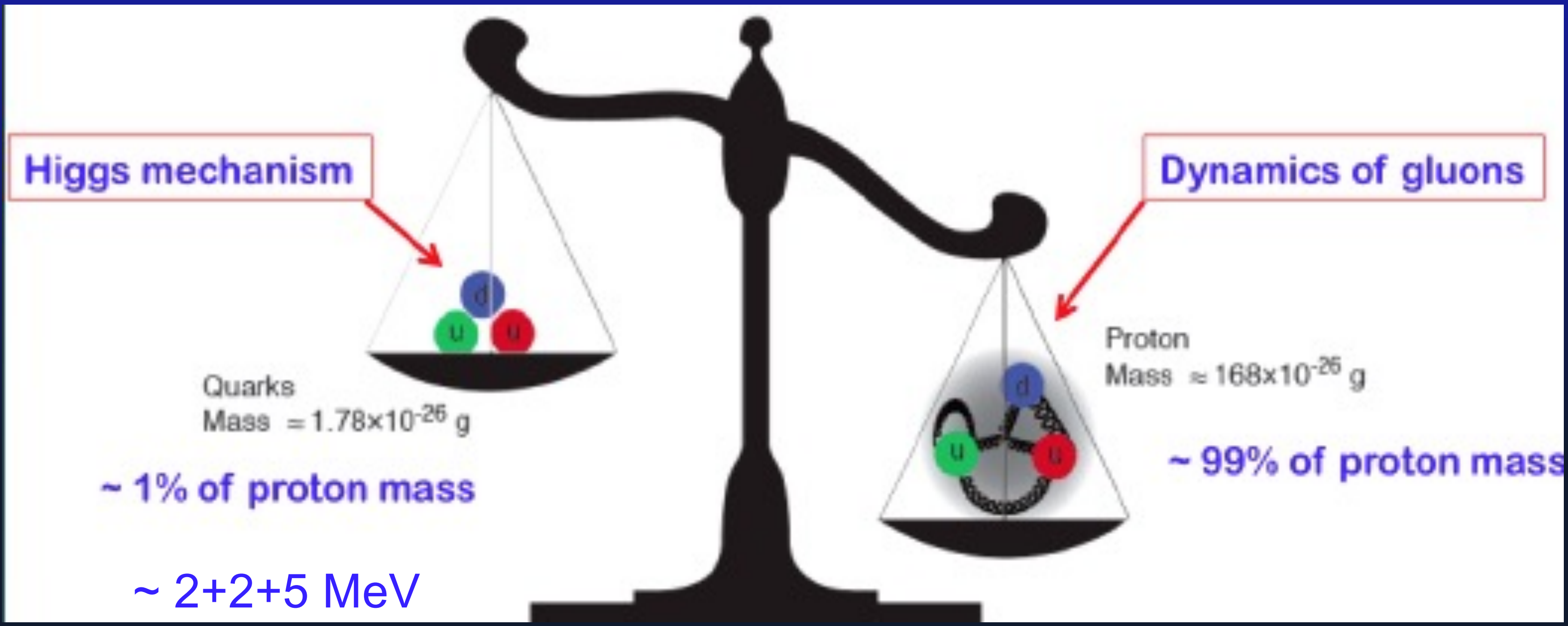
How does dense matter crossover from nucleonic degrees of freedom to quark degrees of freedom at higher density – application to neutron stars



Signals of beyond-the-standard-model physics? Effects of new particles and forces?

How does a nucleon get mass?

m_p = 938 MeV is almost 100 times greater than the masses of its valence quarks. Cannot be understood in terms of the Higgs mechanism!! Higgs => I should weigh ~ 750 gm



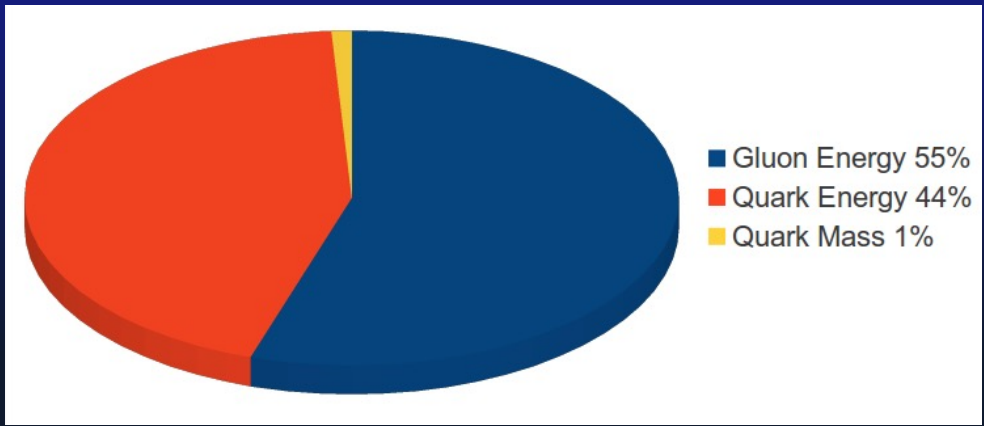
In atoms, mass = total mass of constituents minus binding energy (energy released in chemical reactions) < mass of constituents.

In nucleons, mass >> total mass of constituents.

The mass of the nucleon

Naively, the mass of the nucleon arises from all quark, anti-quark and gluon kinetic energies – from the uncertainly principle when localizing excitations within the nucleon (radius, r_p).

$\Delta E \sim \hbar c / r_p$ per quark or gluon.



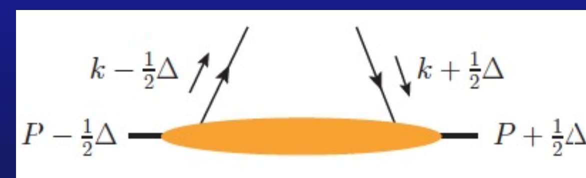
Energy distribution within nucleon

How are these components distributed (in space and momentum) in the nucleon?
And how do nuclei differ from being a simple collection of nucleons – changes of quark and gluon distributions in nuclei?

Parton (quark or gluon) distributions

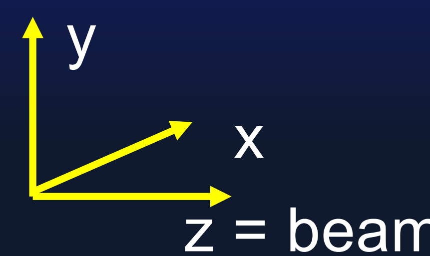
Start with correlation or Green's function $G(1,2) = \langle N|\bar{q}(2)\cdots q(1)|N\rangle$ (here q = quarks). "..." depends on process. |N> = nucleon.

Define momentum distributions in terms of different Fourier transformations.



Fully transform in momentum space

$$f(k,\Delta) \propto \int d^4x_1 d^4x_2 e^{i\Delta(x_1+x_2)/2} e^{-ik(x_1-x_2)} G(1,2)$$



Δ = momentum transfer to nucleon

k = relative momentum transfer

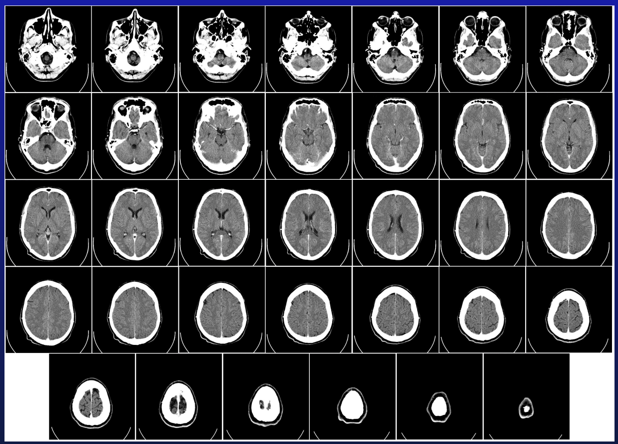
$$\vec{x}_{1\perp} - \vec{x}_{2\perp} \rightarrow \vec{k}_{\perp}$$

transverse momentum distribution alone

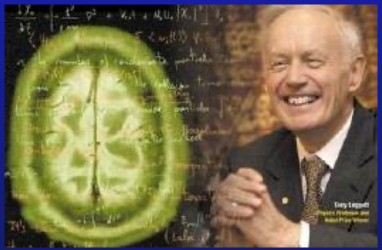
Along direction of electron momentum transfer z_1-z_2 => $x_{Bjorken}$, etc.

Tomography

Build up correlation functions by measuring transverse momentum $\vec{k}_{\perp}$ distributions at varying x. (Requires large luminosity.)

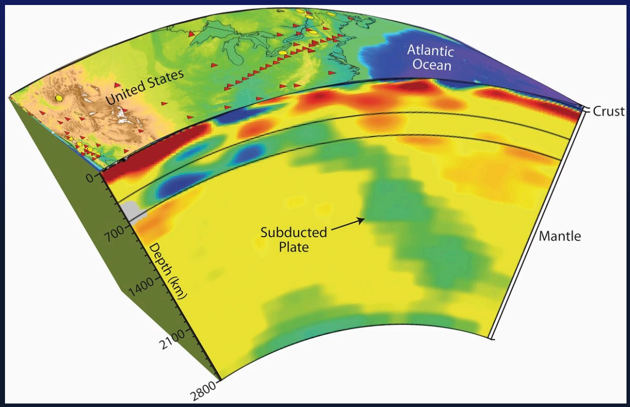


Computer assisted tomography of brain



A.J. Leggett

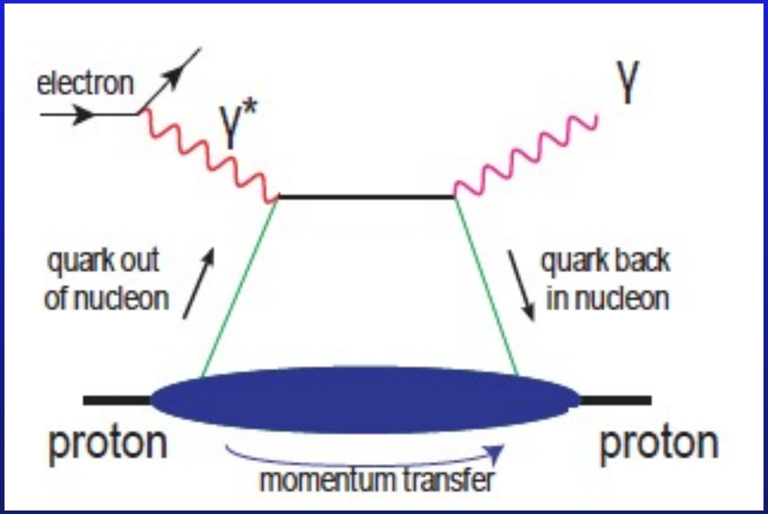
Seismic tomography, slices at varying depth; earthquakes as probes



Measure dependence of cross section on momentum transfer **q** from electron to target => information about transverse position and momentum of struck quarks and gluons.

Measuring quark and gluon transverse distributions.

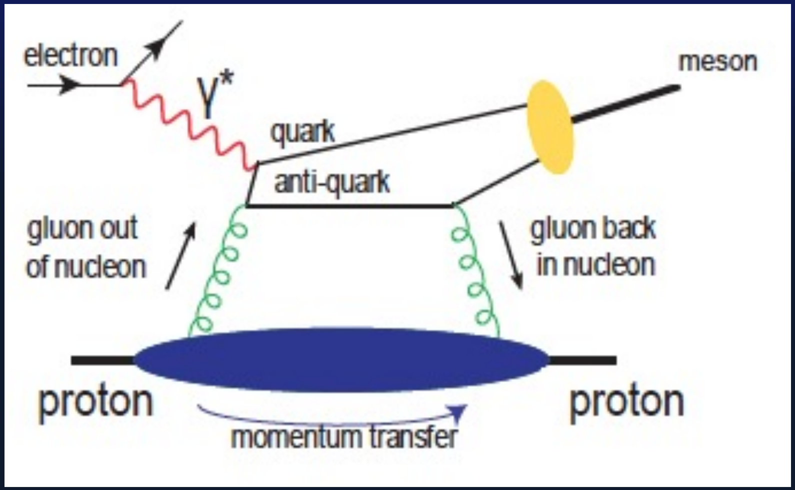
Real photon production:
Deeply virtual Compton scattering
sensitive to transverse
position of quarks.



Real meson production:
gives info on gluon distribution.

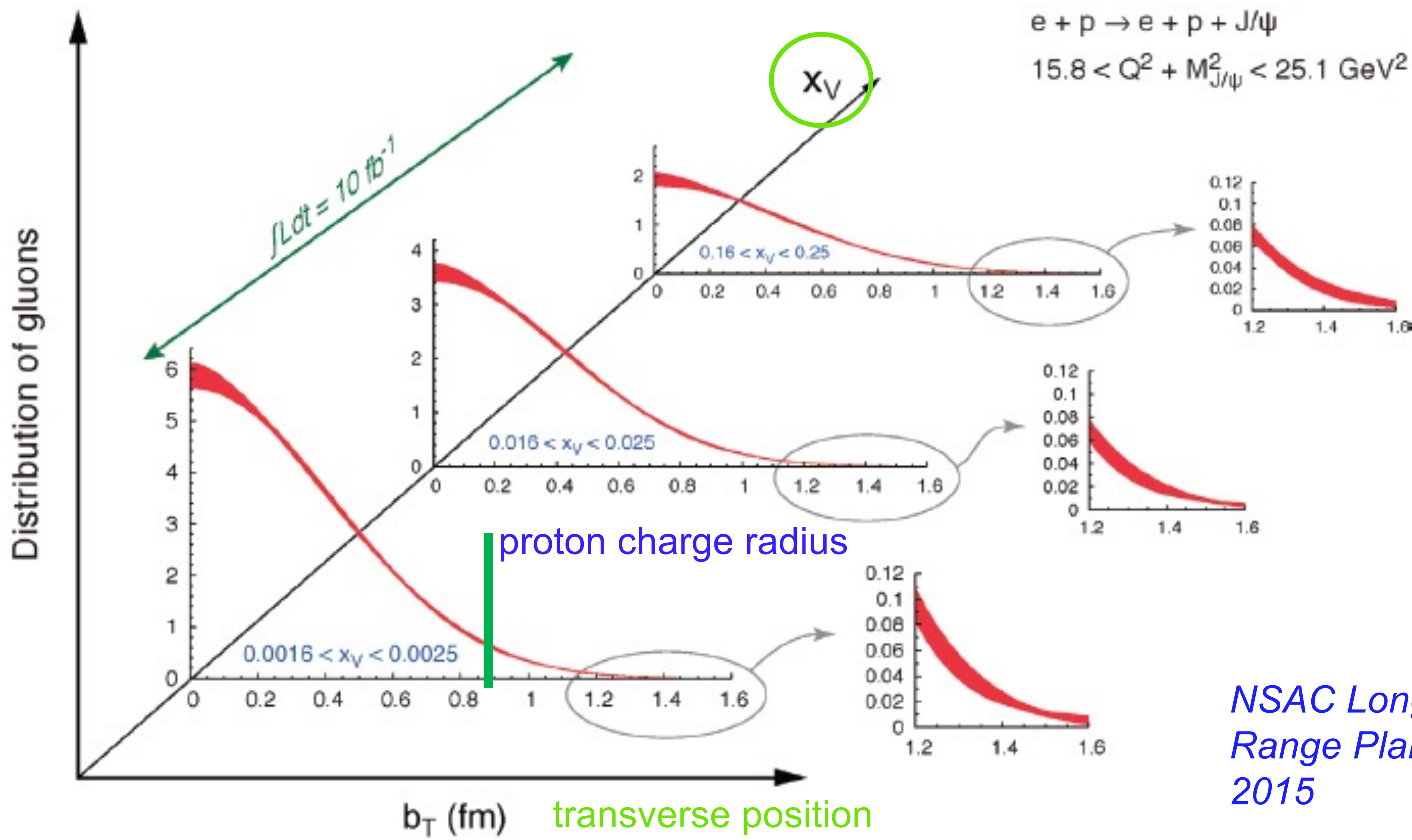
Ex., Heavy mesons J/ψ or Upsilon (Υ):

$$\gamma^* \rightarrow \bar{b}b \rightarrow \Upsilon$$



Expected transverse gluon distributions in space

Slices vs. transverse position b_T at various x

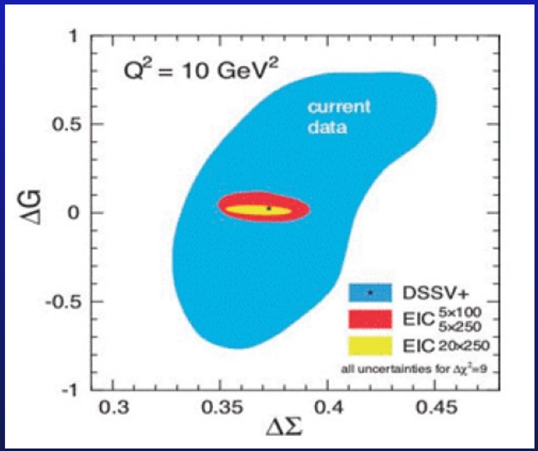


NSAC Long Range Plan 2015

How do the quarks and gluons give the proton its spin?

Proton spin is the basis of nuclear magnetic resonance (MRI) imaging.

Spin crisis:



Current estimate of contributions to spin (in units of ħ/2):

Quarks ΔΣ ~ 30-40%

Gluons ΔG ~ -70 to + 70% ?? (RHIC pp)

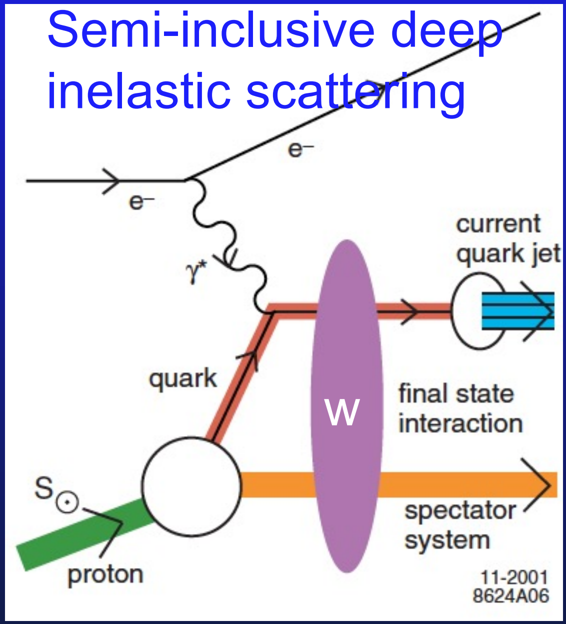
Orbital motion ?? (JLab - 12GeV)

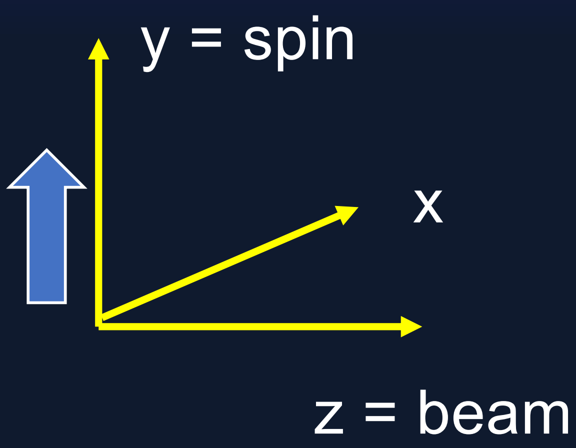
Extract orbital contributions from transverse motion measured by spatial and momentum distributions of components in deep inelastic scattering -- tomography.

Extract gluon contributions from transfer of gluon polarization to q-bar q pair, probed by polarized electrons.

Polarized beams (> 70%) critical!

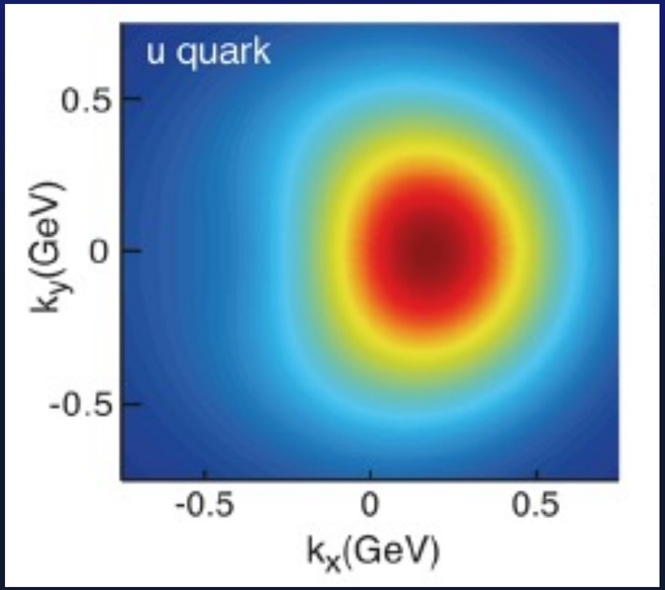
Detecting quark orbital contributions to proton spin





Proton spinning (along y) with orbital motion as well leads to asymmetry of quark distributions in x direction.

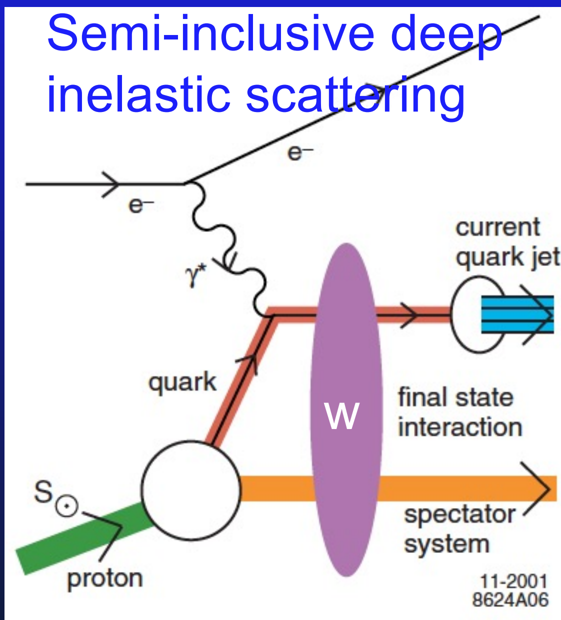
Beam along z.

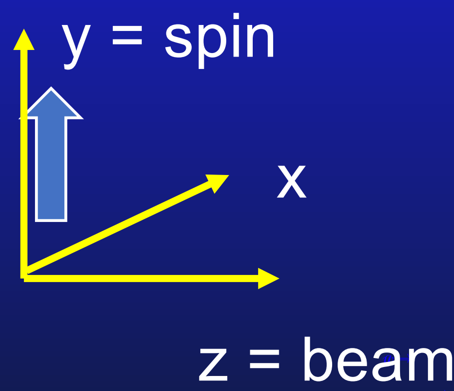


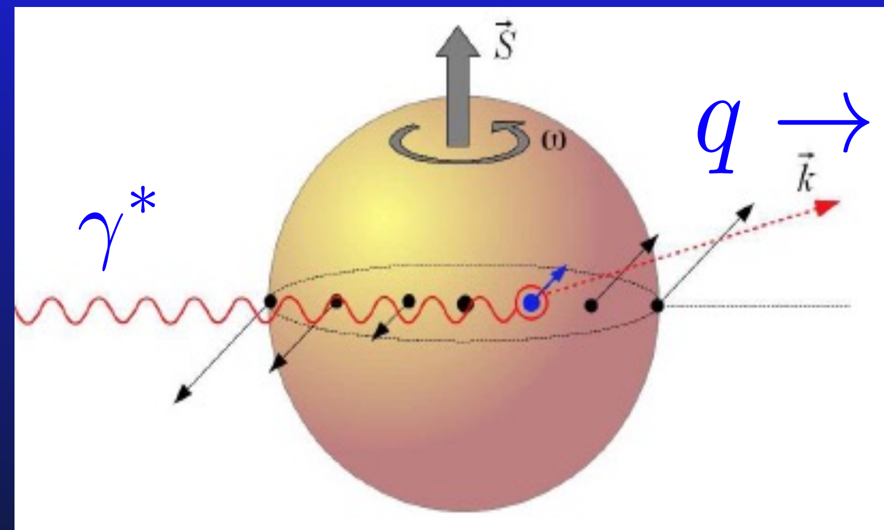
Representative asymmetry of up quark distribution for x = 0.1.

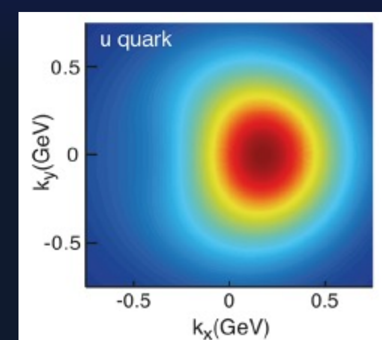
Intuitive quasi-classical scattering on rotating nucleon

Yuri Kovchegov and Matt Sievert, Phys. Rev. D 89, 054035 (2014)





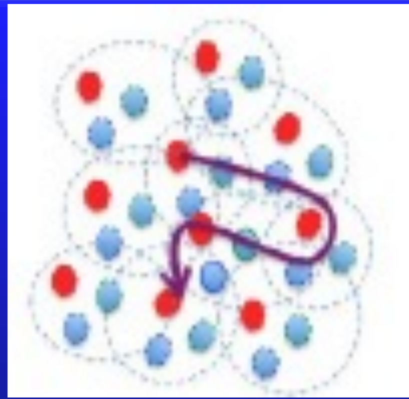




Virtual photon g* excites quark along z direction.
Quark scatters strongly (w) with other partons
=> short mean free path.

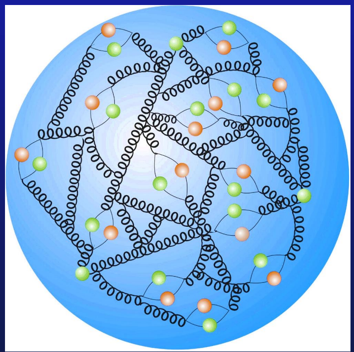
Knocked-out quark more likely emitted from far (right) side,
and thus orbital motion kicks it a bit in the +x direction.

Gluon physics



Gluons in nucleons and nuclei (as well as other hadrons) are like dark matter in the universe– unseen but crucial in holding matter together.

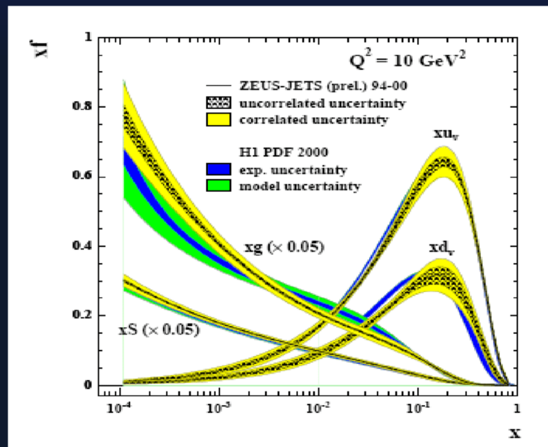
DESY

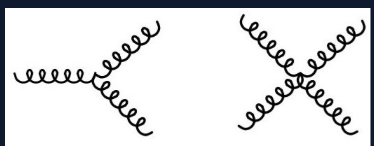


"The most precise picture of the proton"

Nucleons and nuclei are in fact complex interacting many-body systems -- beyond bags of free quarks and free gluons. Ex., nuclei exhibit composite fermions– the nucleons. Confinement!

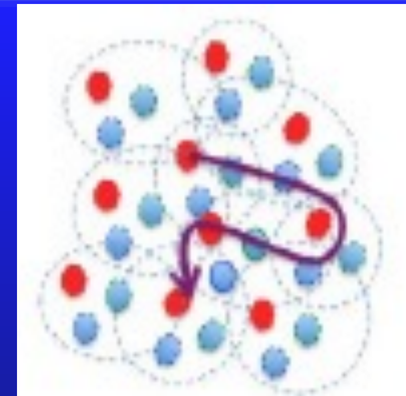
HERA => huge numbers of low momentum gluons in the nucleon -- at low x ($<10^{-4}$).
Low momentum sector (wee partons) dominated by strongly interacting gluons!. The gluon field is highly non-linear!





A new many-body system! New emergent phenomena?

Gluon physics



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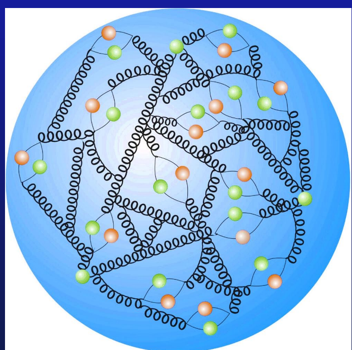
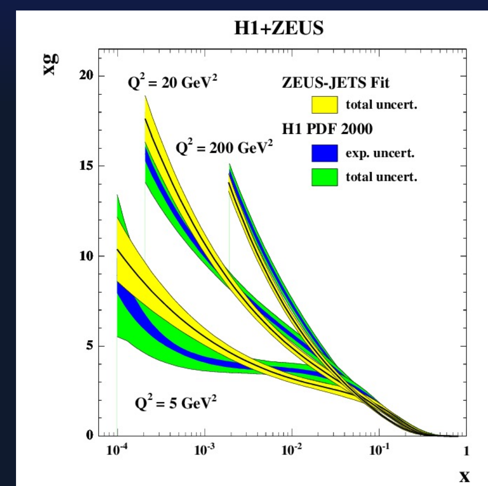
DESY



"The most precise picture of the proton"

Nucleons and nuclei are in fact complex interacting many-body systems -- beyond bags of free quarks and free gluons.   Ex., nuclei exhibit composite fermions (called nucleons). Confinement!



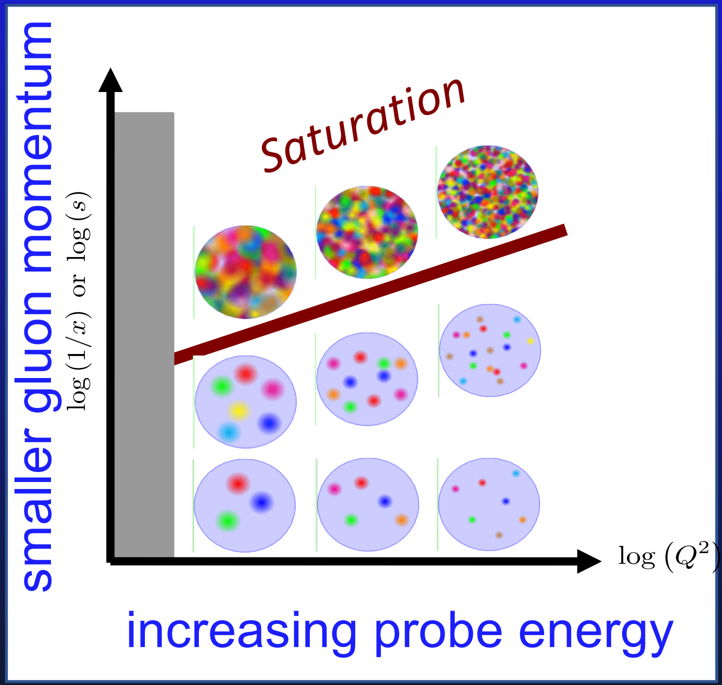
HERA => huge numbers of low momentum gluons in the nucleon -- at low x ($<10^{-4}$).
Low momentum sector (wee partons) dominated by strongly interacting gluons!.   The gluon field is highly non-linear!

A new many-body system! New emergent phenomena?

Gluon self-interactions become important when there are enough: at small x

Scale of saturated gluonic matter: Q_s

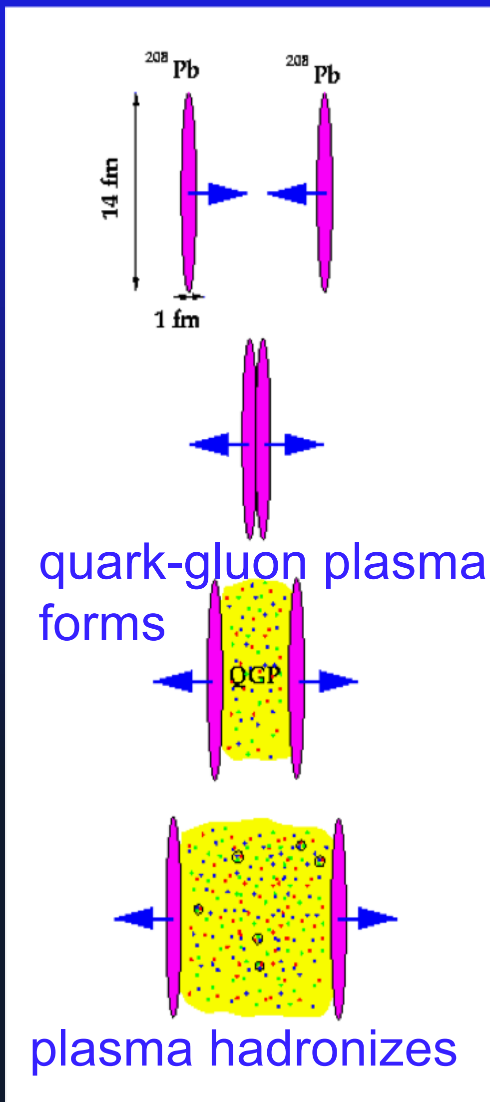
At HERA (318 GeV c.m.) $Q_s \sim 1$ GeV



First approximation, dense cloud of gluons forms a Bose condensate – "color glass condensate."

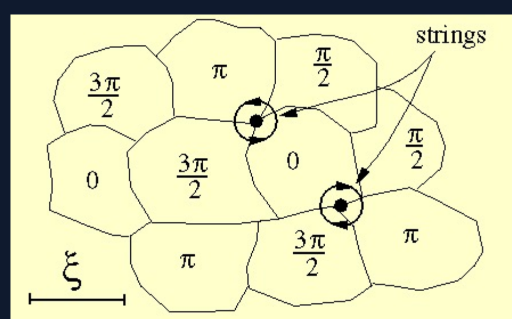
Excitations of saturated gluonic matter? Topology?

Connections to heavy ion collisions:

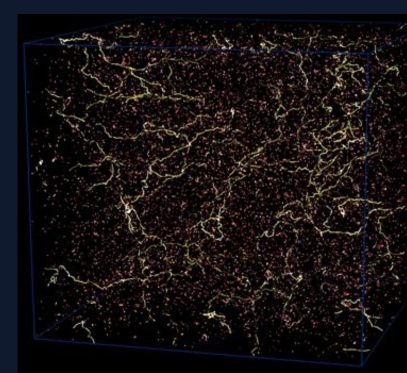


Saturated gluonic matter reachable at a sufficiently energetic EIC. Describes initial state in ultrarelativistic heavy ion collisions. Bose-condensed gluonic matter (color-glass condensate, ...). Condensate is metastable, decaying into quark-gluon plasma.

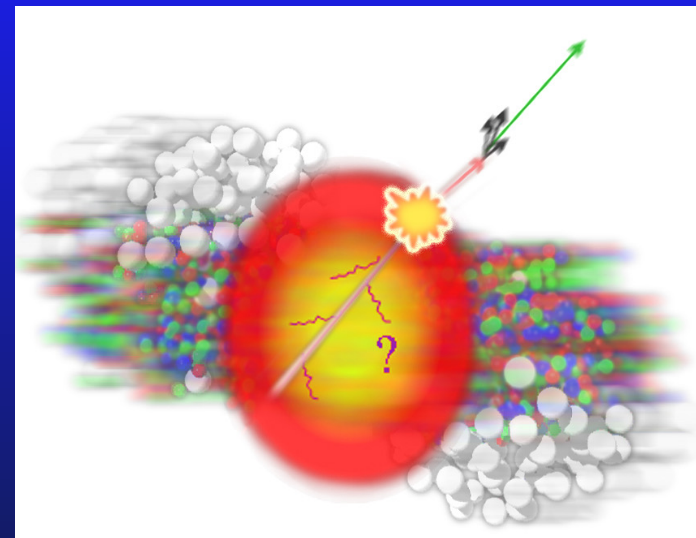
As in early universe, form topological defects, e.g., handedness asymmetry of produced $q \bar{q}$ pairs (chiral magnetic effect) related to the structure of the color field in saturated gluonic matter.



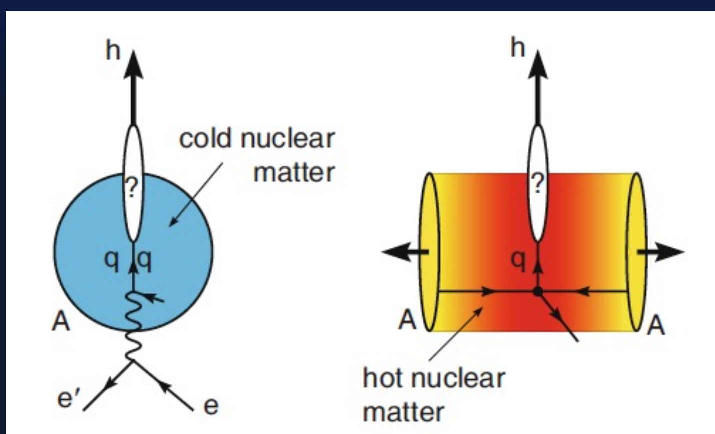
cosmic strings

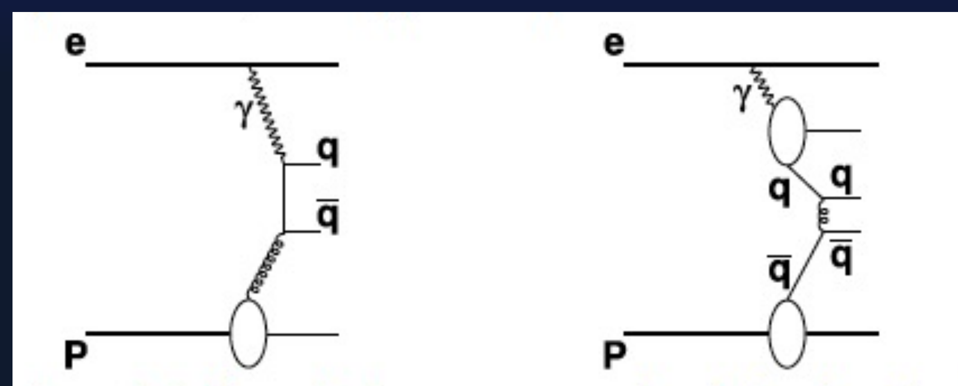


Particle jets in electron-ion and ion-ion collisions



asymmetric jet in AA collision

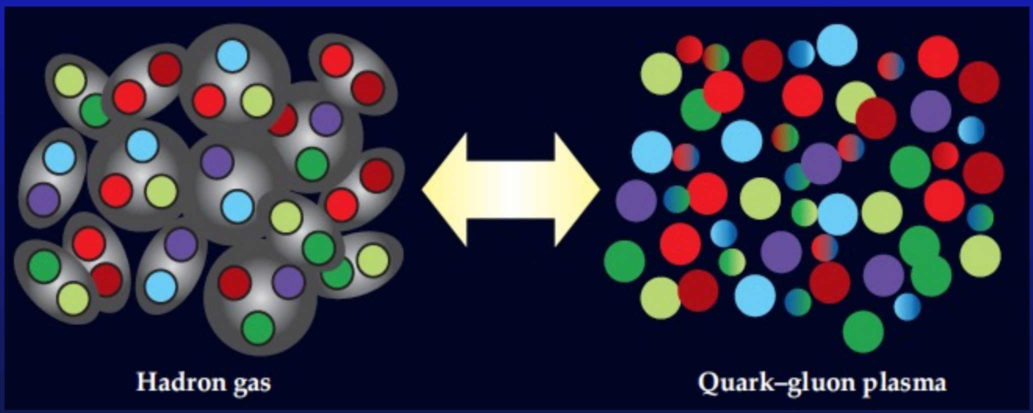


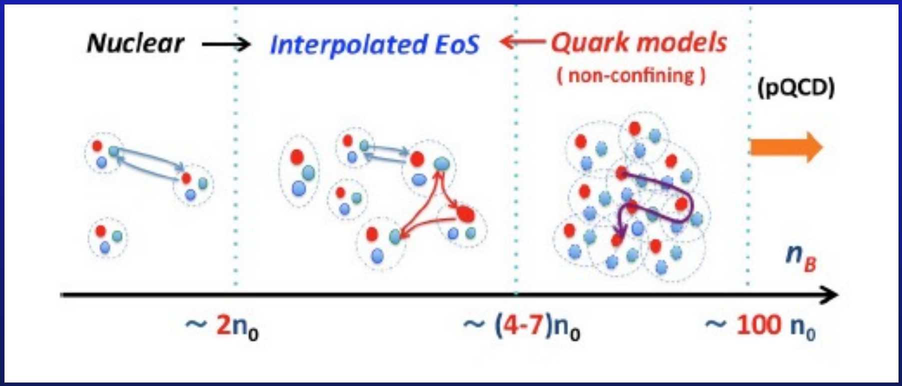


Jets probe quark and gluon distributions

Dense matter and neutron stars:

Study transition from cold nuclear matter to quark matter – vital for neutron stars. What is energy density vs. baryon density?





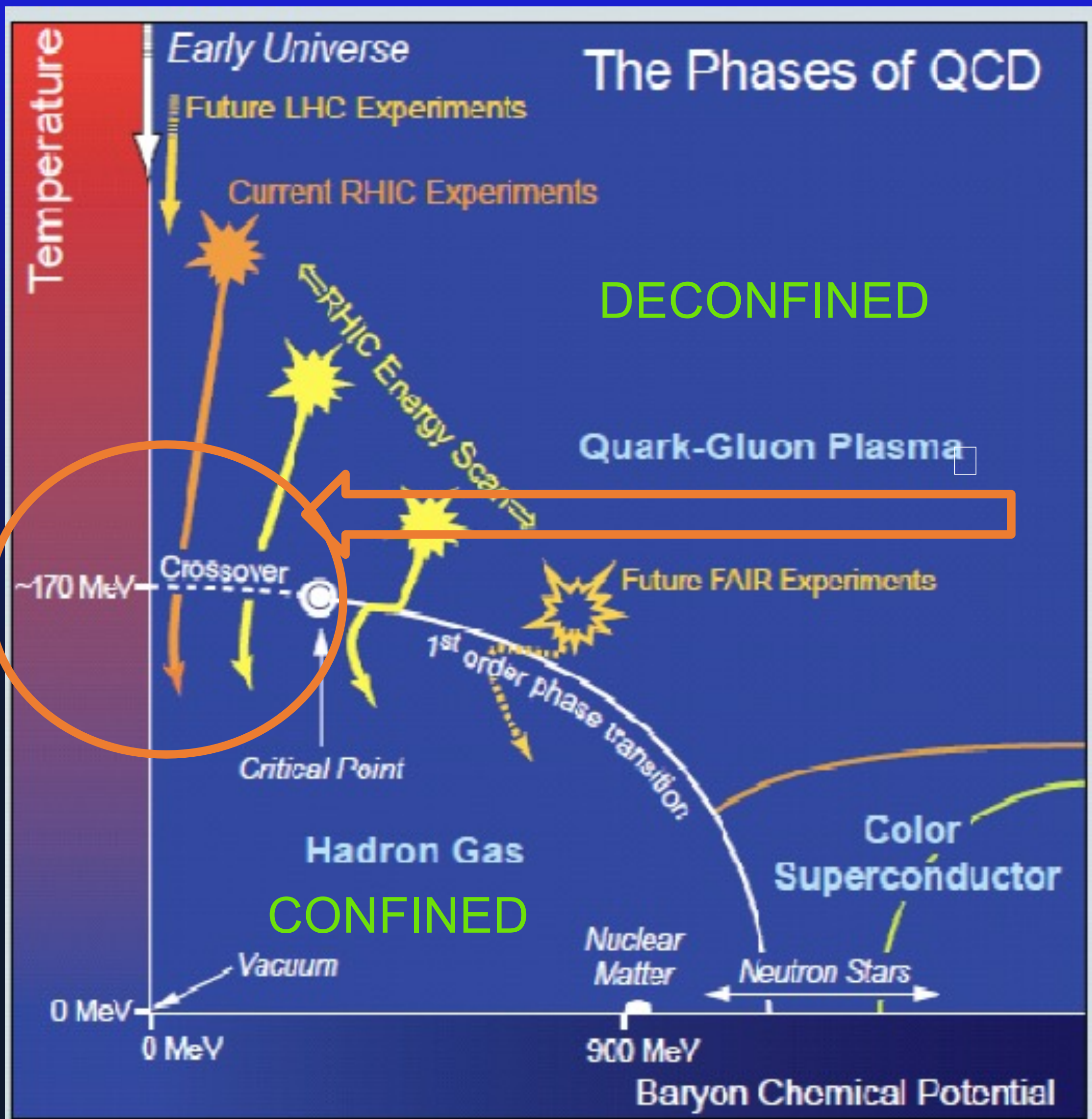
Expect "smooth" transition from nucleons to quarks

Gluon (and quark) distributions in nuclei at finer and finer scales should shed light on transition from nucleonic to quark degrees of freedom as density of matter increases.

(Can mapping of energy-momentum tensor (stress-energy tensor) in eA collisions reveal pressure vs. baryon density in dense matter?)

Modern phase diagram

Crossover at zero net baryon density



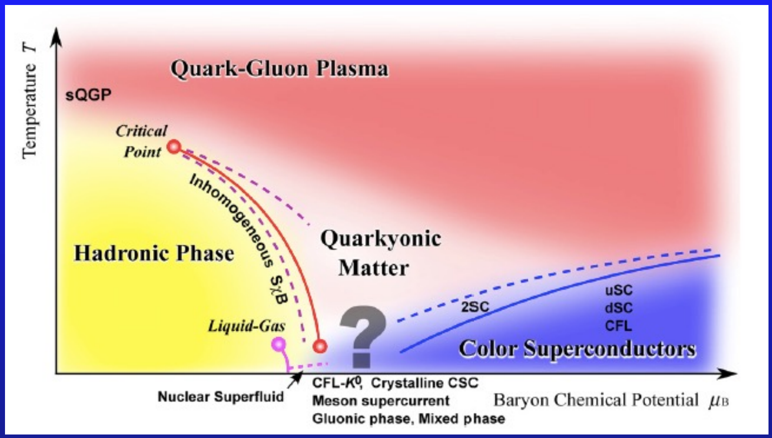
QCD lattice gauge theory – for finite light quark masses -- predicts crossover from confined phase at lower T to deconfined phase at higher T.

Wuppertal-Budapest lattice collab.
S. Borsanyi et al., PLB (2014)
A. Bazavov et al., PRD (2014)

Do quarks roam freely in the deconfined phase?
If so, they must also roam freely at lower T.

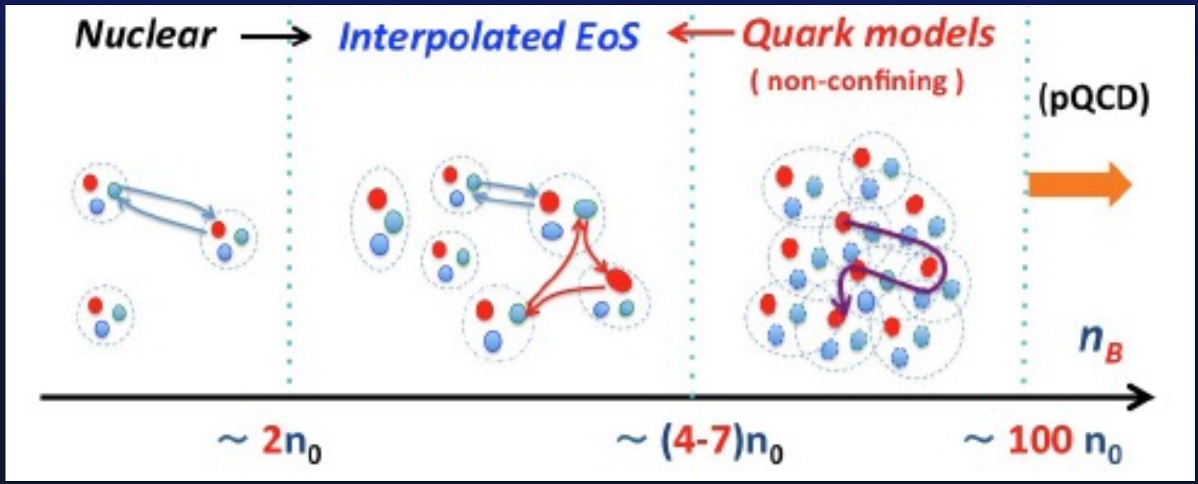
Are free quarks really running about in this room?

No free quarks even above the crossover!



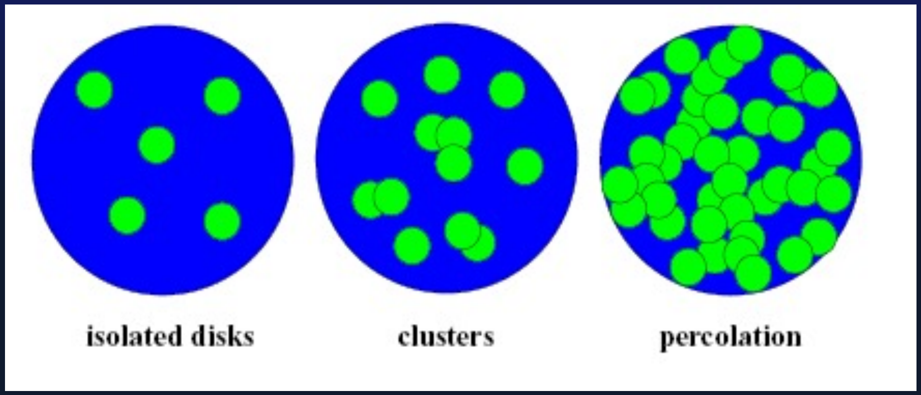
In confined region quarks are inside hadrons. Also have quarks and antiquarks in the QCD forces between hadrons. With temperature, form larger clusters, which percolate at the crossover. Expect similar percolation of clusters with increasing density.

In deconfined regime clusters extend across all of space.



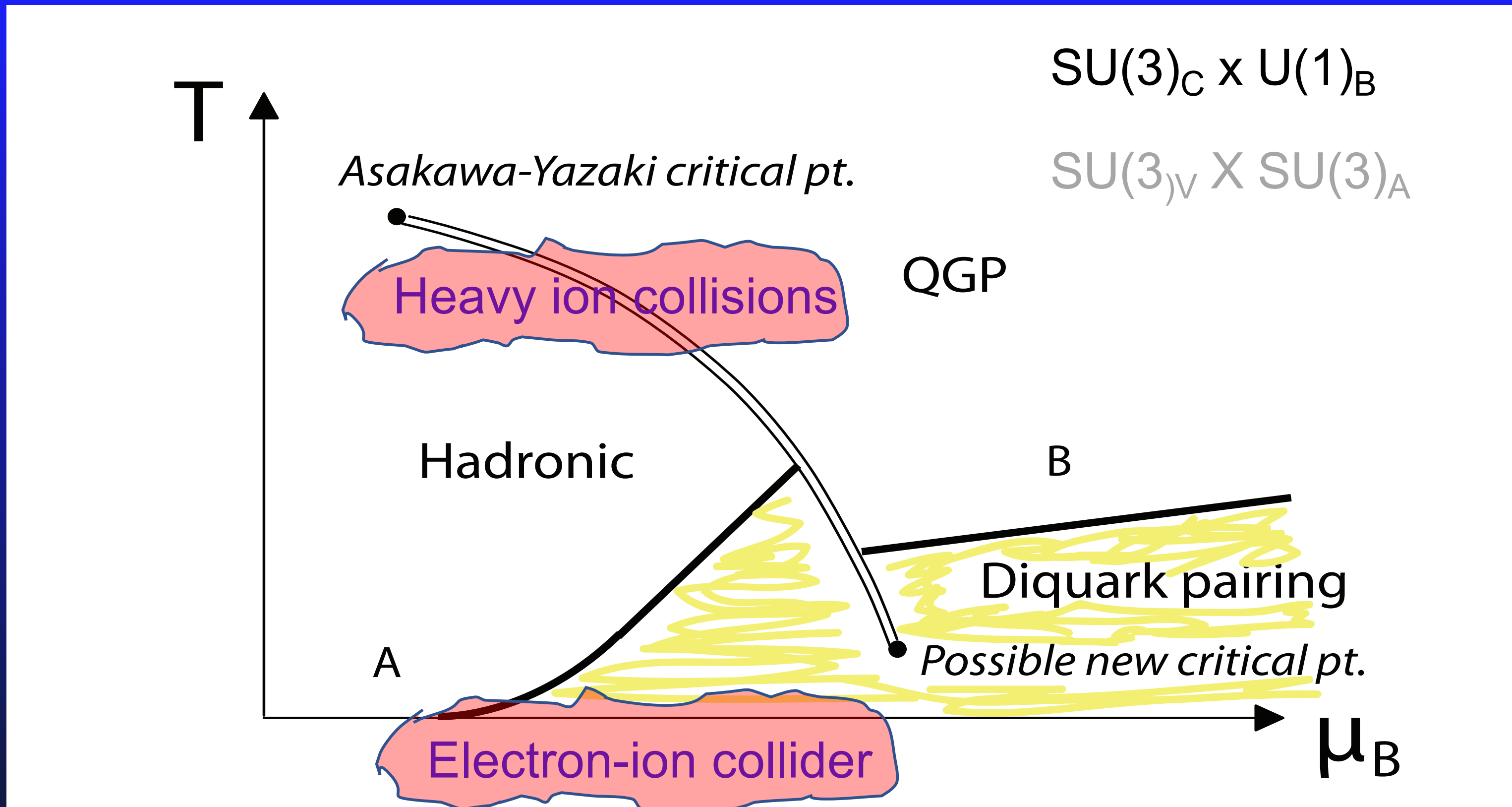
Need theoretical description of this evolution

Percolation

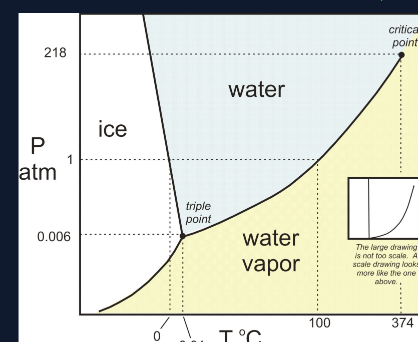


$n_{perc} \sim 0.34\ (3/4\pi\ r_n^3)\ \ fm^{-3}$

Quarks can still be bound even if deconfined.



Critical points similar to those in liquid-gas phase diagram (H_2O). Neither critical point necessary!!



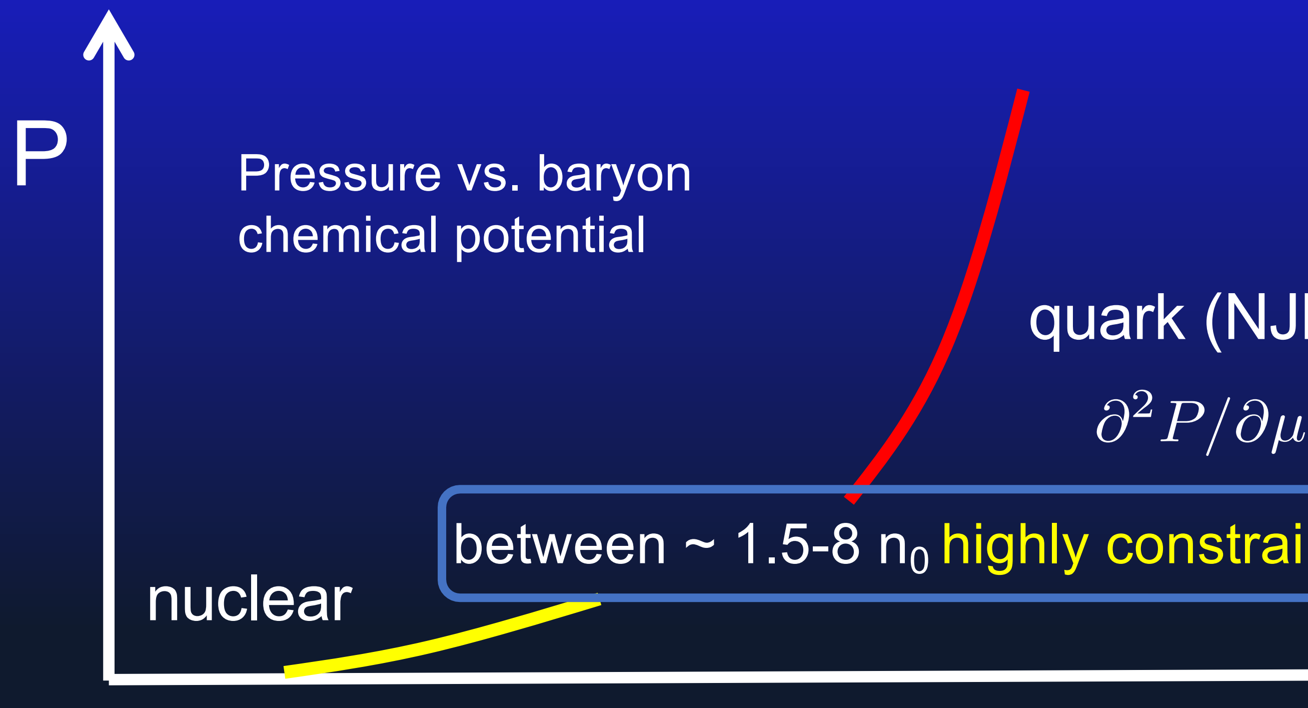
Can go continuously from A to B around the upper critical point. Liquid-gas phase transition.

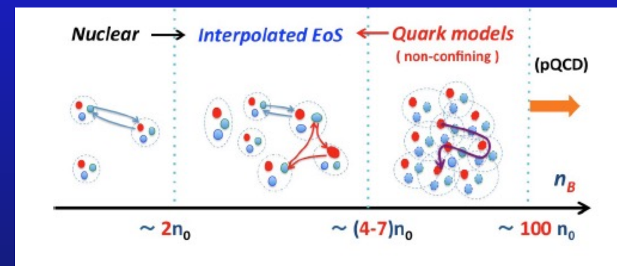
In lower shaded region have BCS pairing of nucleons, of quarks, and possibly other states (meson condensates). Different symmetry structure than at higher T.

QHC21 (quark-hadron crossover) equation of state:

QHC19 --GB, S. Furusawa, T. Hatsuda, T. Kojo & H. Togashi, Ap.J. 885:42 (2019)
QHC21 --T. Kojo, GB, &T. Hatsuda, Ap. J.   934:46 (2022).

Have good idea of equation of state at nuclear and at high densities.





$$\partial^2 P/\partial\mu^2 = \partial n_B/\partial\mu \geq 0$$

between ~ 1.5-8 n_0 highly constrained by continuity & stability

Quarks in Nambu-Jona-Lasinio (NJL) model with universal repulsive short-range qq coupling (*Kunihiro*) and diquark (BCS) pairing interaction

$$\mathcal{L}_V^{(4)} = -g_V\left(\overline{q}\gamma^\mu q\right)^2$$

$$\mathcal{L}_d^{(4)} = H \sum_{A,A'=2,5,7} [(\bar{q}i\gamma_5\tau_A\lambda_{A'}C\bar{q}^T)(q^T Ci\gamma_5\tau_A\lambda_{A'}q)$$

Nuclear matter equation of state below 1.5 n_0

For thermodynamic stability need large diquark pairing interaction, H

$$\mathcal{L}_d^{(4)} = H \sum_{A,A'=2,5,7} [(\bar{q} i\gamma_5 \tau_A \lambda_{A'} C \bar{q}^T)(q^T C i\gamma_5 \tau_A \lambda_{A'} q)$$

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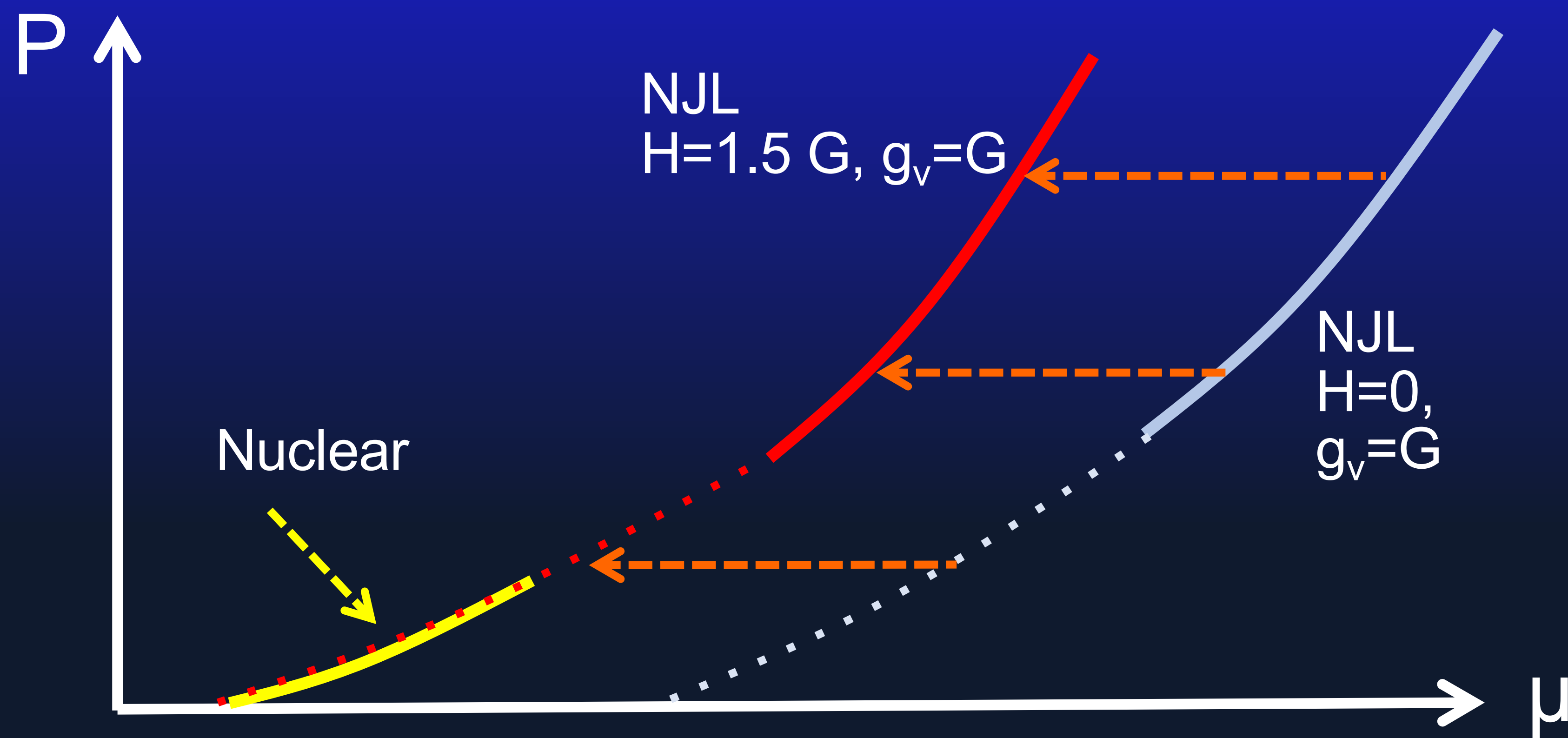
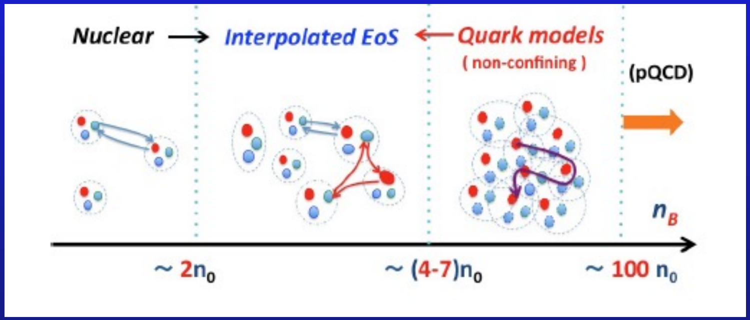


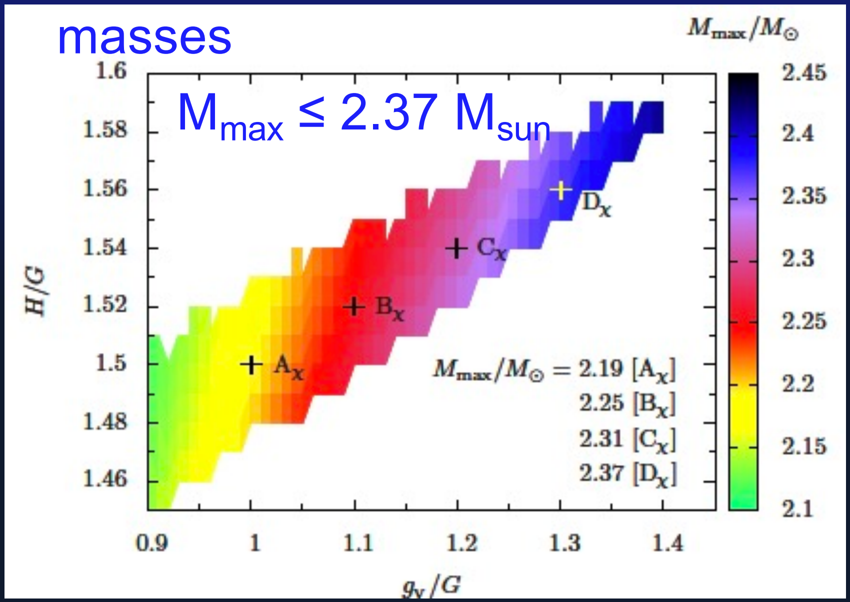
Expect onset of stronger 2-body correlations, with decreasing density, as quark matter comes nearer to becoming confined

QHC21 (quark-hadron crossover) eqn of state

T. Kojo, GB, & T. Hatsuda, Ap. J. 934:46 (2022).

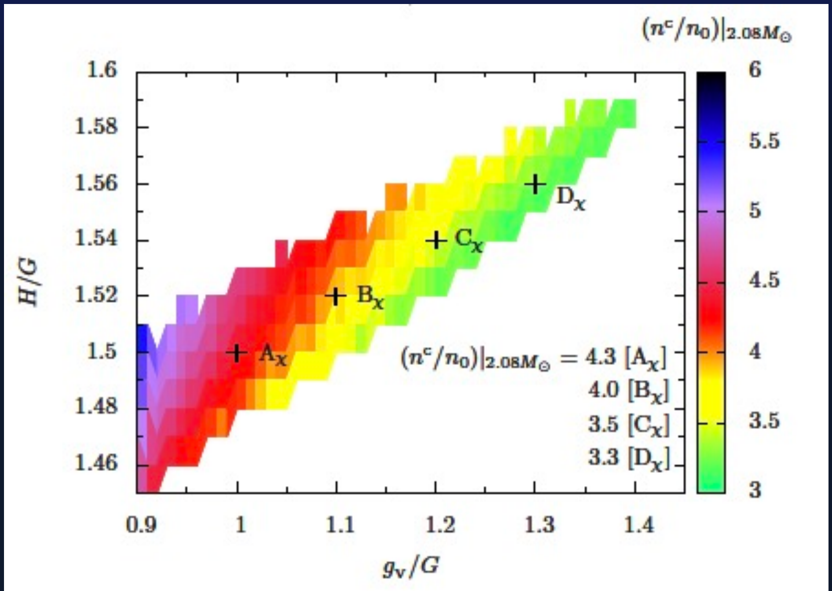


Parameters g_v and H must be in colored region so that
speed of sound ≦ speed of light.



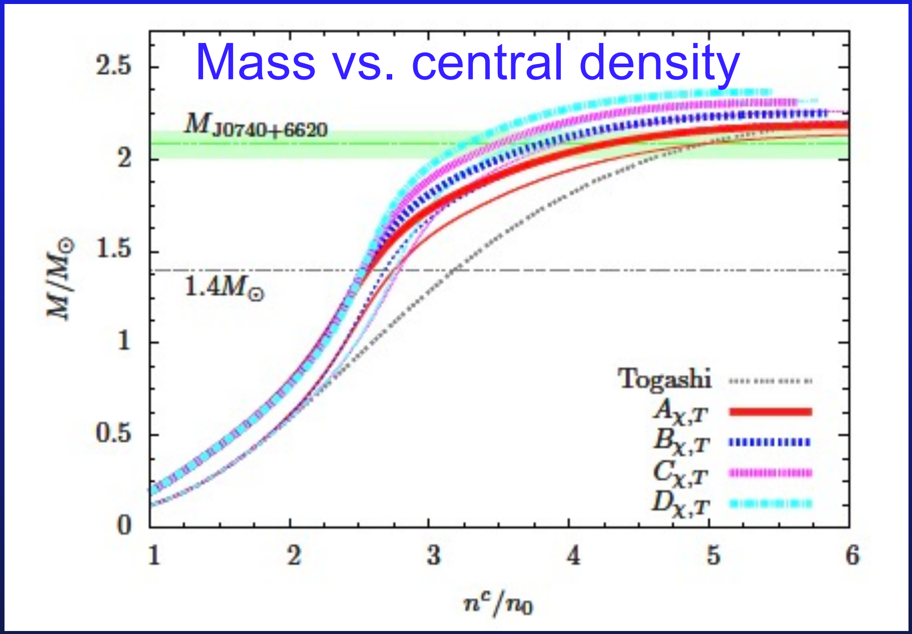
Strange quarks included

Further restricted by maximum neutron star mass > 2.08 M_{sun}



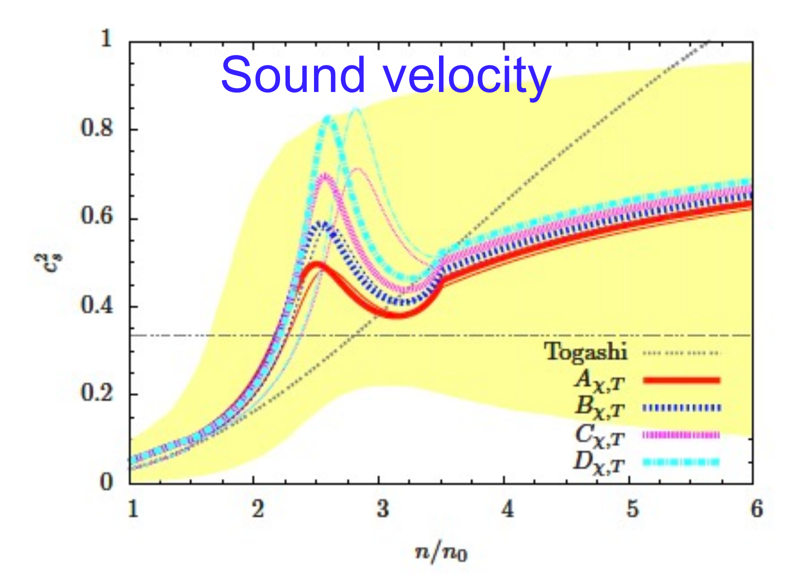
2.08 M_{sun} => $g_v \gtrsim 0.9$ G

Quarks beginning to enter neutron stars



Central density of PSR J0740+6620 ~$5n_0$.
Well above densities where pure hadronic calculations are valid. Entering transition to strongly interacting quark matter.

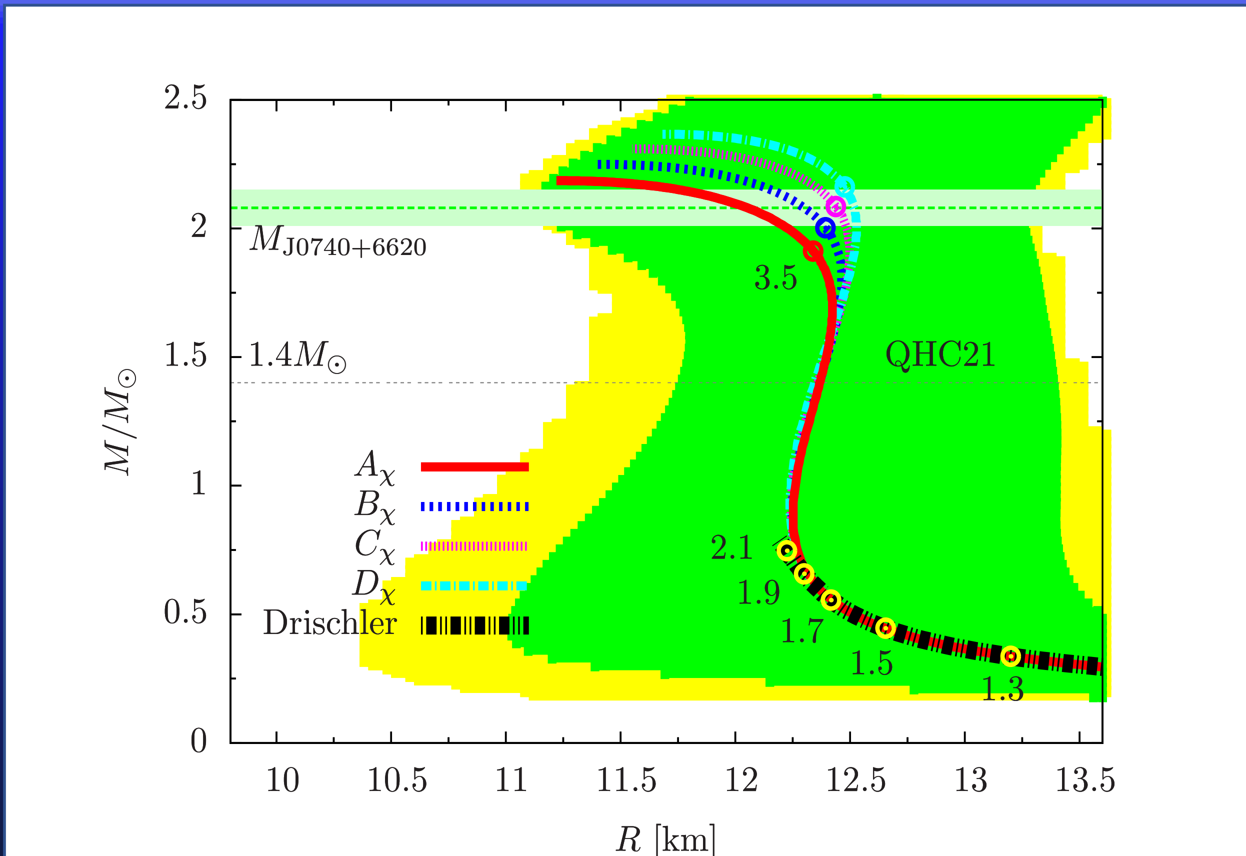
Cores of higher mass stars could reach beyond transition.
Fully developed quark matter in cores? Need fully microscopic calculations of matter undergoing transition from nucleonic to quark degrees of freedom.



Peak not seen in nucleonic eq. of state

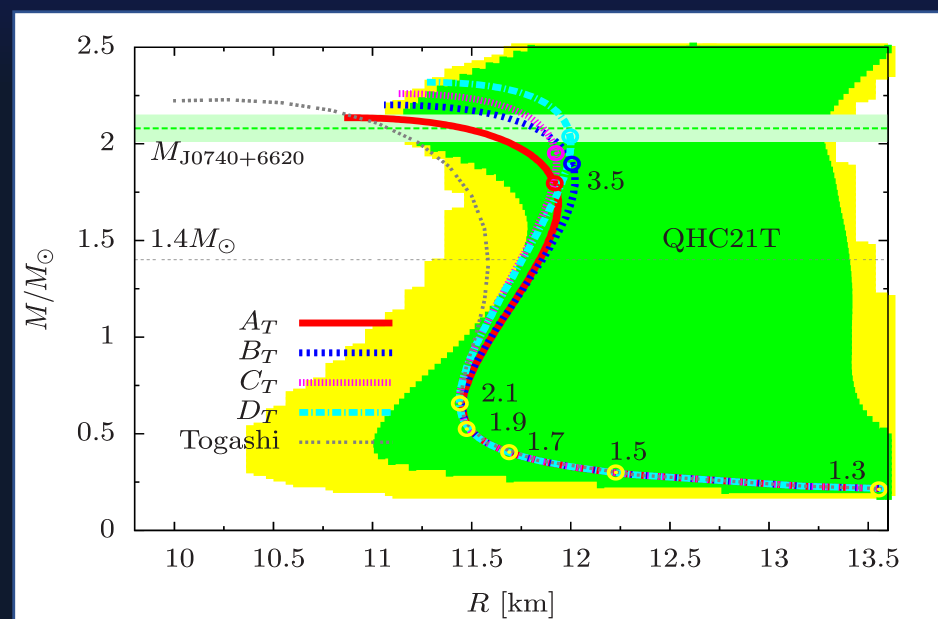
Chiral EFT + transition to quark matter in excellent agreement with NICER inferences of radii. Rapid pressure rise!

Central density for 2.08 M_{sun} ~ 3.6 n_0

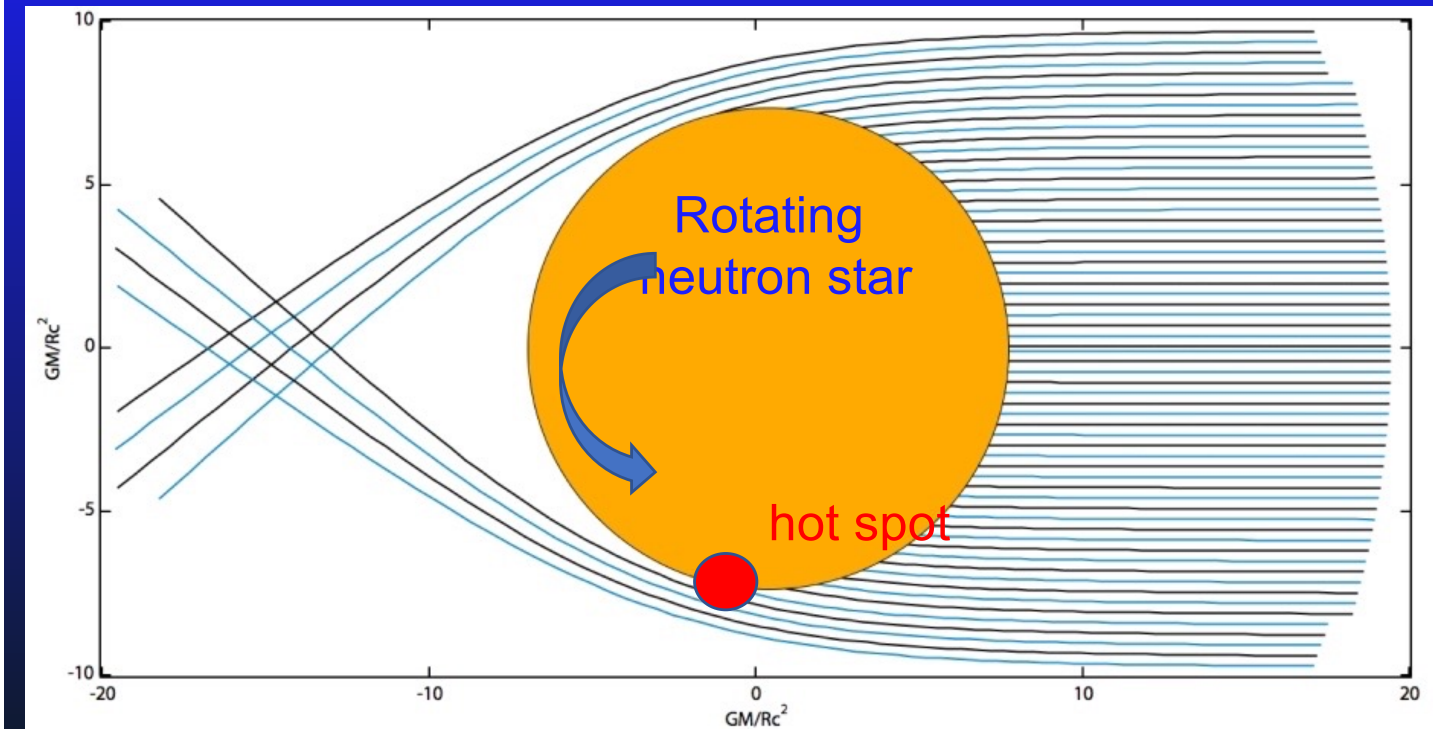


QHC21

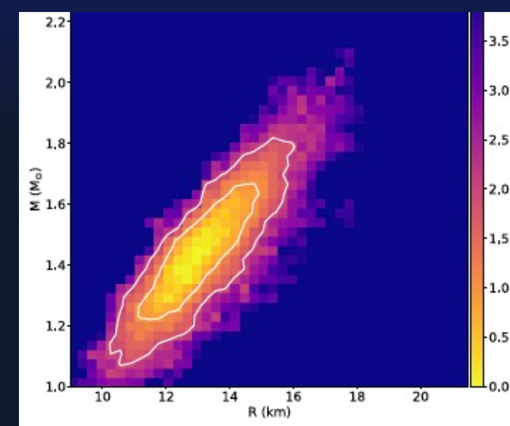
APR-Togashi too soft in nuclear regime.



Track hot spots on neutron star. Light bending by star enables one to see spot "behind" star. Bending depends on M and R.

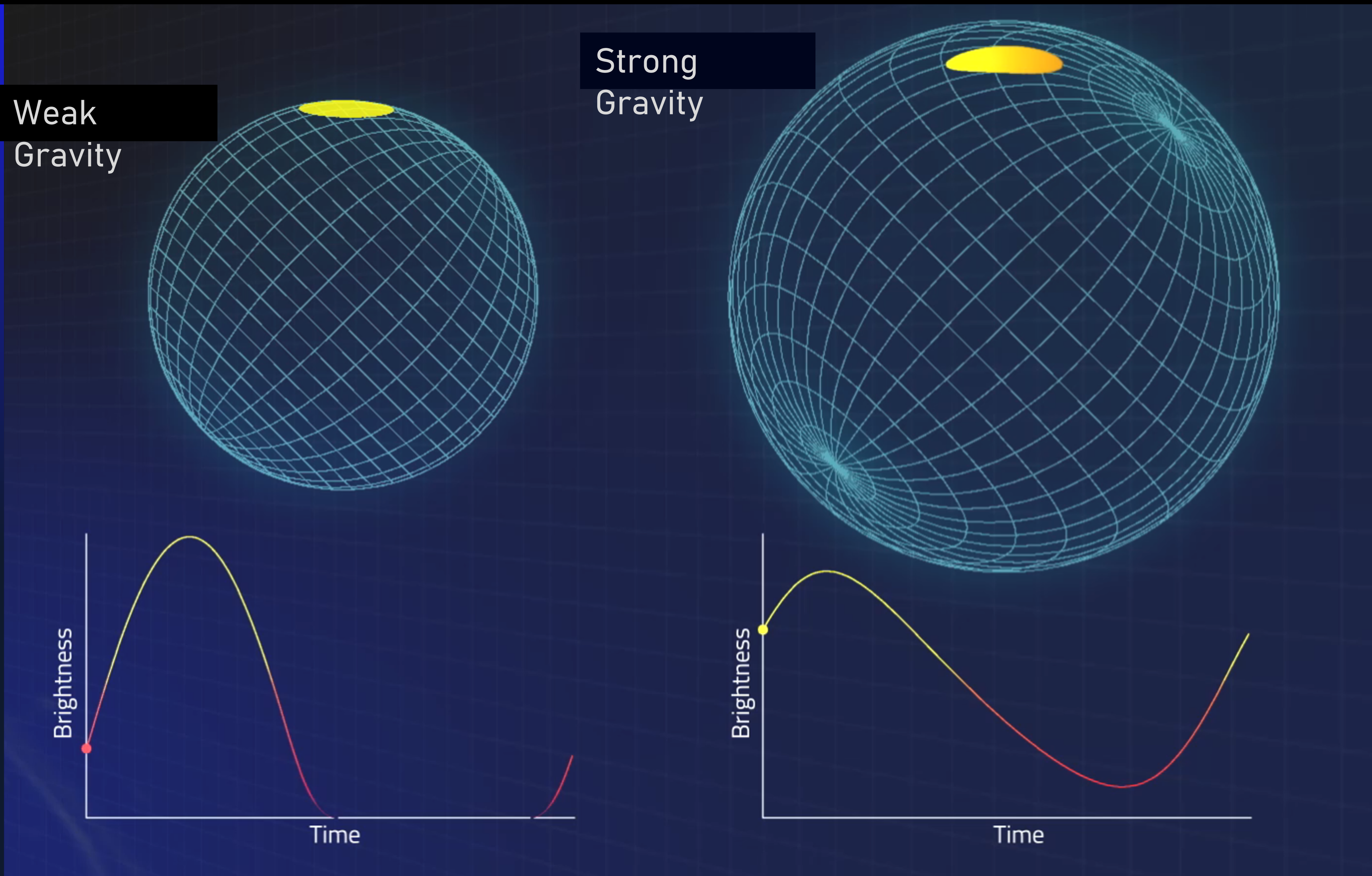


Measure amplitudes and phases
in different frequencies, construct
model of hot spots to interpret data



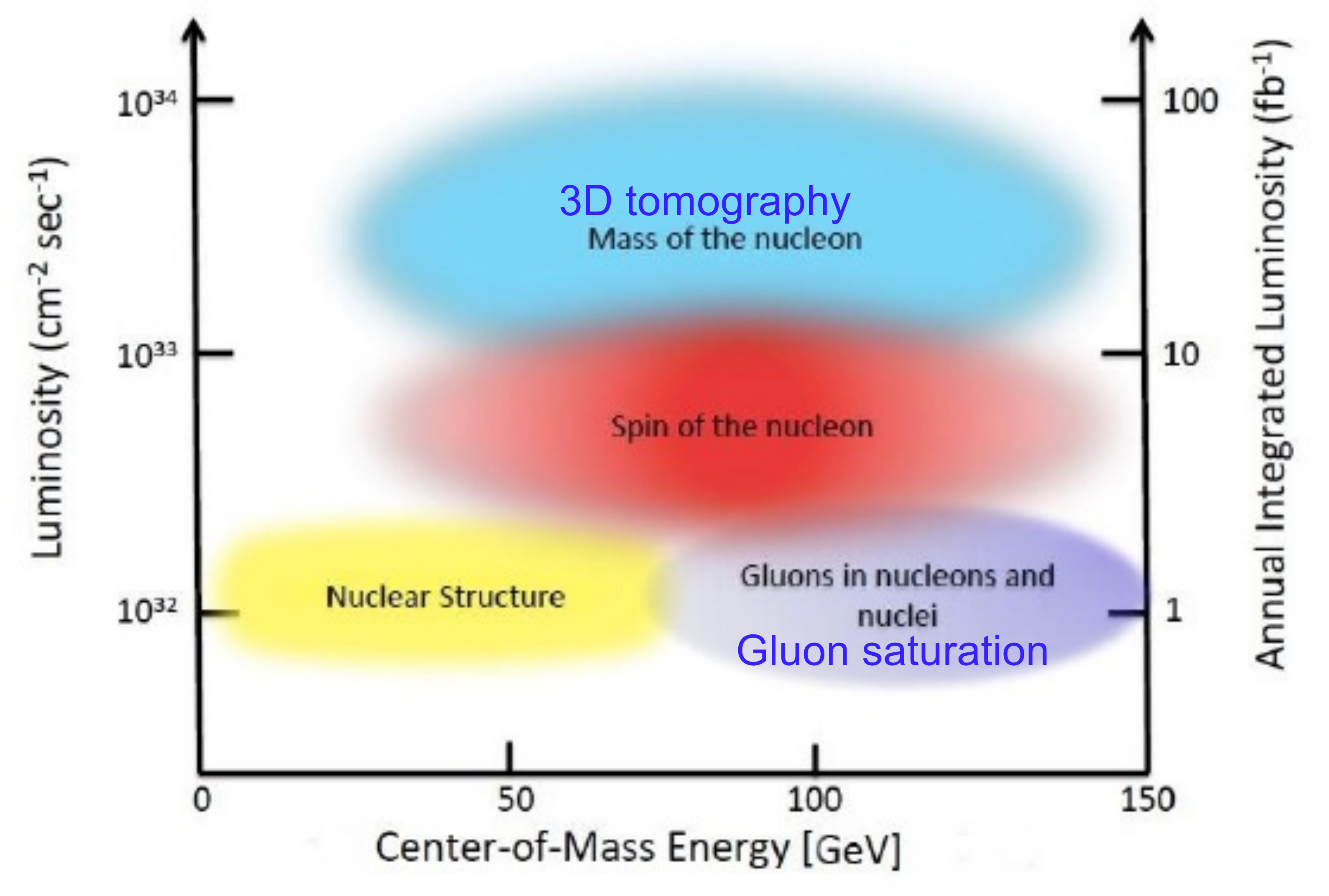
PSR J0030+0451
Mass vs Radius
Miller et al.,Ap.J. (2019)

Pulse Profile Modeling (PPM)



Credit: Morsink/Moir/Arzoumanian/NASA-GSFC

The accelerator



(electron-proton) c.m. energy - luminosity landscape

Luminosity measures the rate of collisions: ($L\sigma$ = event rate)

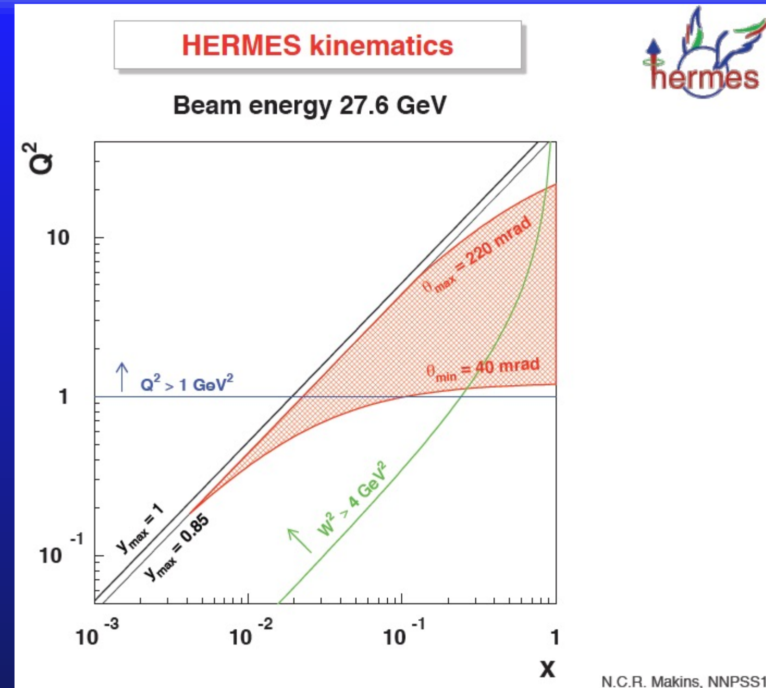
Kinematic range

$Q^2 = x(1- E_{e'}/E_e)s \leqq x\, s$

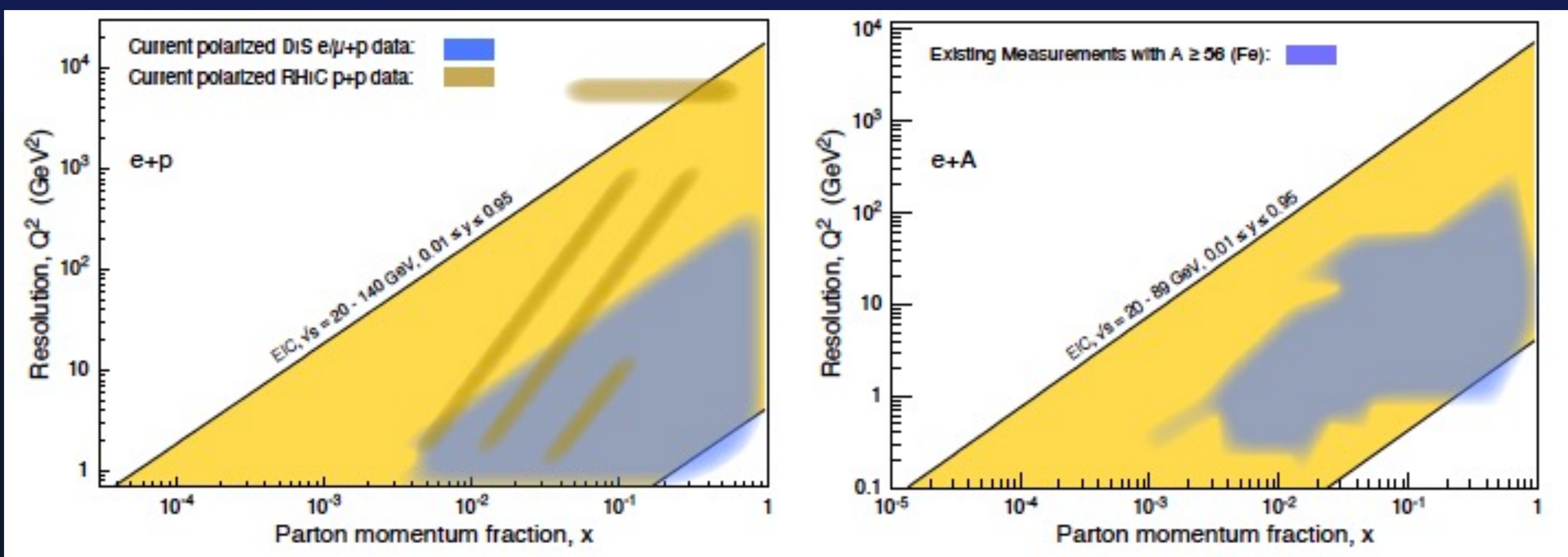
$\sqrt{s}$ = center of mass collision energy

$x \gtrsim 10^{-4}$ for $Q^2 \sim 1$ GeV^2, $\sqrt{s} \sim 100$ GeV

($Q^2 \sim 1$ GeV^2 lower limit for inelastic scattering)

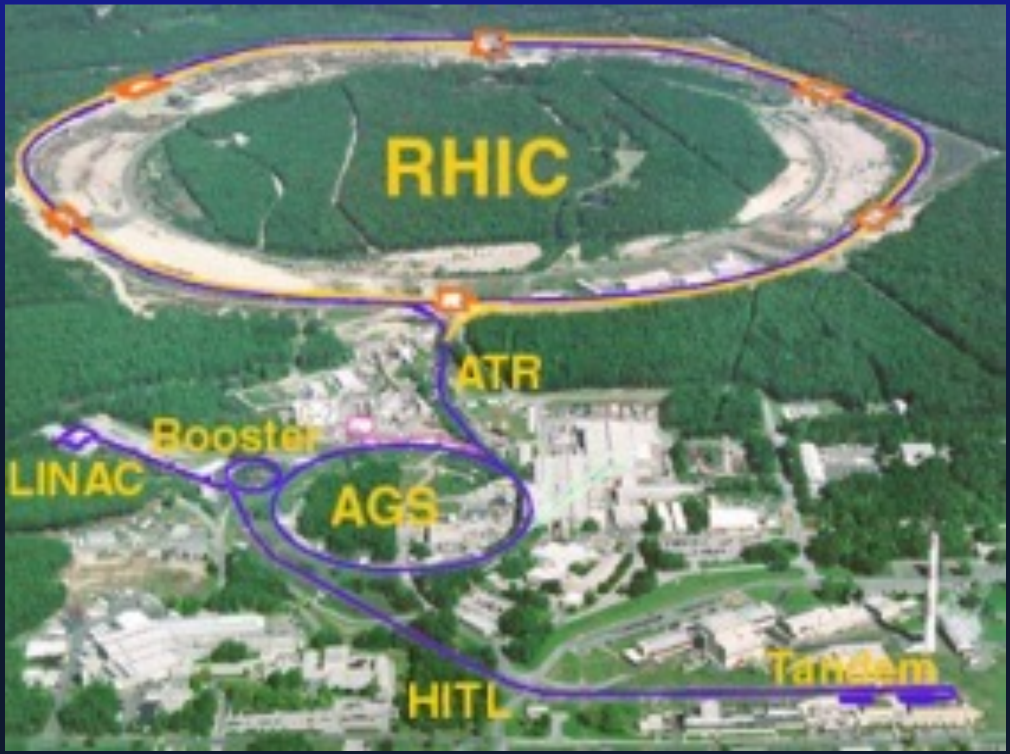


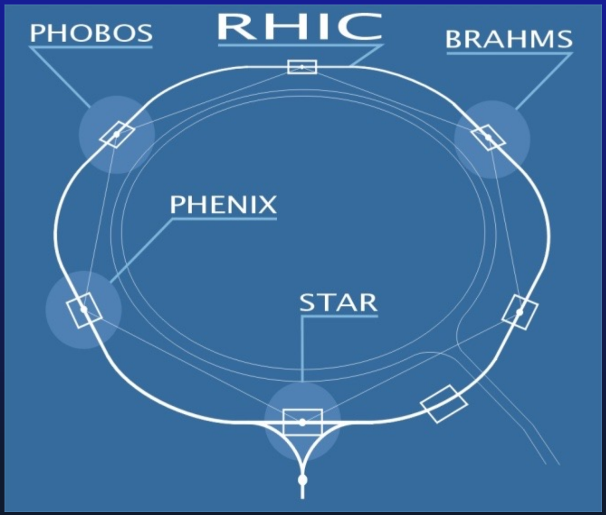
Small x region needs large beam energy



Relativistic Heavy Ion Collider (Brookhaven) since 2000 now becoming the EIC (joint project with Jefferson Lab)

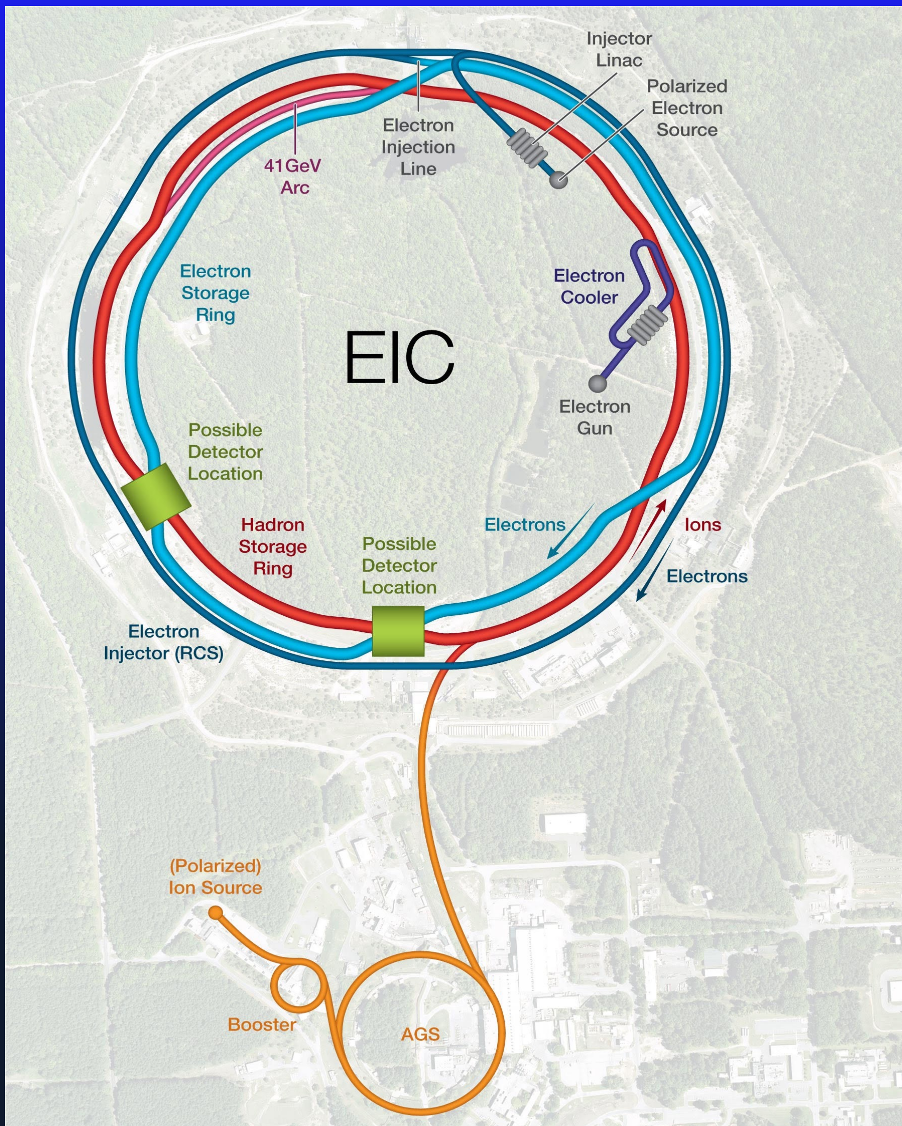
RHIC: Colliding heavy ion beams 100 GeV/A
Colliding polarized proton beams ~255 + 255 GeV





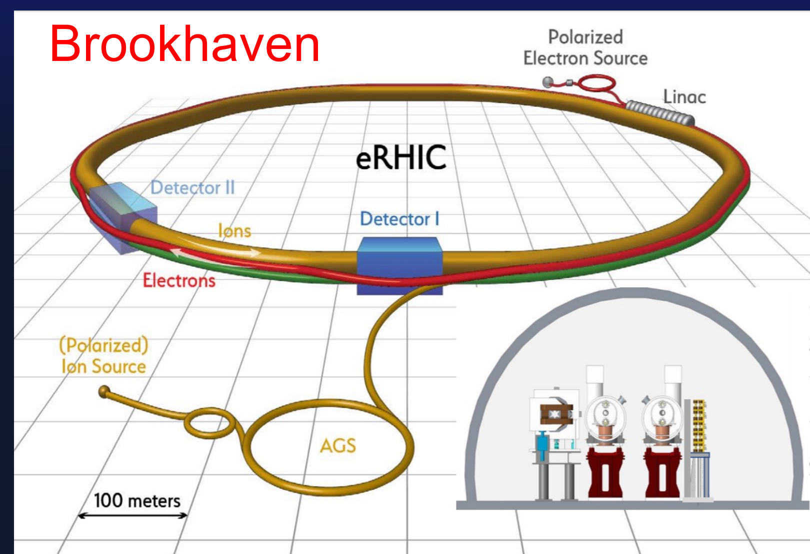
Au(197×100GeV)+Au(197×100GeV)

Brookhaven eRHIC



Add 5-18 GeV electron storage ring in present RHIC tunnel. Collide e with p to 275 GeV (vs 255 GeV now) and ions to 100 GeV/A in one RHIC ring.

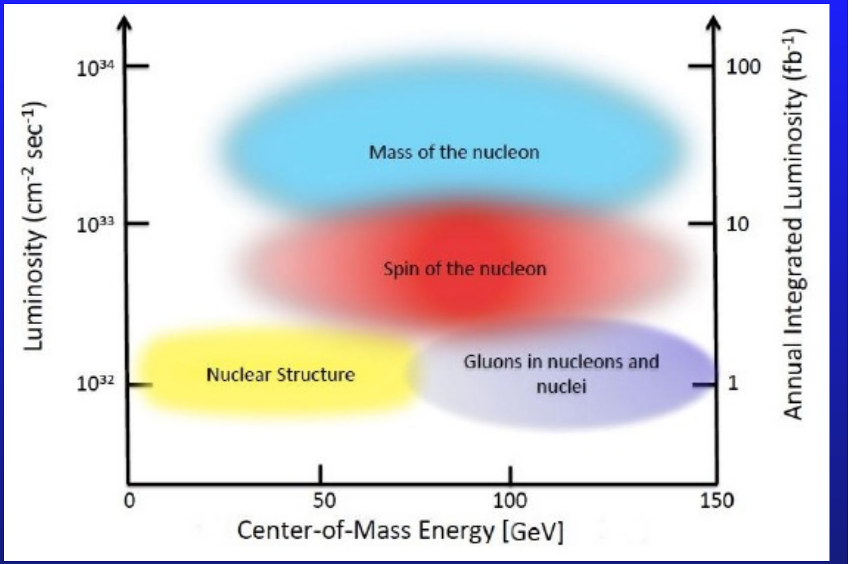
$\sqrt{s} \simeq 2\sqrt{E_e E_p}$ = 75-140 GeV



Accelerator requirements

To map quarks and gluons -- from nuclei (quark-gluon gas) to saturated gluonic matter – designing for varianble c.m. energy range from ~28 to ~120 GeV, upgradeable to 140 GeV.

Reach x down to 10^{-4}



Need ion beams from p to heavy stable nuclei: A = 1 - 208

3D imaging of gluon and sea quark distributions in nucleons and nuclei needs high luminosity:

L up to 10^{34} /cm^2 sec. (~ 10^2 -10^3 X HERA, and LHC at 2 X 10^{34})

To study correlations of gluon and sea quark distributions with spin, need polarized e-, p, (and light-ion beams) each above 70%.

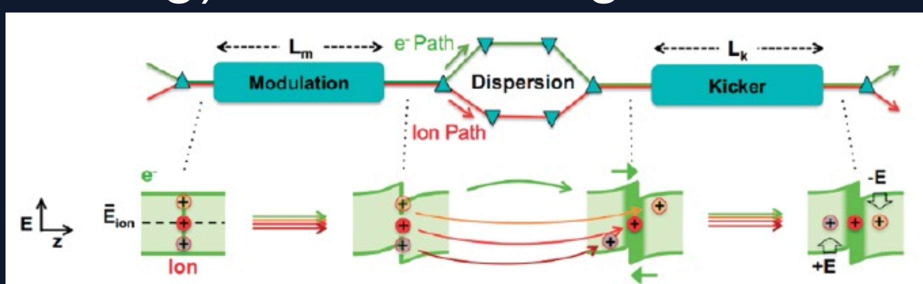
Two intersection regions (IP6 and IP8): allow for two detectors

Accelerator challenges

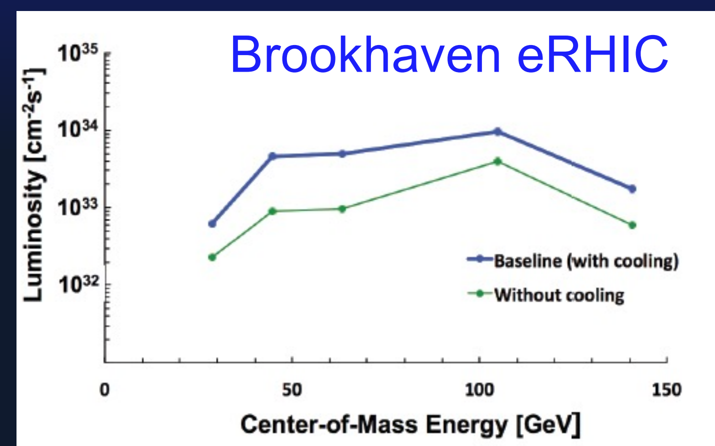
The EIC is the only large scale accelerator project in the U.S.!
Accelerator requirements beyond current technology.

High energy, spin-polarized beams colliding with high luminosity.
Polarized beams in a collider achieved only at
--HERA (polarized e^- or e^+ on unpolarized p) and at the
--Relativistic Heavy Ion Collider (RHIC - pp) with both proton beams polarized

Require strong hadron beam cooling (focusing) to achieve high luminosity

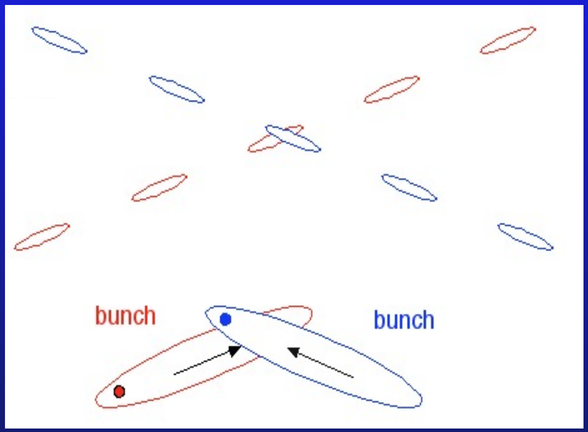


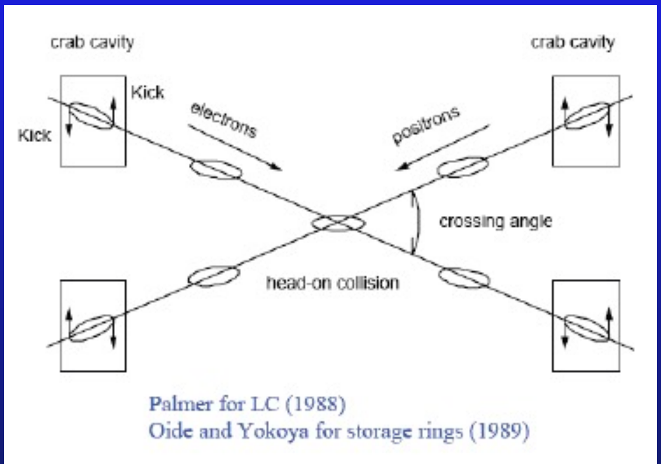
e.g., Coherent Electron Cooling: kick slow hadrons forward, fast ones backward



luminosity increase at eRHIC with hadron cooling

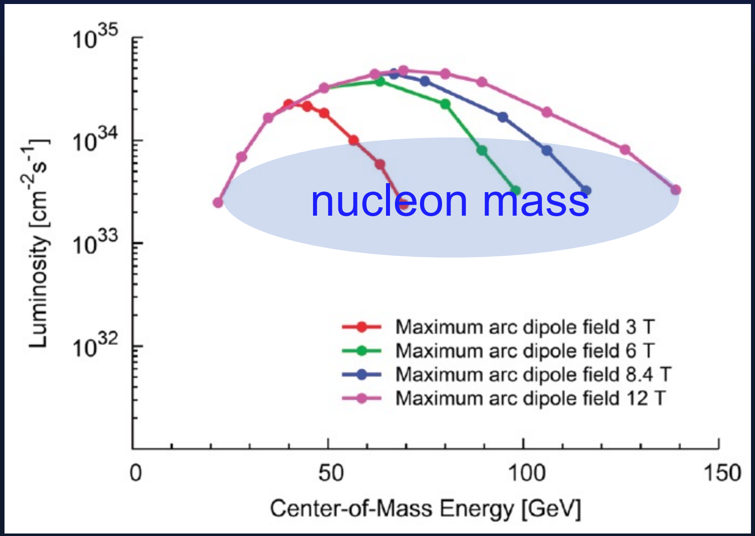
Intersection region design: Crab crossings (beam walking sideways) to maximize collisions





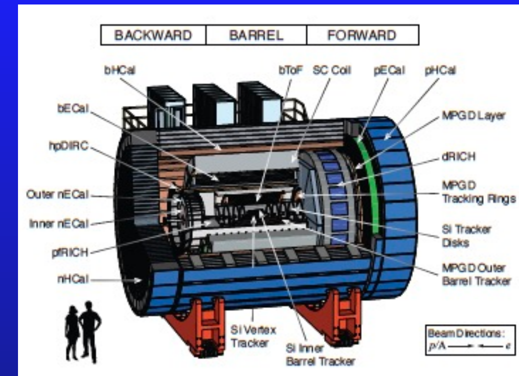
Instead of bunches crossing at an angle, turn them parallel when they collide parallel. At KEK-B (e^+e^-), but never used in hadron beams.

Magnet technology



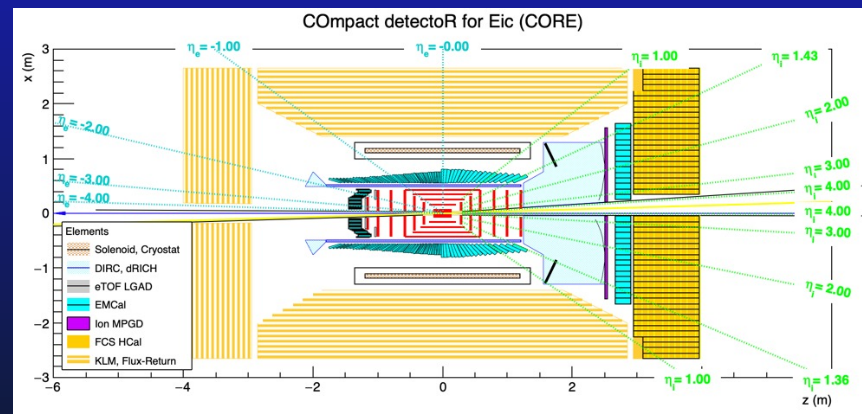
luminosity and cm ep energy vs magnetic fields (JLEIC). Red curve ~ 10 GeV e +100 GeV p

Detectors



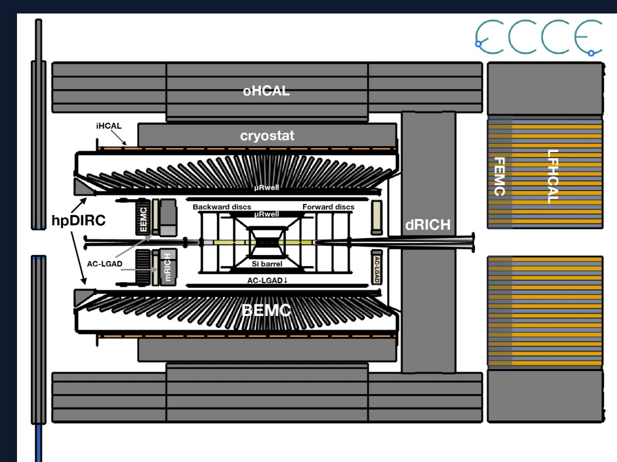
Three proposed detectors at IP6:

ATHENA (A Totally Hermetic Electron Nucleus Apparatus) *J. Adam et al. J. Instrument. 17, P10019 (2022)*



CORE - a COmpact detectoR for the EIC (arXiv:2209.00496)

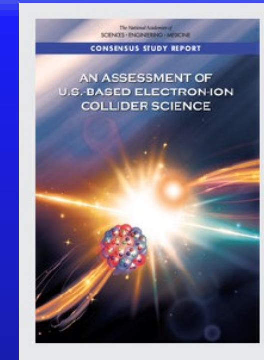
ECCE - EIC Comprehensive Chromodynamics Experiment (arXiv:2209.02580)



EPIC (Electron Proton/Ion Collisions) collaboration finalizing detector design.

Second detector working group formed

Timeline of EIC



Summer 2018: National Academy of Sciences report issued

Sept 2019: EIC enters U.S. budget

Dec. 19. 2019: Department of Energy first approval (CD-0)

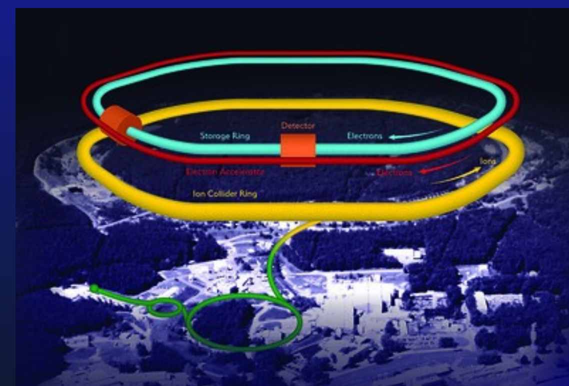
Jan. 9, 2020: Site selection -- Brookhaven

June 2021: Preliminary plan approved (CD-1)

Early 2024: Approval of final plan (CD-2)

2025: Start of construction (CD-3)

Early 2030's: Beams!! (CD-4)



Department of Energy (DOE) approval steps

PROJECT ACQUISITION PROCESS AND CRITICAL DECISIONS					
Project Planning Phase		Project Execution Phase			Mission
Preconceptual Planning	Conceptual Design	Preliminary Design	Final Design	Construction	Operations
CD-0 Approve Mission Need	CD-1 Approve Preliminary Baseline Range	CD-2 Approve Performance Baseline	CD-3 Approve Start of Construction	CD-4 Approve Start of Operations or Project Closeout	

EIC Users Group

EIC Users Group now has over 1200 physicists from some 250 labs worldwide.

http://www.eicug.org/

Link to EIC Users Group White Paper (Nov. 2022):

https://www.dropbox.com/s/yk6pvhoxd9irdz9/LRP_EIC_White_Paper_Draft.pdf?dl=0