



Recent Results of GPD Measurements at COMPASS and JLab Hall A

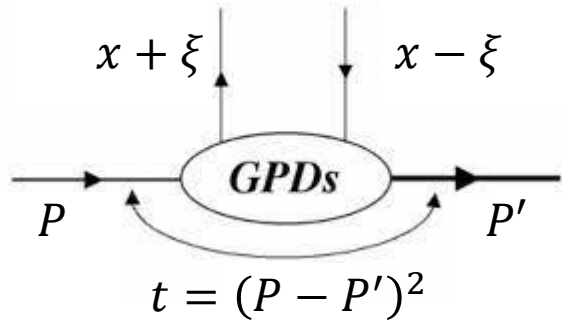
NCU workshop on EIC physics and detectors

November 09, 2022

Po-Ju Lin

Institute of Physics, Academia Sinica

Generalized Parton Distributions (GPDs)



➤ At fixed Q^2 , the GPDs depend on the following variables:

x : average longitudinal momentum fraction

ξ : longitudinal momentum difference

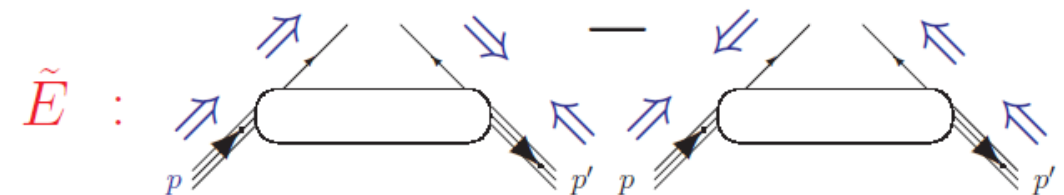
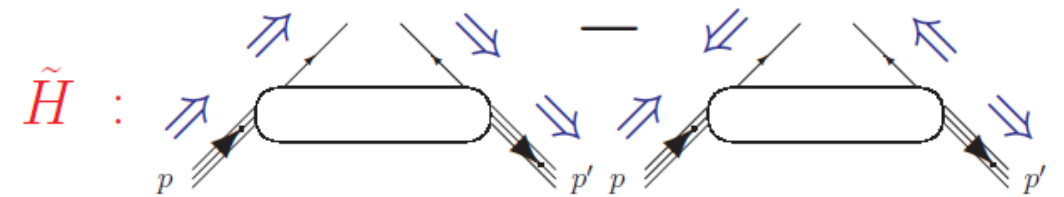
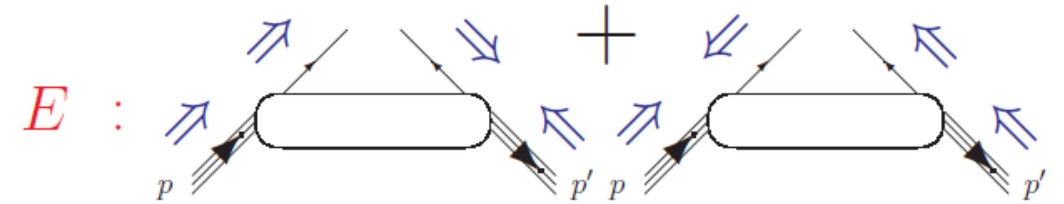
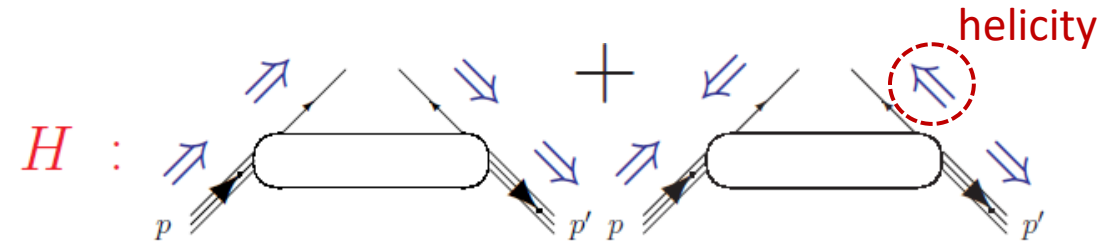
t : four momentum transfer

(correlated to b_{\perp} via Fourier transform)

➤ A total of 8 GPDs for a specific parton

4 Chiral-even (parton helicity unchanged): $H, E, \tilde{H}, \tilde{E}$

4 Chiral-odd (parton helicity changed): $H_T, E_T, \tilde{H}_T, \tilde{E}_T$



Generalized Parton Distributions (GPDs)

Dirac FF

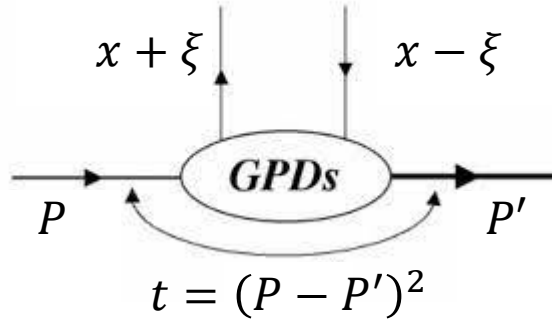
$$\int H_f(x, \xi, t) dx = F_1^f(t)$$

First Moment

$$\int E_f(x, \xi, t) dx = F_2^f(t)$$

Form Factors

Pauli FF



$\xi = t = 0$
PDFs

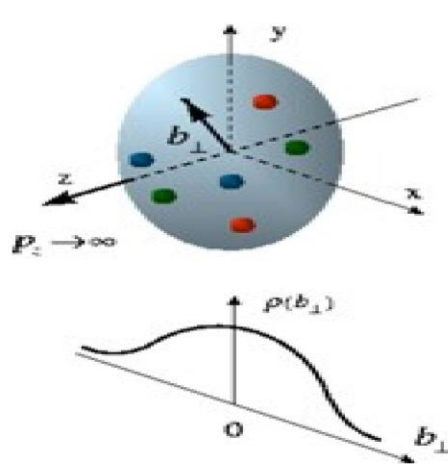
$$H_f(x, 0, 0) = q_f(x)$$

Unpolarized

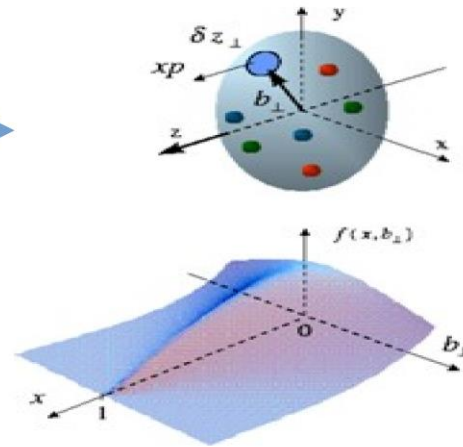
$$\tilde{H}_f(x, 0, 0) = \Delta q_f(x)$$

Polarized

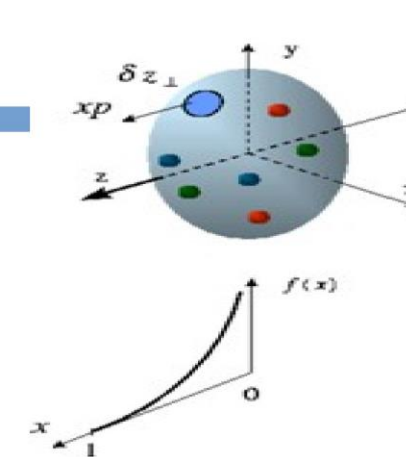
➤ GPDs embody both PDFs and FFs



- Form Factors (FFs)
- ✓ Spatial distribution
- ✗ Momentum distribution

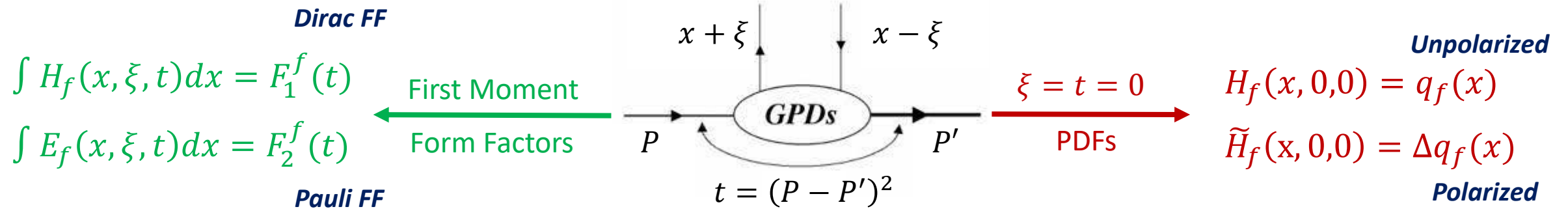


- Generalized Parton Distributions (GPDs)
- ✓ Spatial distribution
- ✓ Longitudinal momentum distribution



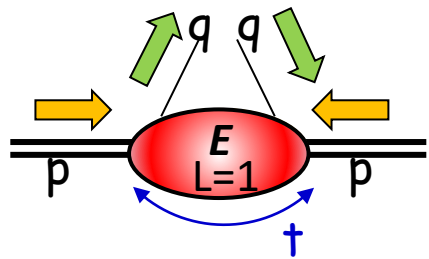
- Parton Distribution Functions (PDFs)
- ✓ Longitudinal momentum distribution
- ✗ Spatial distribution

Generalized Parton Distributions (GPDs)



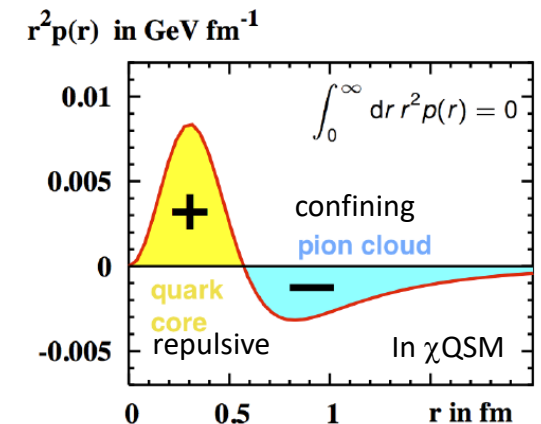
Provides information on the interesting properties of the nucleon.

- Mapping the transverse plane distribution of parton
- Pressure distribution inside nucleon
- Angular momentum of parton



$$J_q = \frac{1}{2} \int_{-1}^1 dx x [H^q(x, \xi, 0) + E^q(x, \xi, 0)]$$

Ji's Sum Rule

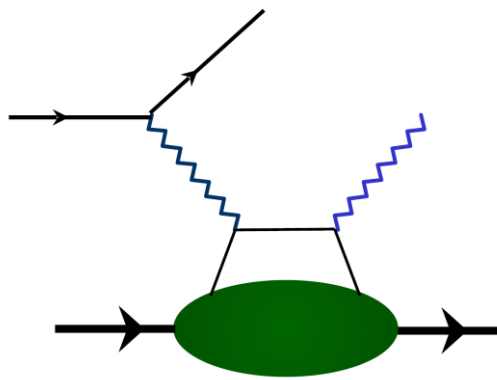


M. Polyakov, P. Schweitzer, Int.J.Mod.Phys. A33 (2018)

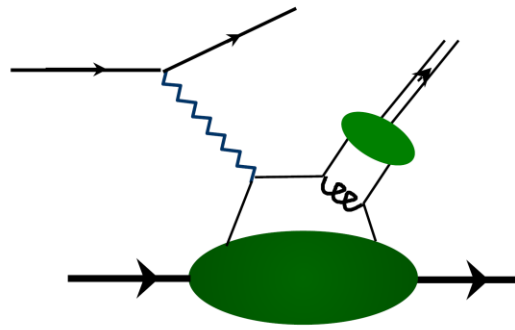
Exclusive Process

- Use **exclusive processes**, where all final state particles are “identified”, to access the multi-variable dependence of GPDs, and constrain the GPD parameterization with measurements in various phase space.
- Processes:

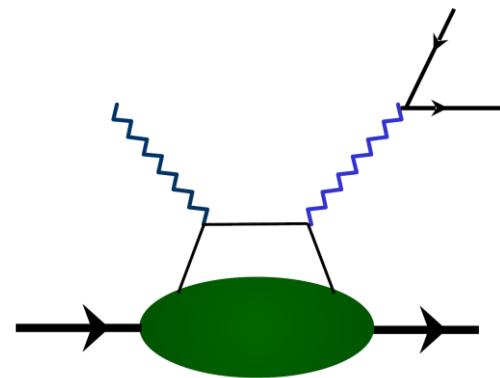
- Deeply Virtual Compton Scattering (DVCS)
- Deeply Virtual Meson Production (DVMP)
- Time-like Compton Scattering (TCS)
- Double DVCS (DDVCS)
- ...



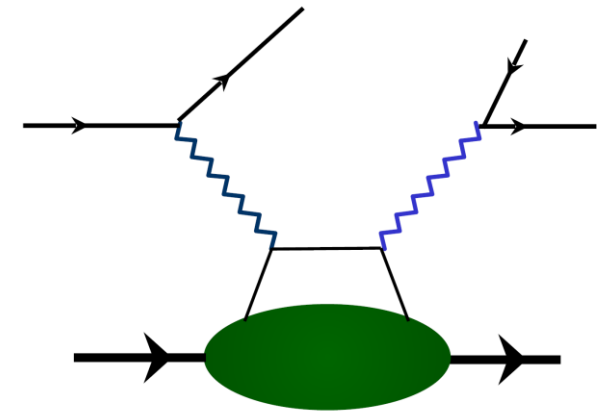
DVCS



DVMP

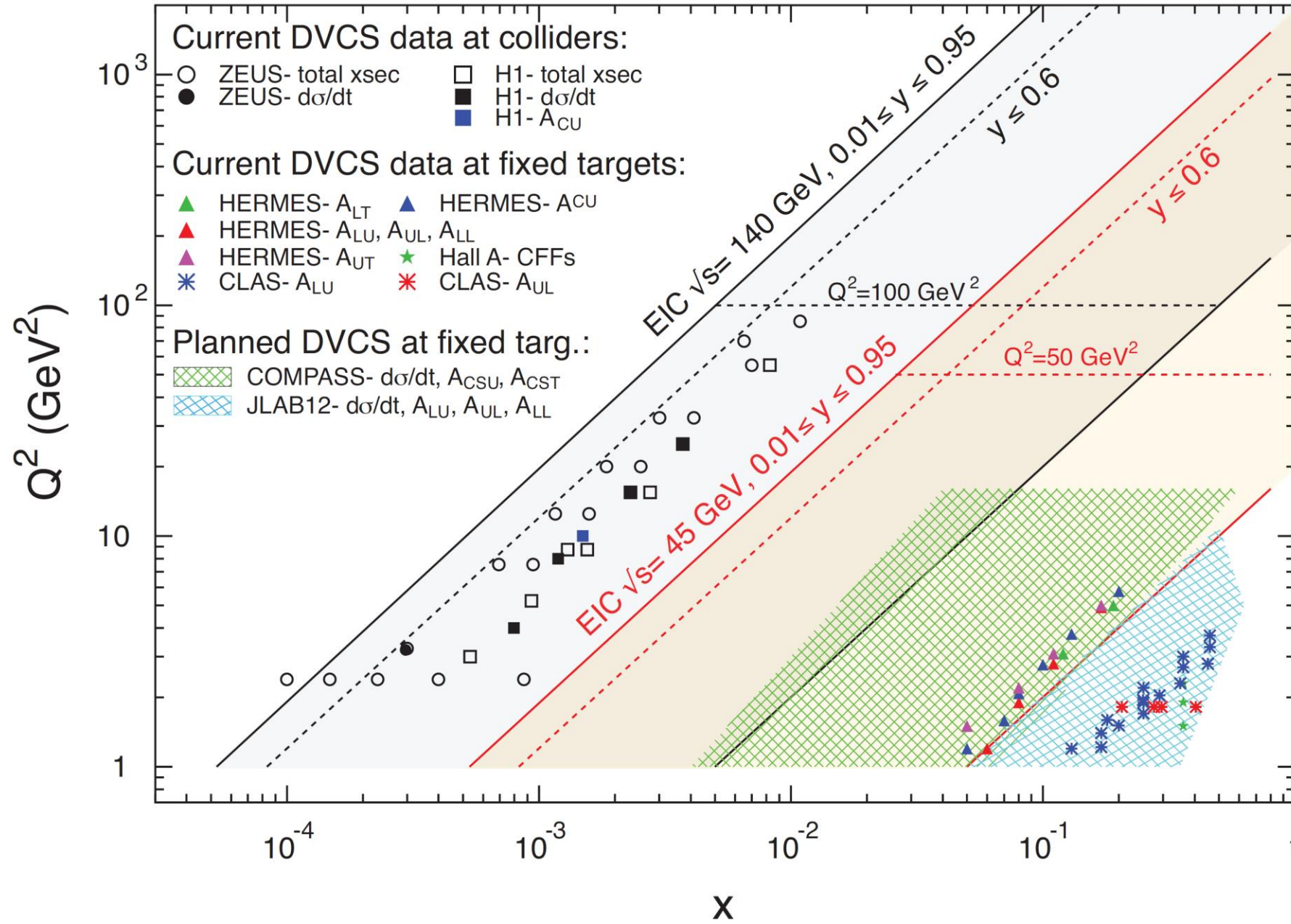


TCS



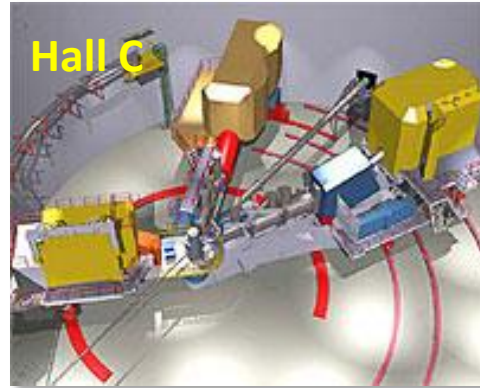
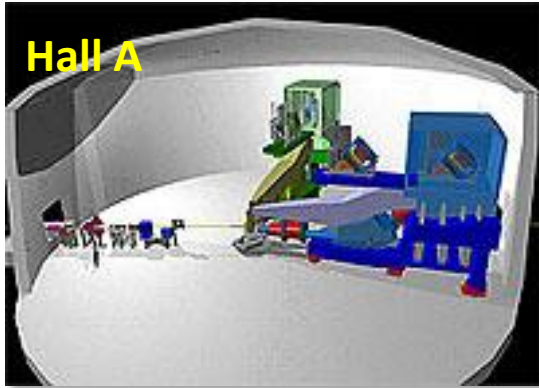
DDVCS

Landscape – Global Programs of DVCS

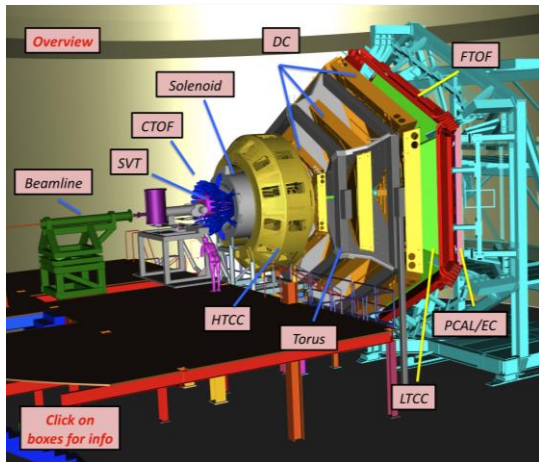


JLab & COMPASS

- **Jlab: Hall A, C, CLAS** High Luminosity Polar. **6 & 12 GeV e-** Long., (Trans.) polarised p & nuclear target
Missing mass technique (A,C) and complete detection (CLAS)

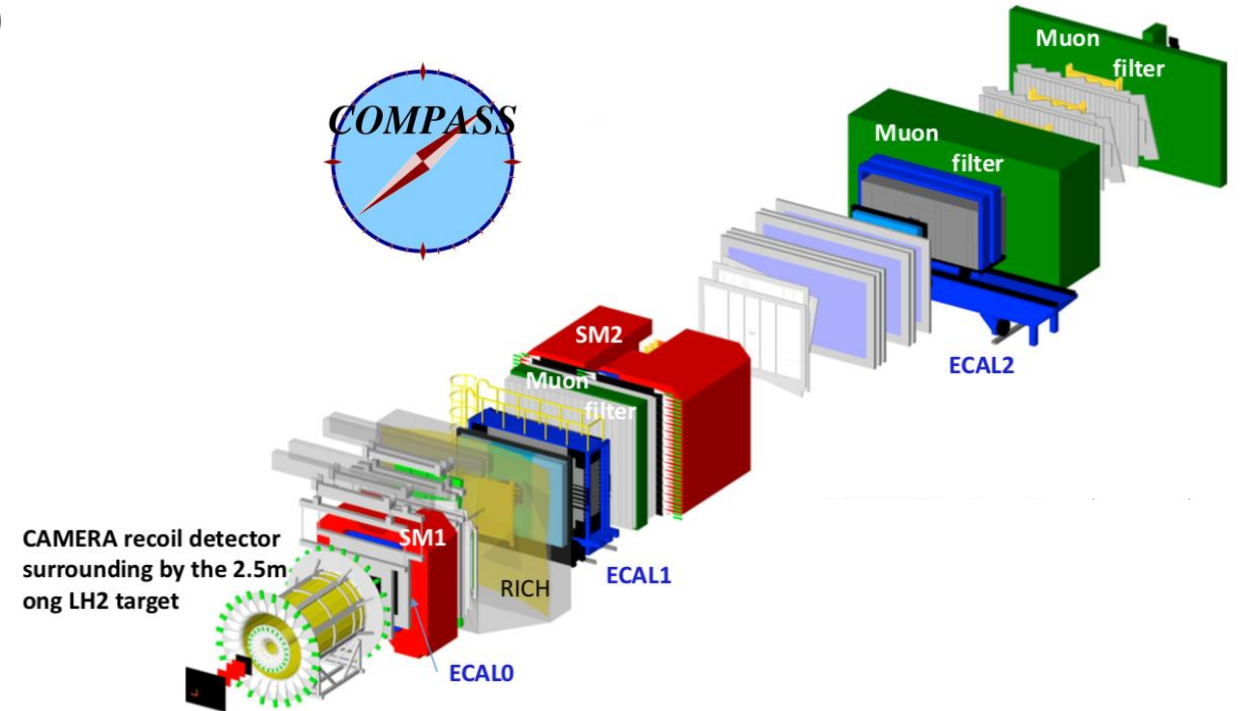


- High luminosity, limited kinematic coverage
→ Test the validity of theoretical formalism

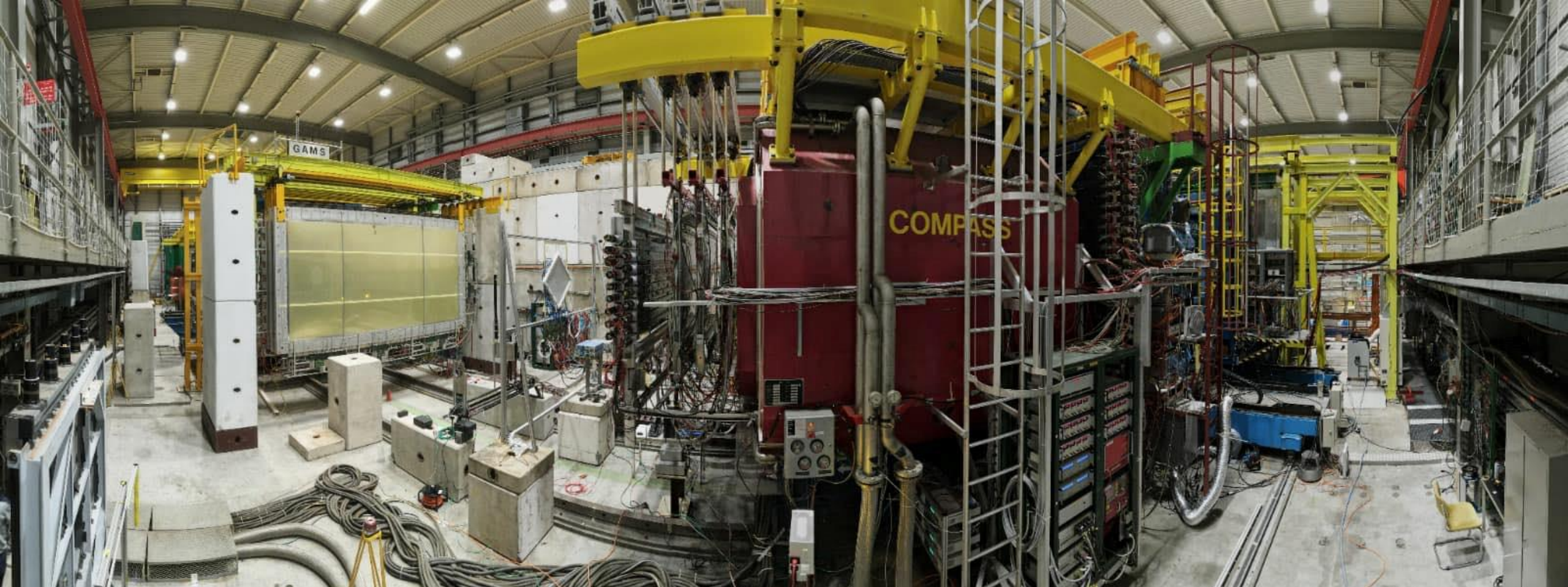


Hall B – CLAS 12

- Lower luminosity, wide kinematic coverage
→ Map the GPDs



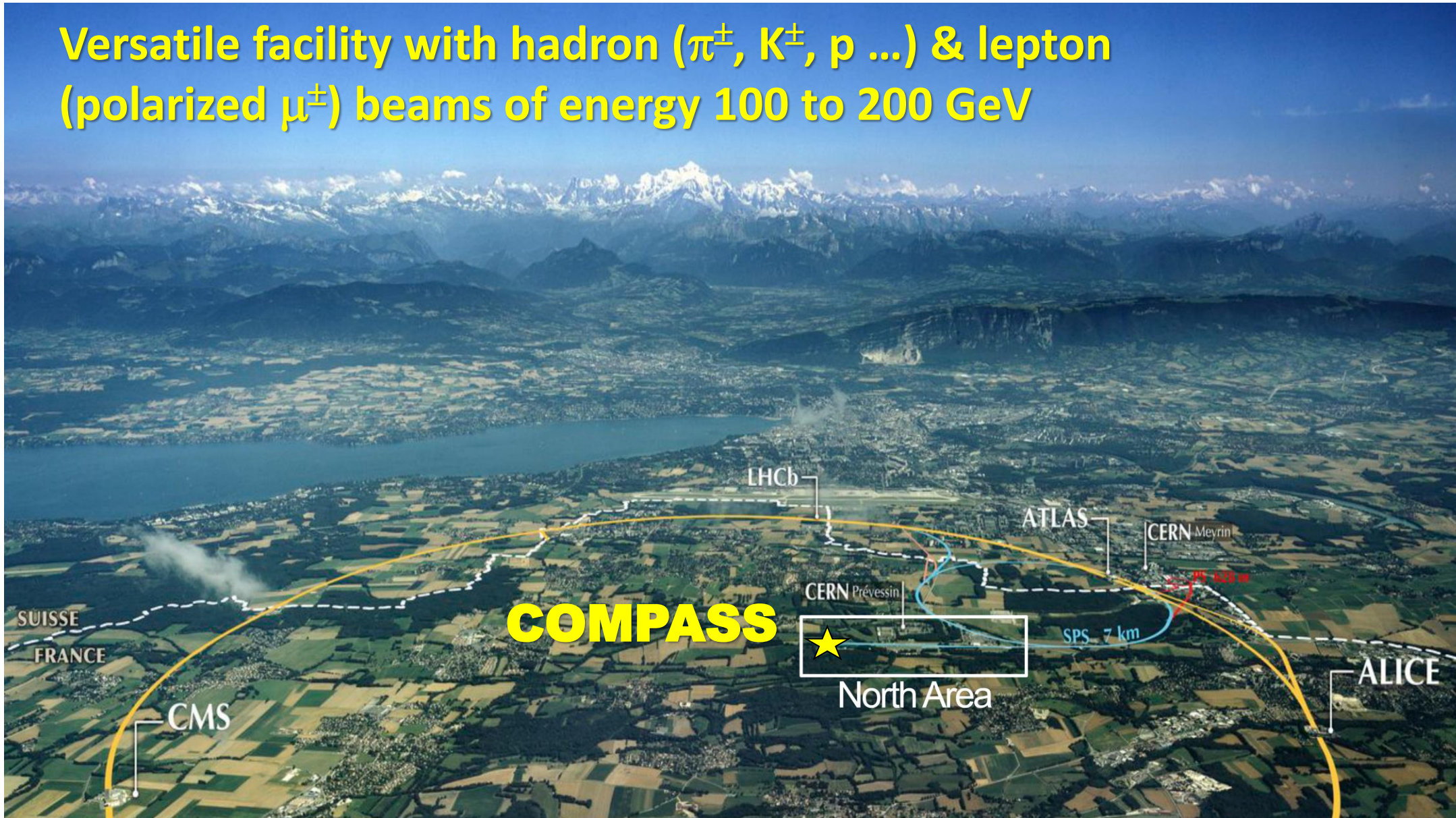
- **COMPASS @ CERN:** Polarised **160 GeV μ^+/μ^-** p target, (Trans.) polarised *target with recoil p detection*



Deeply Virtual Compton Scattering @ COMPASS

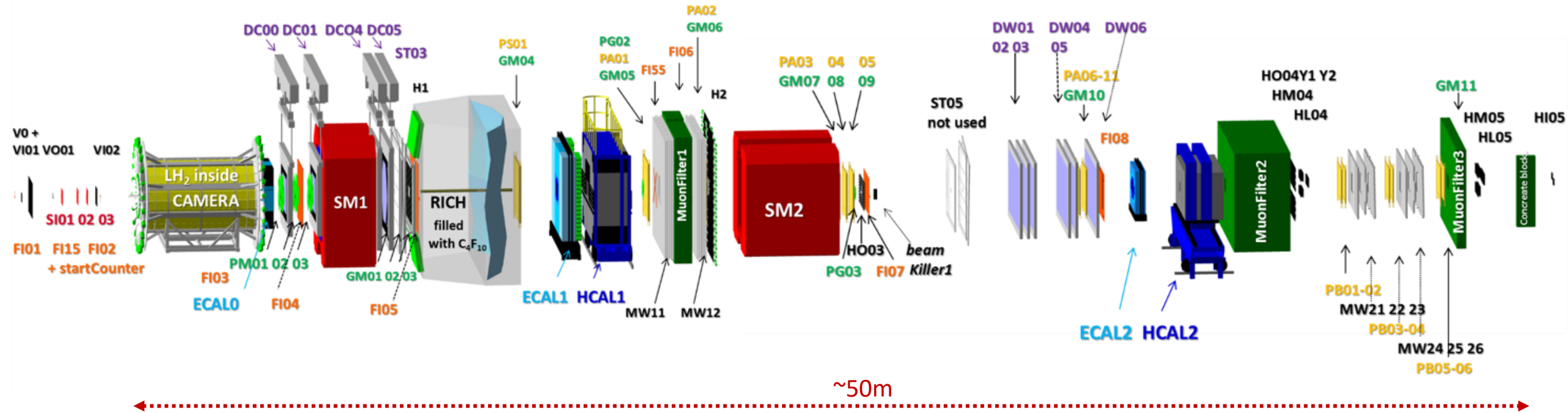
COMPASS Experiment

Versatile facility with hadron (π^\pm , K^\pm , p ...) & lepton (polarized μ^\pm) beams of energy 100 to 200 GeV



COmmun
MUon and
PRoton
APparatus for
STructure and
SPectroscopy

COMPASS Setup for Hard Exclusive Measurements



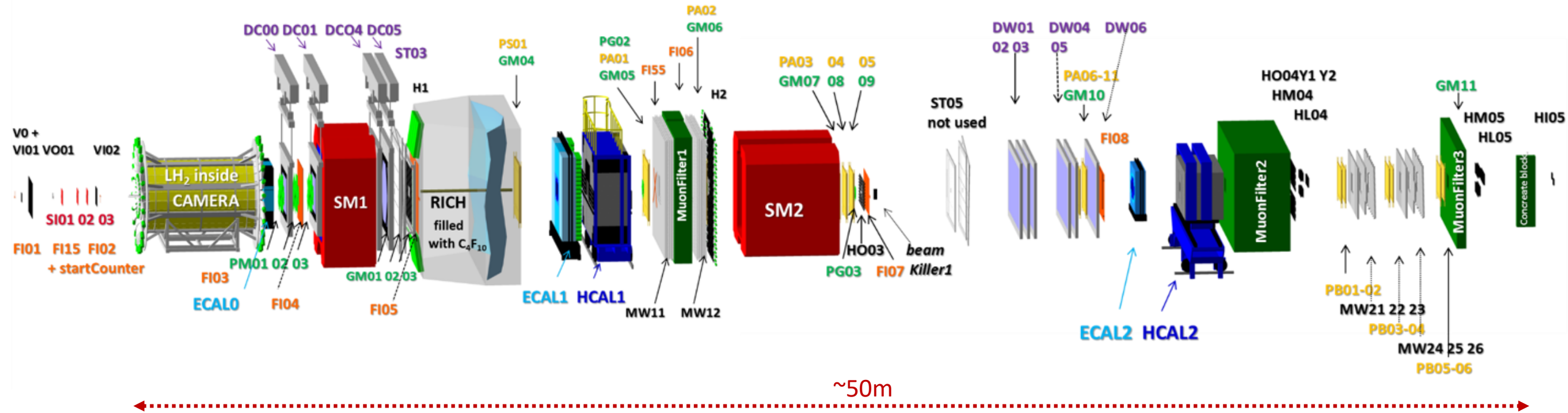
Muon Beams

- μ^+ & μ^- with opposite polarisation
- About $\pm 80\%$ polarisation
- Momentum: 160 GeV/c

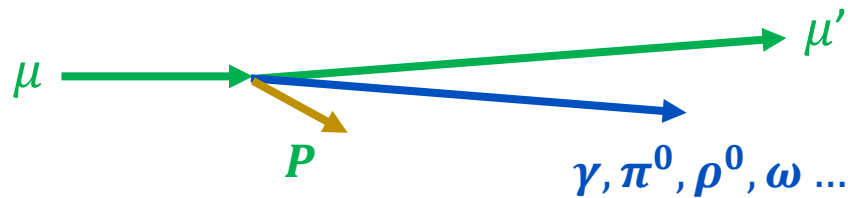
Two-stage, large angle, and wide momentum range spectrometer. PID including hadron absorbers, RICH, HCALs, ECALs, and muon filters.

❖ NIM A 577 (2007) & NIM A 779 (2015) 69

COMPASS Setup for Hard Exclusive Measurements



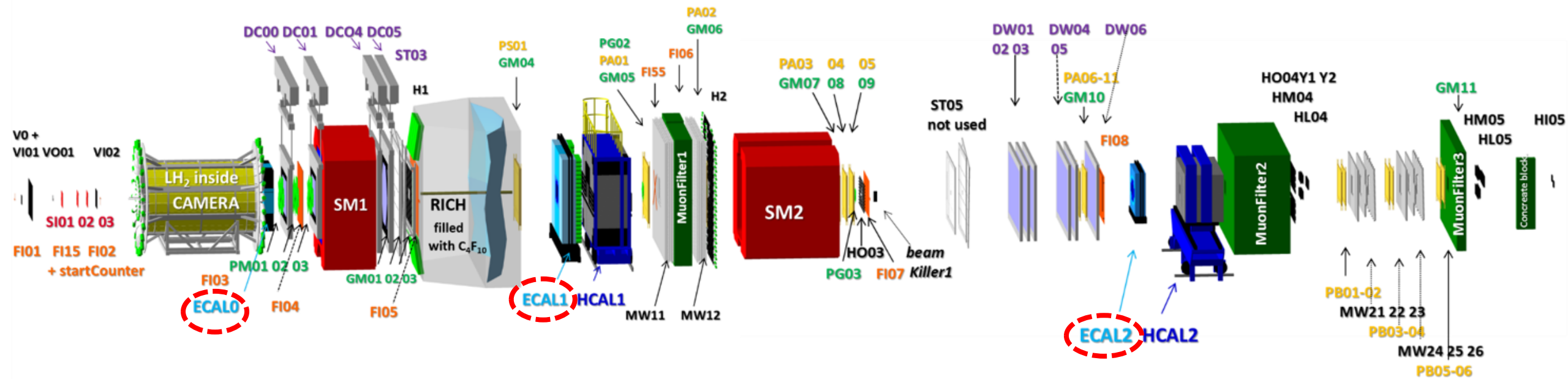
Exclusive Muoproduction



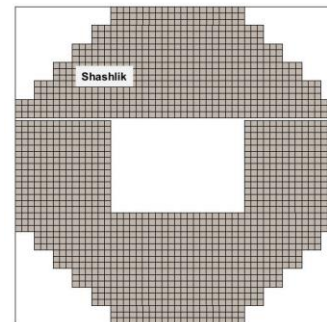
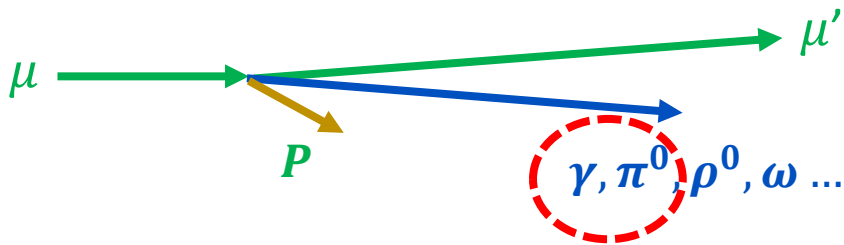
Two-stage, large angle, and wide momentum range spectrometer. PID including hadron absorbers, RICH, HCALs, ECALs, and muon filters.

❖ NIM A 577 (2007) & NIM A 779 (2015) 69

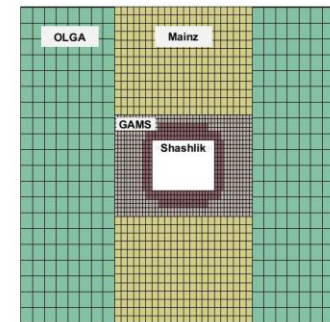
COMPASS Setup for Hard Exclusive Measurements



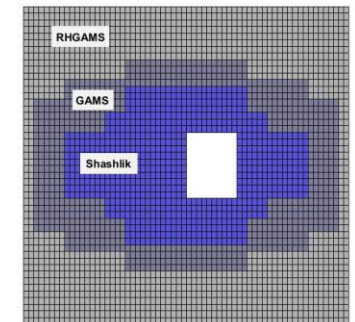
Exclusive Muoproduction



ECAL 0
2×2 m²
1746 cells

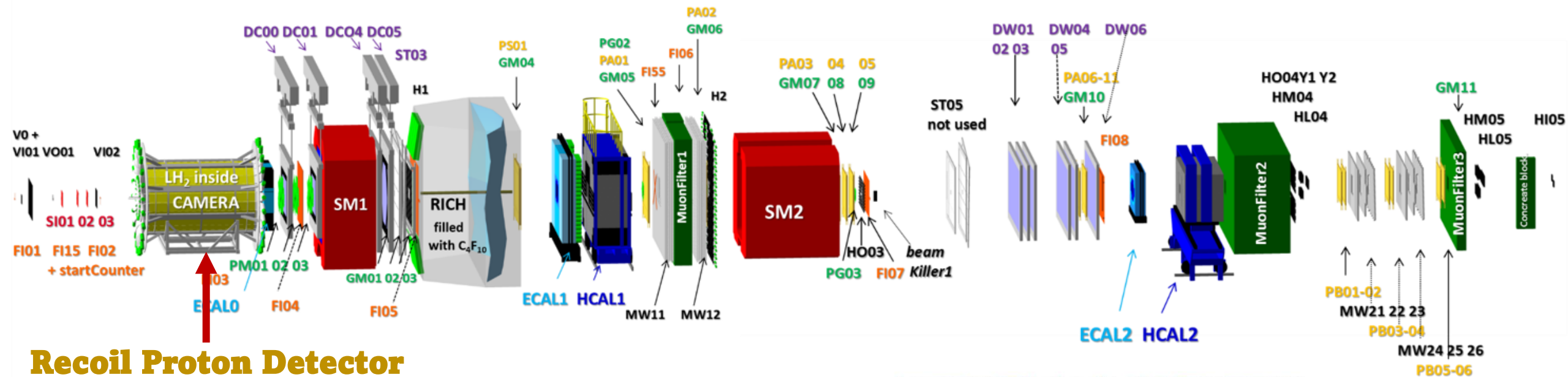


ECAL 1
4×3 m²
1708 cells



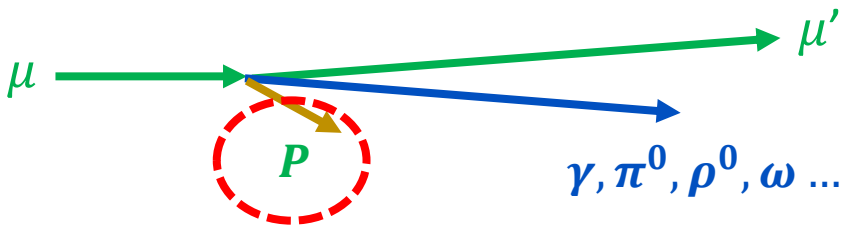
ECAL 2
2.5×2 m²
2972 cells

COMPASS Setup for Hard Exclusive Measurements

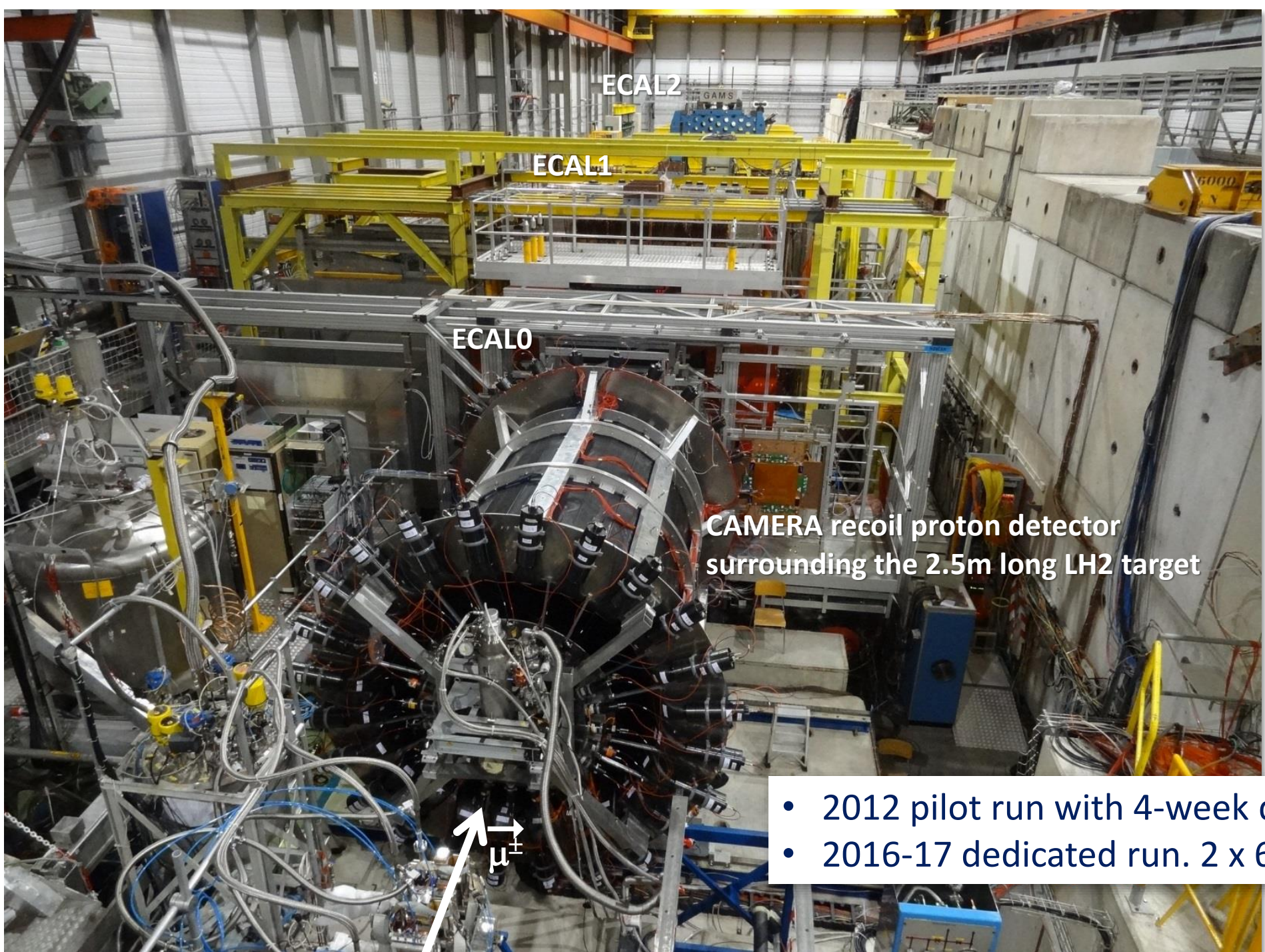


Recoil Proton Detector (CAMERA)

Exclusive Muoproduction



transverse picture of CAMERA



ECAL2

ECAL1

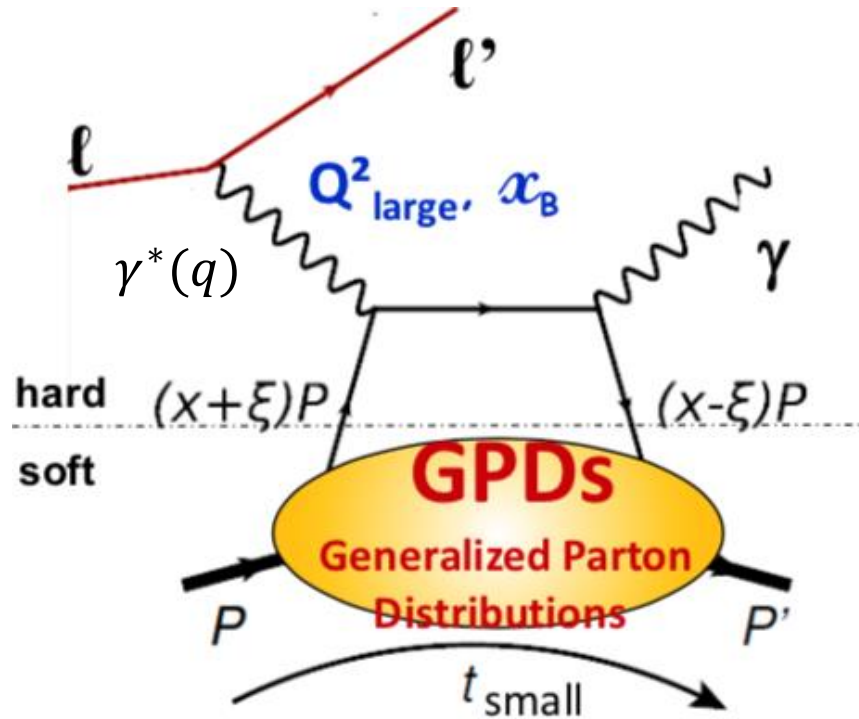
ECAL0

CAMERA recoil proton detector
surrounding the 2.5m long LH2 target

μ^\pm

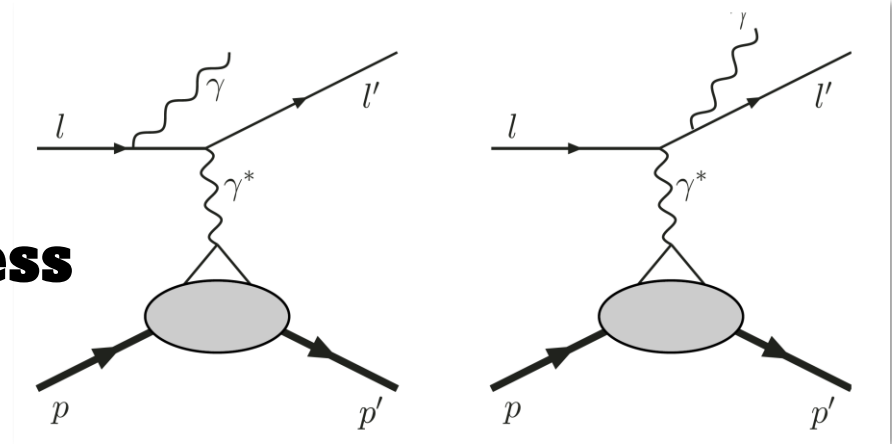
- 2012 pilot run with 4-week data taking
- 2016-17 dedicated run. 2 x 6 months.

Deeply Virtual Compton Scattering



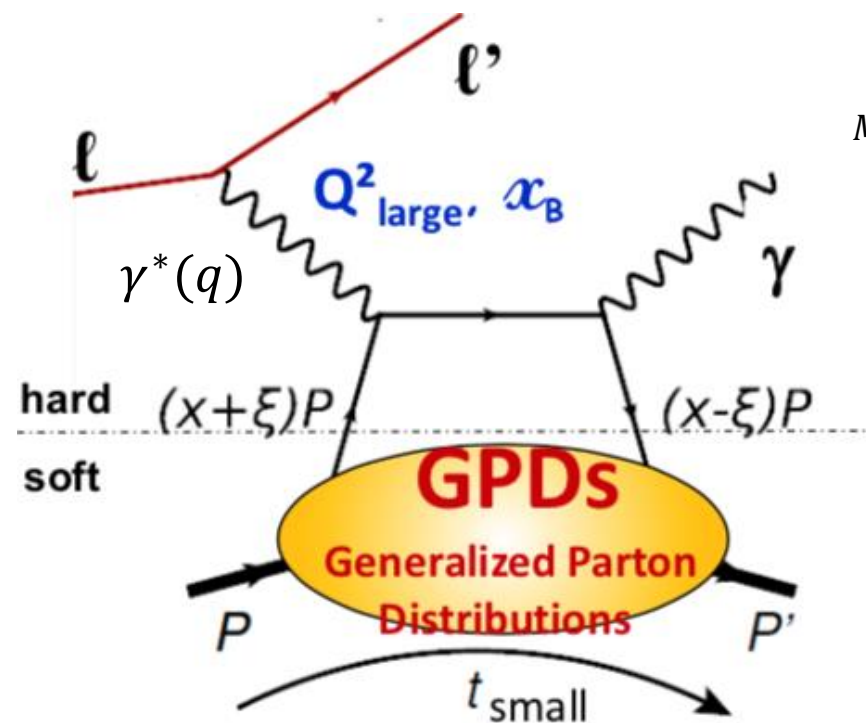
$$\text{DVCS: } l + p \rightarrow l' + p' + \gamma$$

**BH
Process**



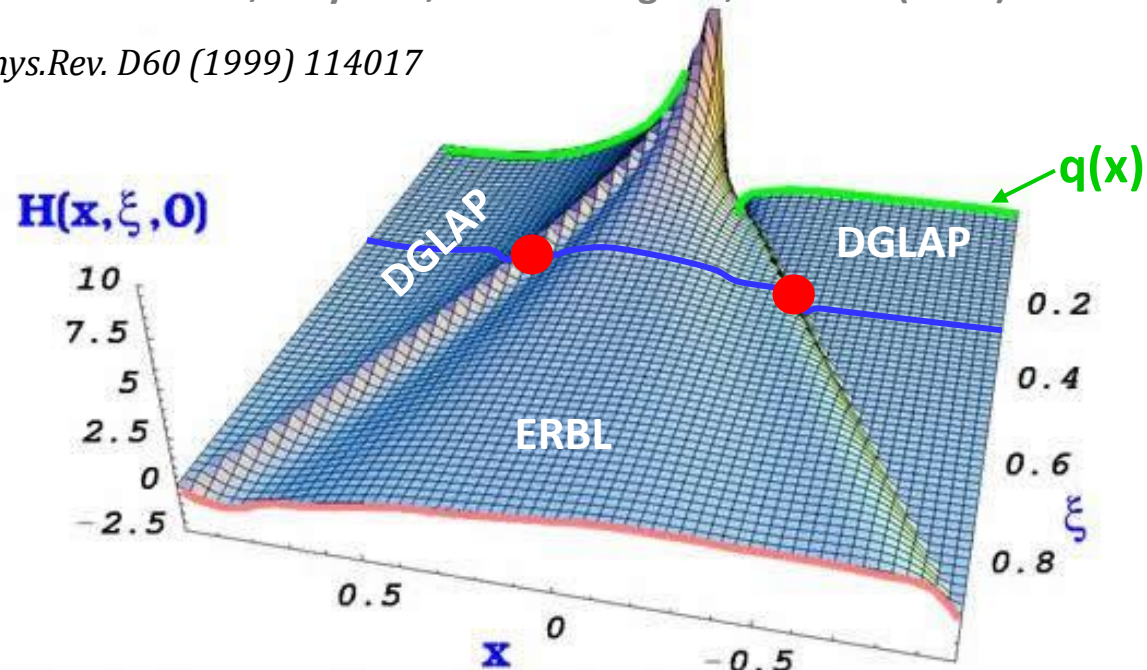
- DVCS is regarded as the golden channel and gives access to four chiral-even GPDs $H, \tilde{H}, E, \tilde{E}(x, \xi, t)$. Its interference with the well-understood Bethe-Heitler process gives access to more info.
- With LH_2 target and small x_B coverage \rightarrow focuses on **H** at COMPASS

Compton Form Factors (CFFs)



From Goeke, Polyakov, Vanderhaeghen, PNPP47 (2001)

M. Polyakov, C. Weiss, Phys.Rev. D60 (1999) 114017



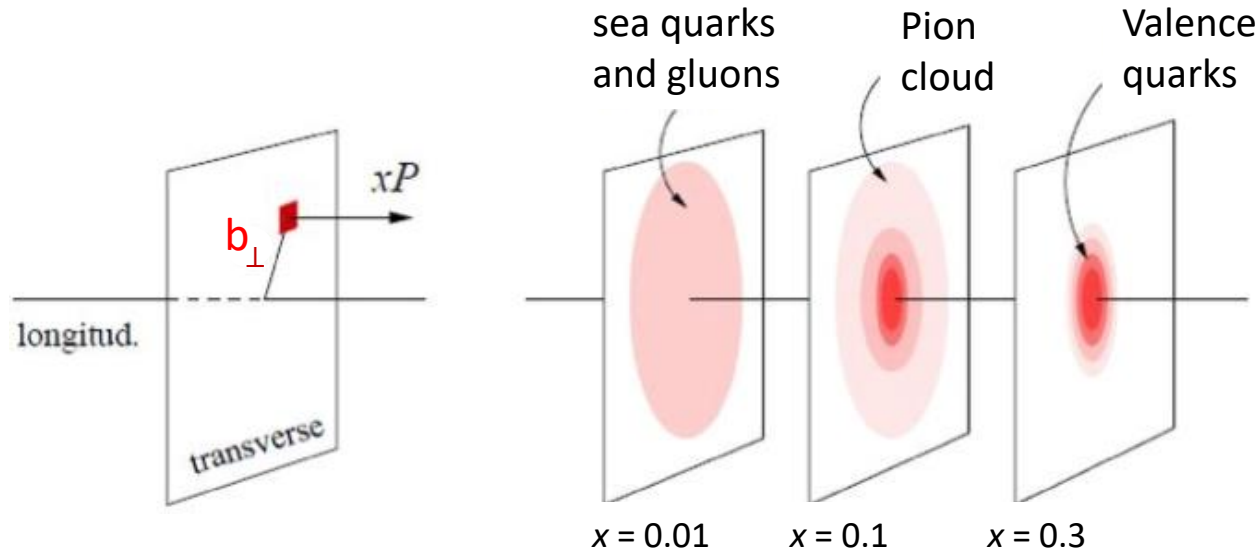
$$\overset{\text{CFF}}{\mathcal{H}(\xi, t)} \overset{\text{GPD}}{=} \int_{-1}^{+1} dx \frac{\mathbf{H}(x, \xi, t)}{x - \xi + i\epsilon} + \dots = \mathcal{P} \int_{-1}^{+1} dx \frac{\mathbf{H}(x, \xi, t)}{x - \xi} - i\pi \mathbf{H}(x = \pm \xi, \xi, t) + \dots$$

$$\text{Re } \mathcal{H}(\xi, t) = \mathcal{P} \int dx \frac{\text{Im } \mathcal{H}(x, t)}{x - \xi} + \Delta(t)$$

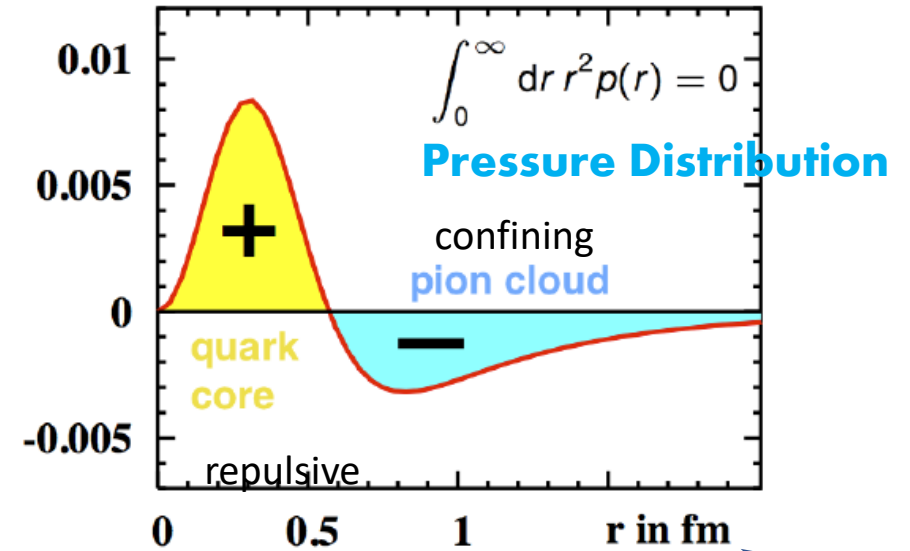
Transverse Imaging and Pressure Distribution

M. Polyakov, P. Schweitzer, Int.J.Mod.Phys. A33 (2018)

Mapping in the transverse plane



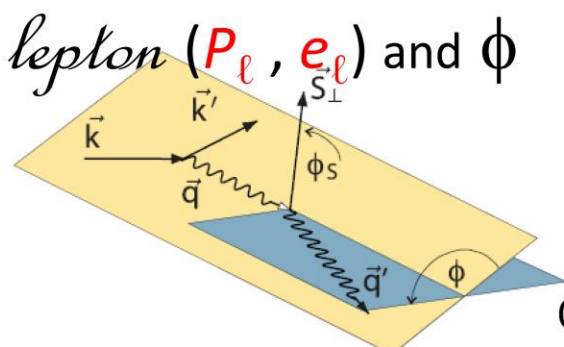
$r^2 p(r)$ in GeV fm^{-1}



$$\begin{aligned}
 & \text{CFP} \rightarrow \mathcal{H}(\xi, t) = \int_{-1}^{+1} dx \frac{\mathbf{H}(x, \xi, t)}{x - \xi + i\epsilon} + \dots \\
 & \text{GPD} \rightarrow \mathcal{H}(\xi, t) = \mathcal{P} \int_{-1}^{+1} dx \frac{\mathbf{H}(x, \xi, t)}{x - \xi} - i\pi \mathbf{H}(x = \pm \xi, \xi, t) + \dots \\
 & \text{REAL part} \rightarrow \text{Re } \mathcal{H}(\xi, t) = \mathcal{P} \int dx \frac{\text{Im } \mathcal{H}(x, t)}{x - \xi} + \Delta(t) \\
 & \text{Imaginary part} \rightarrow \text{Im } \mathcal{H}(\xi, t) = \dots \\
 & \text{D-term} \rightarrow d_1(t)
 \end{aligned}$$

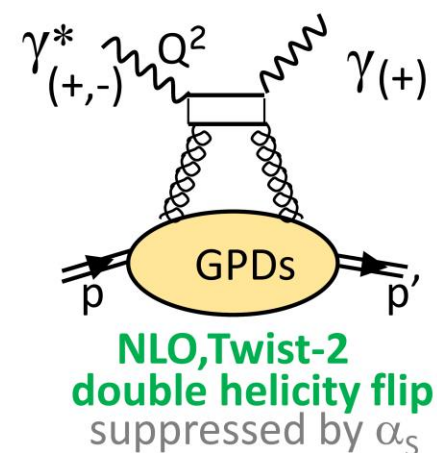
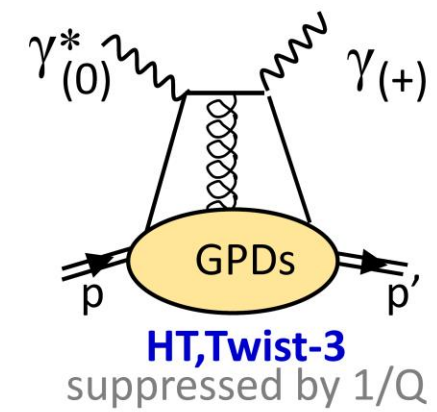
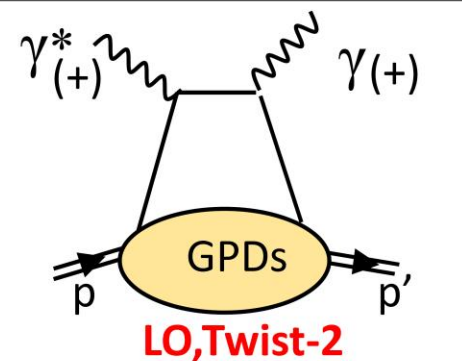
\leftarrow FT of $H(x, \xi=0, t)$

Azimuthal Dependence of BH & DVCS



$$d\sigma = \left| \left[\text{BH} + \text{DVCS} \right] \right|^2 + \text{Interference Term}$$

$$\frac{d^4\sigma(\ell p \rightarrow \ell p \gamma)}{dx_B dQ^2 d|t| d\phi} = \underbrace{d\sigma^{BH}}_{\text{Well known}} + \left(d\sigma_{unpol}^{DVCS} + P_\ell d\sigma_{pol}^{DVCS} \right) + (e_\ell \text{Re } I + e_\ell P_\ell \text{Im } I)$$



With unpolarized target:

Belitsky, Müller, Kirner, NPB629 (2002)

$$d\sigma^{BH} \propto c_0^{BH} + c_1^{BH} \cos \phi + c_2^{BH} \cos 2\phi$$

$$d\sigma_{unpol}^{DVCS} \propto c_0^{DVCS} + c_1^{DVCS} \cos \phi + c_2^{DVCS} \cos 2\phi$$

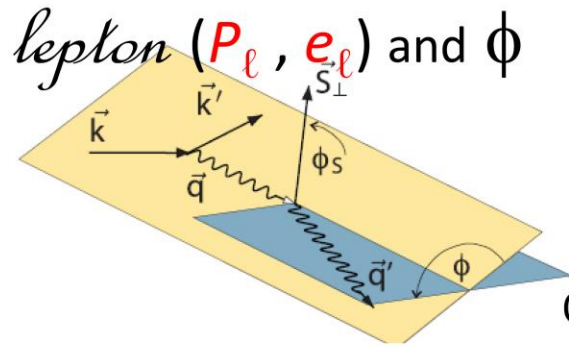
$$d\sigma_{pol}^{DVCS} \propto s_1^{DVCS} \sin \phi$$

$$\text{Re } I \propto c_0^I + c_1^I \cos \phi + c_2^I \cos 2\phi + c_3^I \cos 3\phi$$

$$\text{Im } I \propto s_1^I \sin \phi + s_2^I \sin 2\phi$$

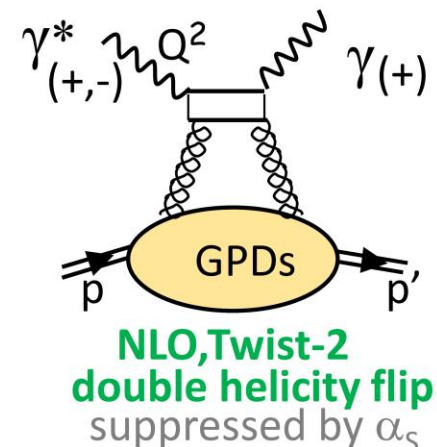
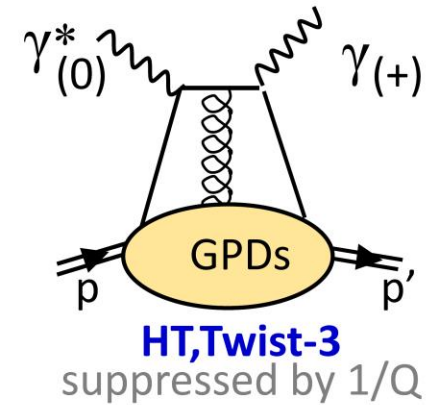
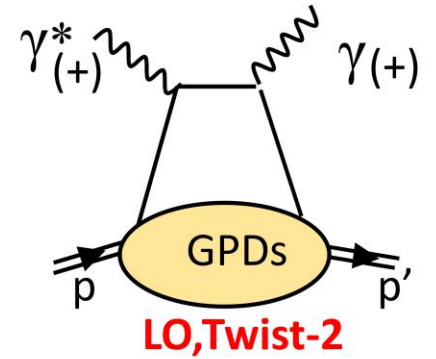
➤ Each term involves different form factors.

Azimuthal Dependence of BH & DVCS



$$d\sigma = \left| \left[\text{BH} + \text{DVCS} \right] \right|^2 + \text{Interference Term}$$

$$\frac{d^4\sigma(\ell p \rightarrow \ell p \gamma)}{dx_B dQ^2 d|t| d\phi} = \underbrace{d\sigma^{BH}}_{\text{Well known}} + \left(d\sigma_{\text{unpol}}^{DVCS} + P_\ell d\sigma_{\text{pol}}^{DVCS} \right) + (e_\ell \text{Re } I + e_\ell P_\ell \text{Im } I)$$



With unpolarized target:

Belitsky, Müller, Kirner, NPB629 (2002)

① beam spin diff

$$\Delta \equiv d\sigma^{\leftarrow} - d\sigma^{\rightarrow}$$

$$d\sigma^{BH} \propto c_0^{BH} + c_1^{BH} \cos \phi + c_2^{BH} \cos 2\phi$$

$$d\sigma_{\text{unpol}}^{DVCS} \propto c_0^{DVCS} + c_1^{DVCS} \cos \phi + c_2^{DVCS} \cos 2\phi$$

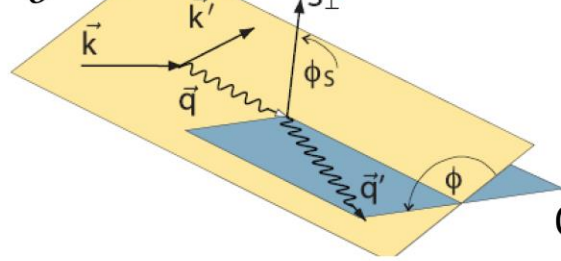
$$d\sigma_{\text{pol}}^{DVCS} \propto s_1^{DVCS} \sin \phi$$

$$\text{Re } I \propto c_0^I + c_1^I \cos \phi + c_2^I \cos 2\phi + c_3^I \cos 3\phi$$

$$\text{Im } I \propto s_1^I \sin \phi + s_2^I \sin 2\phi$$

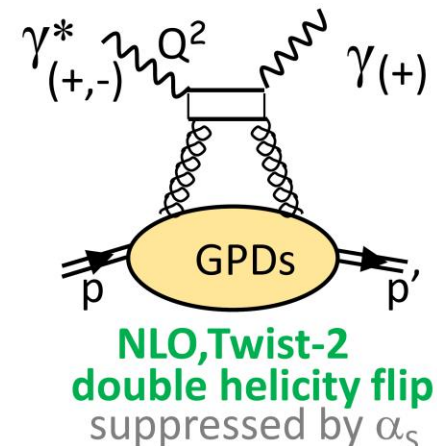
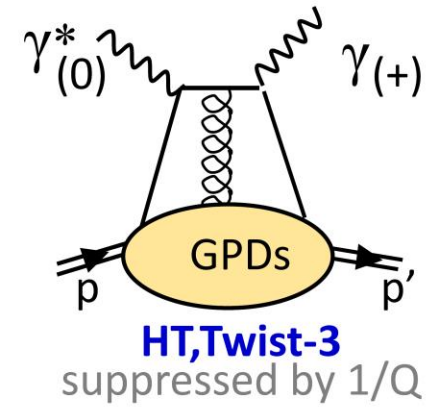
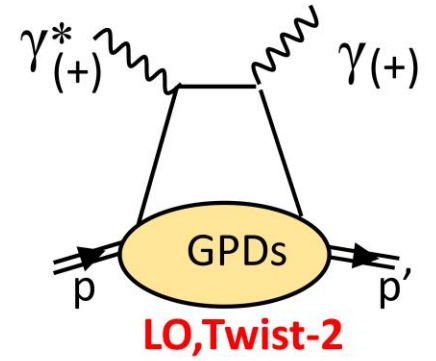
Azimuthal Dependence of BH & DVCS

lepton (P_ℓ, e_ℓ) and ϕ



$$d\sigma = \left| \underbrace{\text{BH}}_{\text{Well known}} + \underbrace{\text{DVCS}}_{\text{Well known}} \right|^2 + \text{Interference Term}$$

$$\frac{d^4\sigma(\ell p \rightarrow \ell p \gamma)}{dx_B dQ^2 d|t| d\phi} = d\sigma^{BH} + \left(d\sigma_{unpol}^{DVCS} + P_\ell d\sigma_{pol}^{DVCS} \right) + (e_\ell \text{Re } I + e_\ell P_\ell \text{Im } I)$$



With unpolarized target:

Belitsky, Müller, Kirner, NPB629 (2002)

① beam spin diff

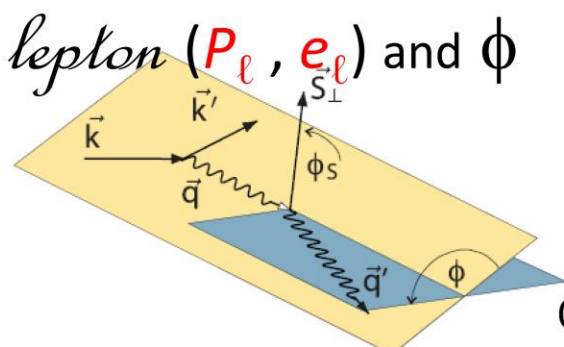
$$\Delta \equiv d\sigma^{\leftarrow} - d\sigma^{\rightarrow}$$

② beam spin sum

$$\Sigma \equiv d\sigma^{\leftarrow} + d\sigma^{\rightarrow}$$

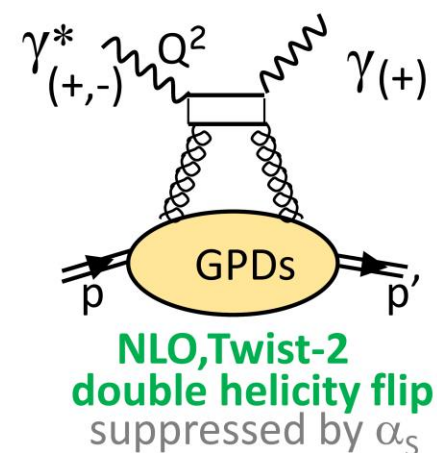
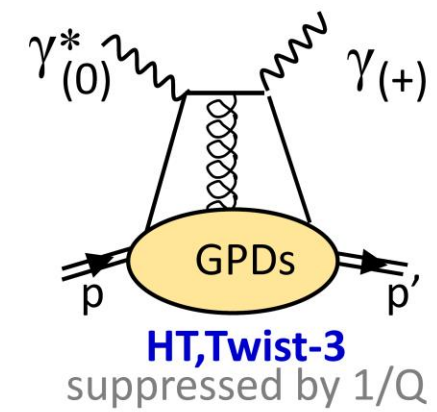
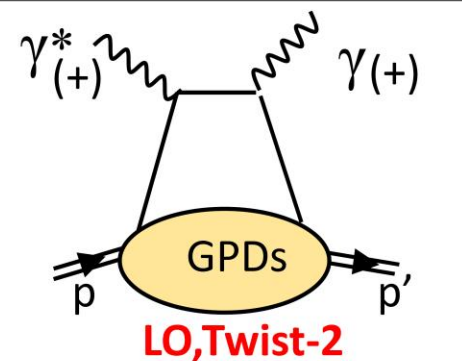
$$\begin{aligned} d\sigma^{BH} &\propto c_0^{BH} + c_1^{BH} \cos \phi + c_2^{BH} \cos 2\phi \\ d\sigma_{unpol}^{DVCS} &\propto c_0^{DVCS} + c_1^{DVCS} \cos \phi + c_2^{DVCS} \cos 2\phi \\ d\sigma_{pol}^{DVCS} &\propto s_1^{DVCS} \sin \phi \\ \text{Re } I &\propto c_0^I + c_1^I \cos \phi + c_2^I \cos 2\phi + c_3^I \cos 3\phi \\ \text{Im } I &\propto s_1^I \sin \phi + s_2^I \sin 2\phi \end{aligned}$$

Azimuthal Dependence of BH & DVCS



$$d\sigma = \left| \left[\text{BH} + \text{DVCS} \right] \right|^2 + \text{Interference Term}$$

$$\frac{d^4\sigma(\ell p \rightarrow \ell p \gamma)}{dx_B dQ^2 d|t| d\phi} = \underbrace{d\sigma^{BH}}_{\text{Well known}} + \left(d\sigma_{unpol}^{DVCS} + P_\ell d\sigma_{pol}^{DVCS} \right) + (e_\ell \text{Re } I + e_\ell P_\ell \text{Im } I)$$



With unpolarized target:

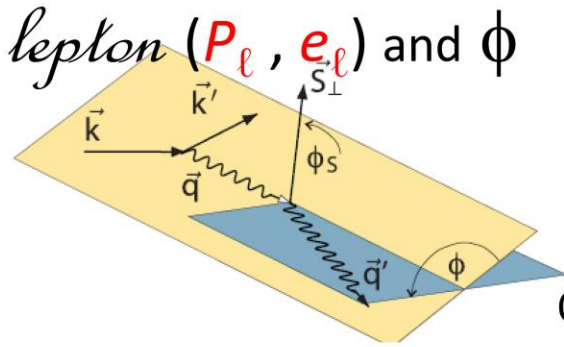
Belitsky, Müller, Kirner, NPB629 (2002)

① beam charge-spin sum

$$\Sigma \equiv d\sigma^{\leftarrow+} + d\sigma^{\rightarrow-}$$

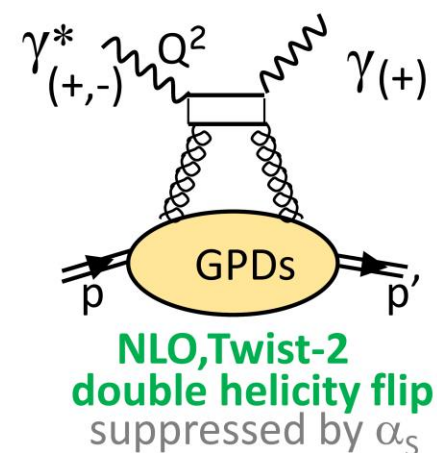
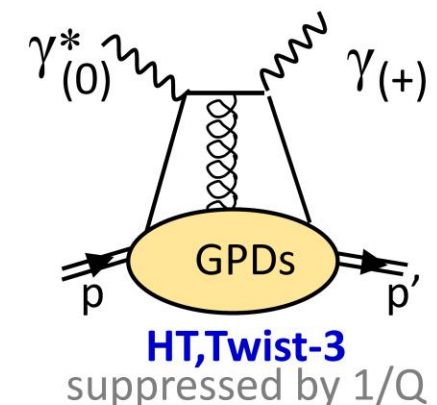
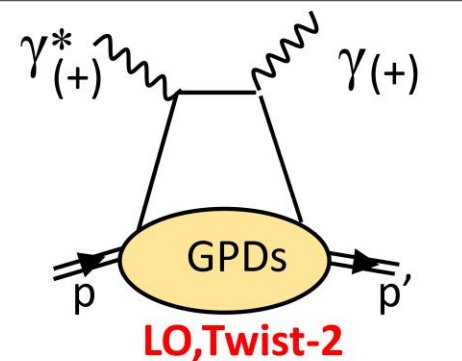
$$\begin{aligned} d\sigma^{BH} &\propto c_0^{BH} + c_1^{BH} \cos \phi + c_2^{BH} \cos 2\phi \\ d\sigma_{unpol}^{DVCS} &\propto c_0^{DVCS} + c_1^{DVCS} \cos \phi + c_2^{DVCS} \cos 2\phi \\ d\sigma_{pol}^{DVCS} &\propto s_1^{DVCS} \sin \phi \\ \text{Re } I &\propto c_0^I + c_1^I \cos \phi + c_2^I \cos 2\phi + c_3^I \cos 3\phi \\ \text{Im } I &\propto s_1^I \sin \phi + s_2^I \sin 2\phi \end{aligned}$$

Azimuthal Dependence of BH & DVCS



$$d\sigma = \left| \left[\text{BH} + \text{DVCS} \right] \right|^2 + \text{Interference Term}$$

$$\frac{d^4\sigma(\ell p \rightarrow \ell p \gamma)}{dx_B dQ^2 d|t| d\phi} = \underbrace{d\sigma^{BH}}_{\text{Well known}} + \left(d\sigma_{unpol}^{DVCS} + P_\ell d\sigma_{pol}^{DVCS} \right) + (e_\ell \text{Re } I + e_\ell P_\ell \text{Im } I)$$



With unpolarized target:

Belitsky, Müller, Kirner, NPB629 (2002)

1 beam charge-spin sum

$$\Sigma \equiv d\sigma^{\leftarrow+} + d\sigma^{\rightarrow-}$$

2 difference

$$\Delta \equiv d\sigma^{\leftarrow+} - d\sigma^{\rightarrow-}$$

$$d\sigma^{BH} \propto c_0^{BH} + c_1^{BH} \cos \phi + c_2^{BH} \cos 2\phi$$

$$d\sigma_{unpol}^{DVCS} \propto c_0^{DVCS} + c_1^{DVCS} \cos \phi + c_2^{DVCS} \cos 2\phi$$

$$d\sigma_{pol}^{DVCS} \propto s_1^{DVCS} \sin \phi$$

$$\text{Re } I \propto c_0^I + c_1^I \cos \phi + c_2^I \cos 2\phi + c_3^I \cos 3\phi$$

$$\text{Im } I \propto s_1^I \sin \phi + s_2^I \sin 2\phi$$

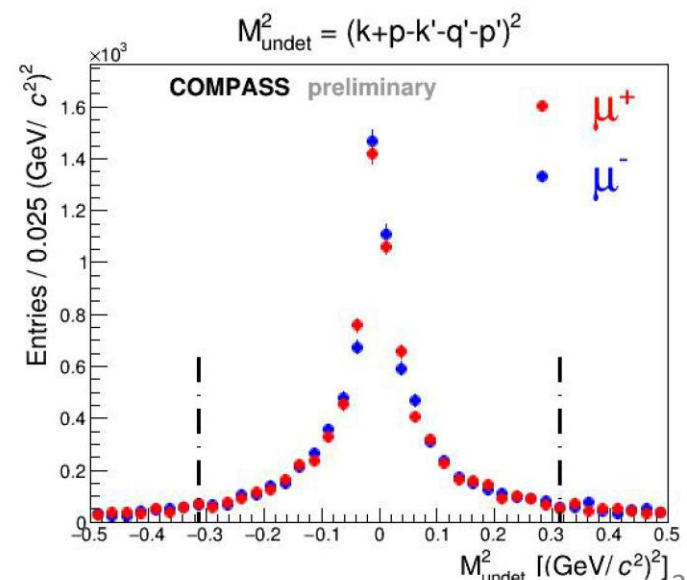
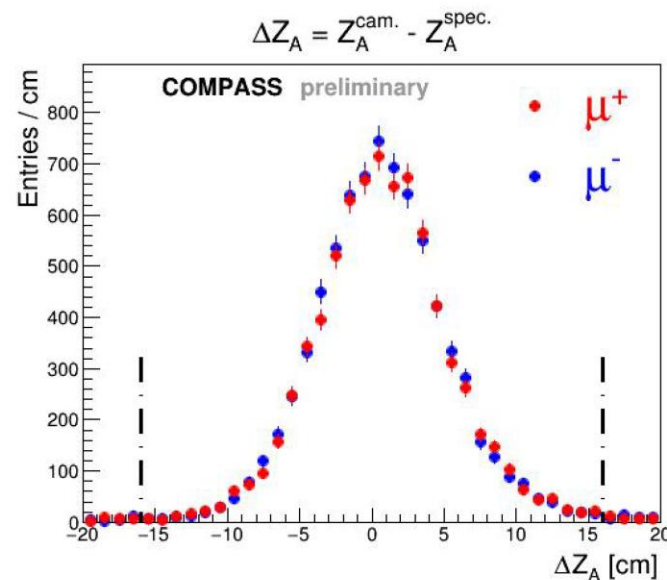
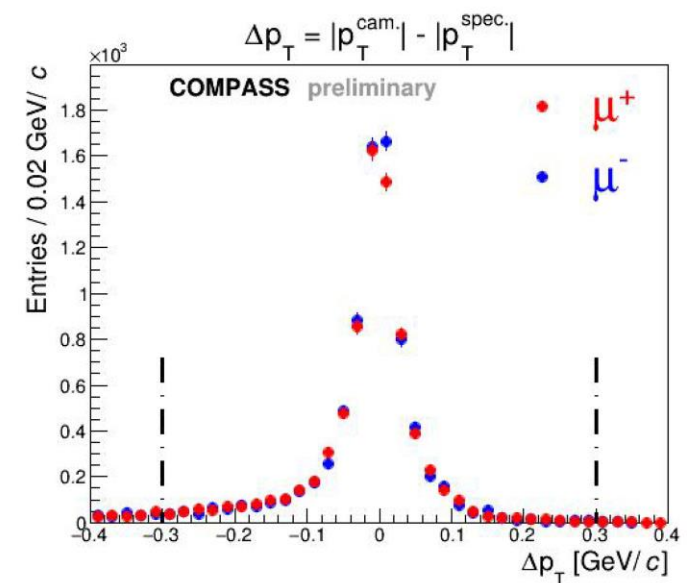
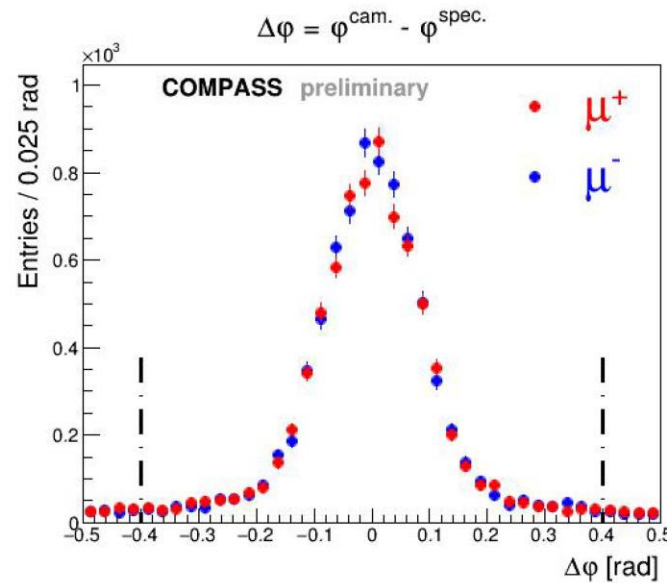
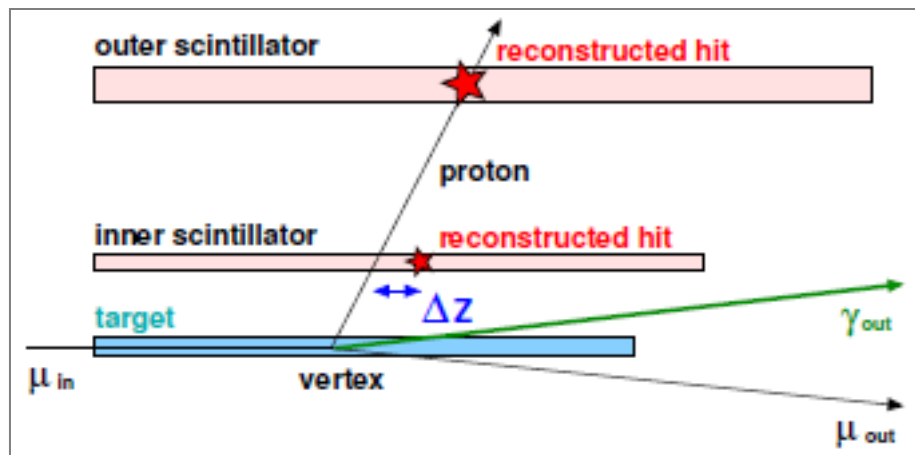
COMPASS 2016 Preliminary Results

$$\Delta\varphi = \varphi^{\text{cam.}} - \varphi^{\text{spec.}}$$

$$\Delta p_T = |p_T^{\text{cam.}}| - |p_T^{\text{spec.}}|$$

$$\Delta z_A = z_A^{\text{cam.}} - z_A^{\text{spec.}}$$

$$M_{\text{undet}}^2 = (k + p - k' - q' - p')^2$$



COMPASS 2016 Preliminary Results

➤ Main background of exclusive single photon events: π^0 decay

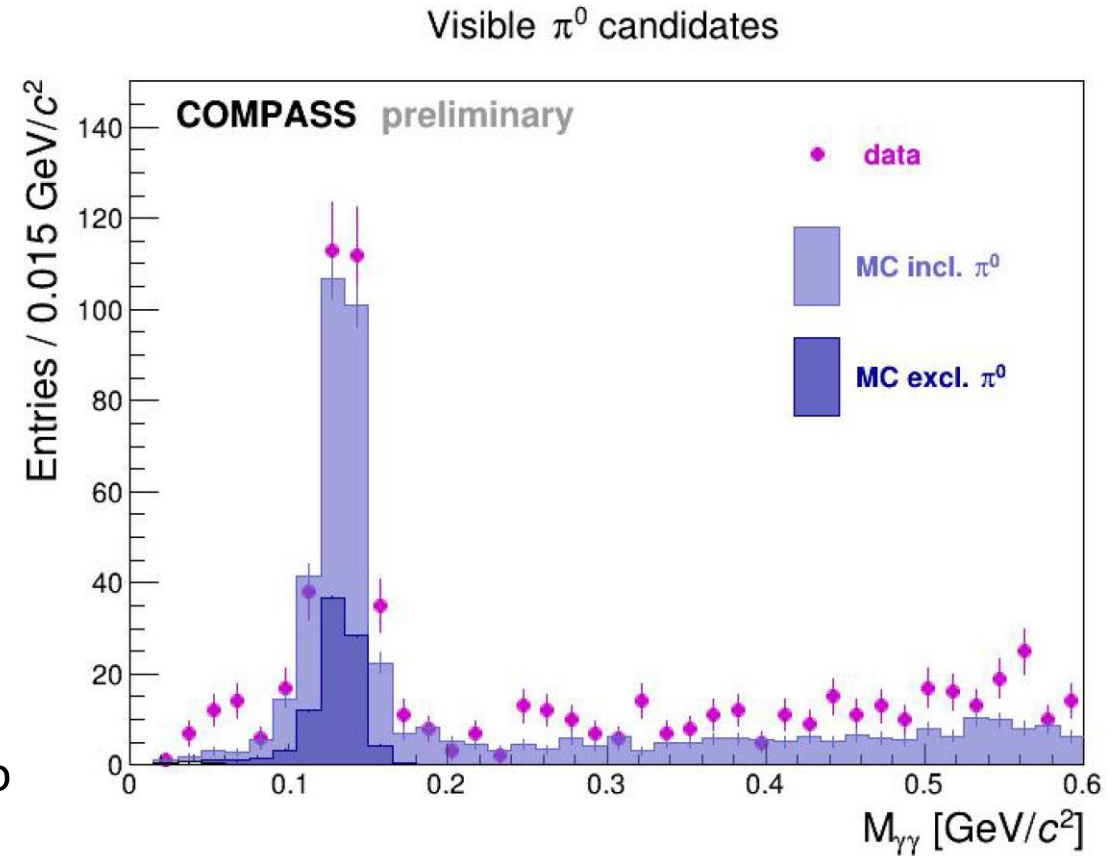
➤ **Visible** (both γ detected) – subtracted

A high-energy DVCS photon candidate is combined with all detected photons with energies lower than the DVCS threshold: (4,5) GeV in Ecal (0,1) respectively

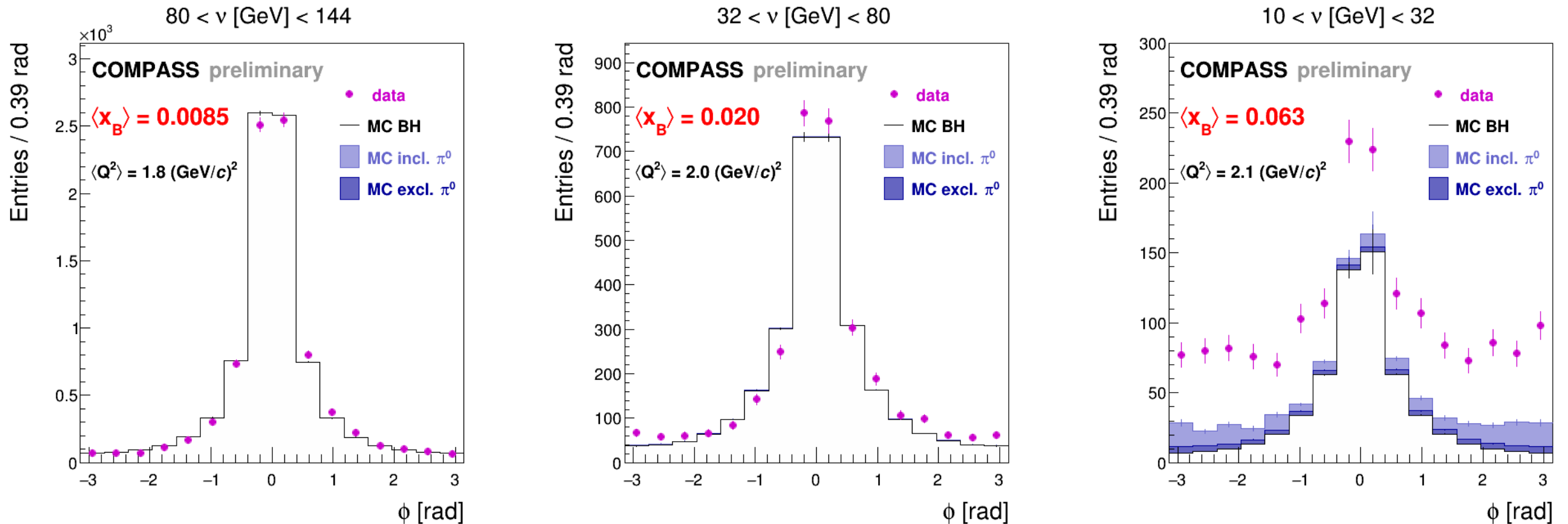
➤ **Invisible** (one γ lost) – estimated by MC

- Semi-inclusive LEPTO 6.1
- Exclusive HEPGEN π^0 (GK model)

The sum of LEPTO and HEPGEN contributions is normalized to the π^0 peak in $M_{\gamma\gamma}$ of the real data



COMPASS 2016 Preliminary Results



➤ Beam charge-spin sum

$$\begin{aligned}
 S_{CS, U}(\phi) &\equiv d\sigma(\mu^{+\leftarrow}) + d\sigma(\mu^{-\rightarrow}) = 2[d\sigma^{BH} + d\sigma_{unpol}^{DVCS} + \text{Im } I] \\
 &= 2[d\sigma^{BH} + c_0^{DVCS} + c_1^{DVCS} \cos \phi + c_2^{DVCS} \cos 2\phi + s_1^I \sin \phi + s_2^I \sin 2\phi]
 \end{aligned}$$

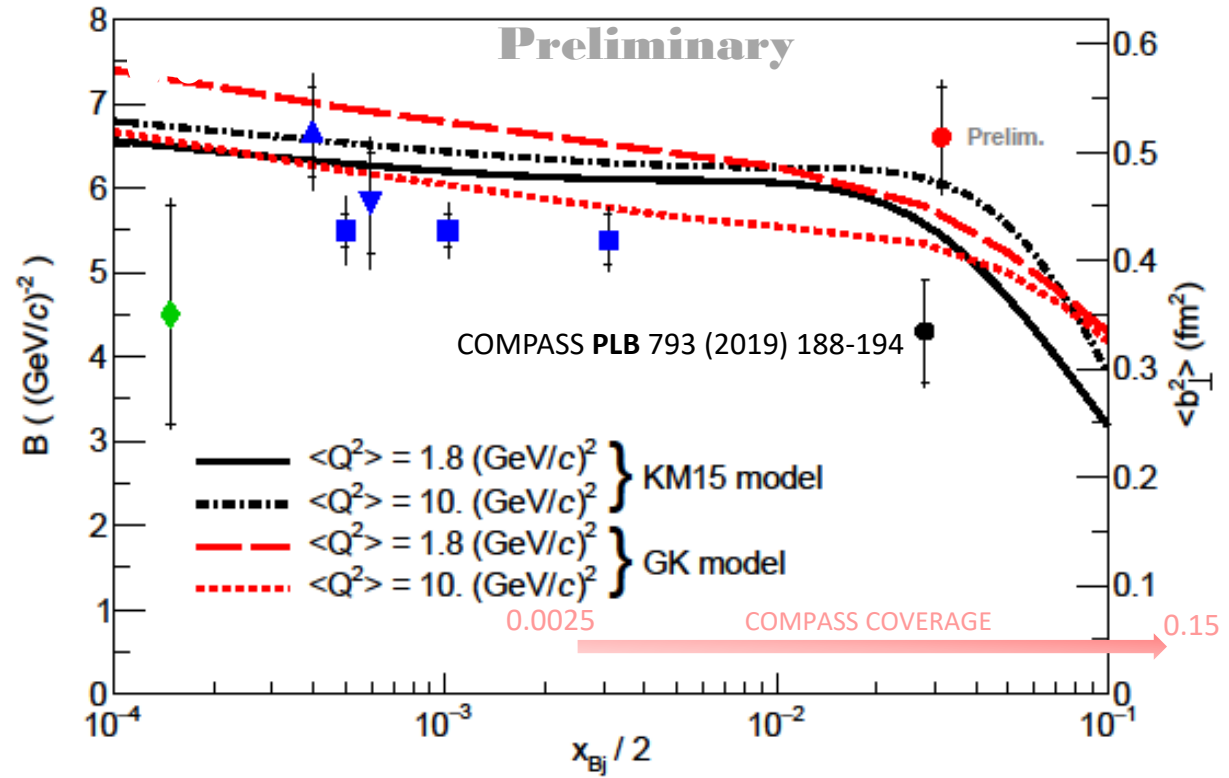
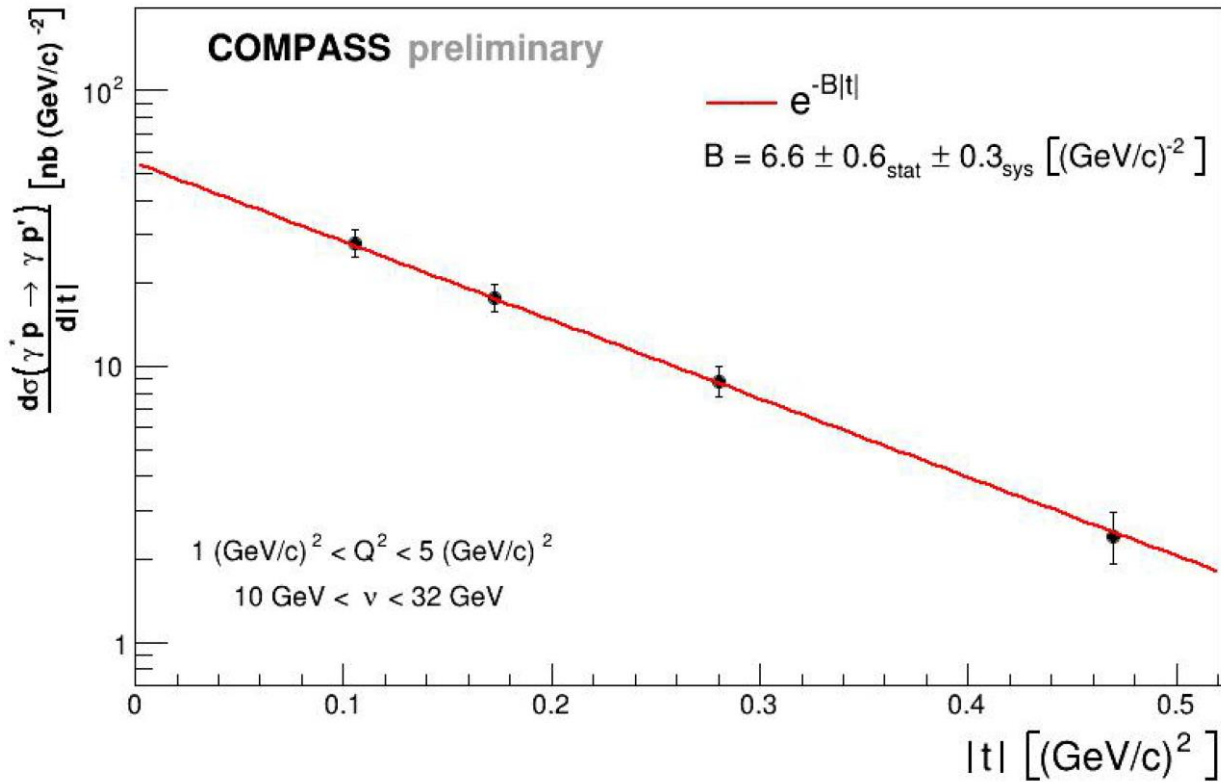
$$c_0^{DVCS} \underset{\text{small } x_{Bj}}{\propto} 4(\mathcal{H}\mathcal{H}^* + \tilde{\mathcal{H}}\tilde{\mathcal{H}}^*) + \frac{t}{M^2} \mathcal{E}\mathcal{E}^* \rightarrow 4 (\text{Im } \mathcal{H})^2 \underset{\text{model dependent}}{}$$

Transverse extension of partons – 2016 data

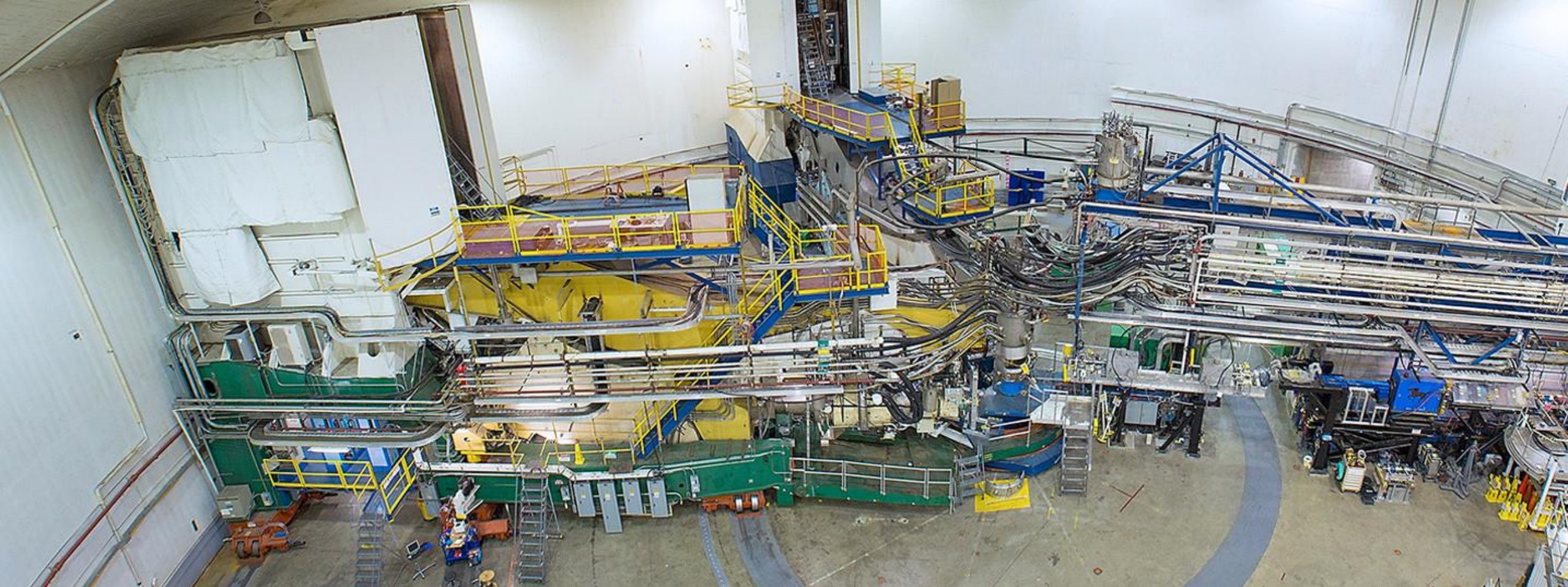
$$d\sigma^{DVCS}/d|t| \propto e^{-B|t|}$$

$$\langle r_{\perp}^2(x_B) \rangle \approx 2B(x_B) \quad \text{At small } x_B$$

- COMPASS: $\langle Q^2 \rangle = 1.8 \text{ (GeV/c)}^2$
- ◆ ZEUS: $\langle Q^2 \rangle = 3.2 \text{ (GeV/c)}^2$
- ▲ H1: $\langle Q^2 \rangle = 4.0 \text{ (GeV/c)}^2$
- ▼ H1: $\langle Q^2 \rangle = 8.0 \text{ (GeV/c)}^2$
- H1: $\langle Q^2 \rangle = 10. \text{ (GeV/c)}^2$



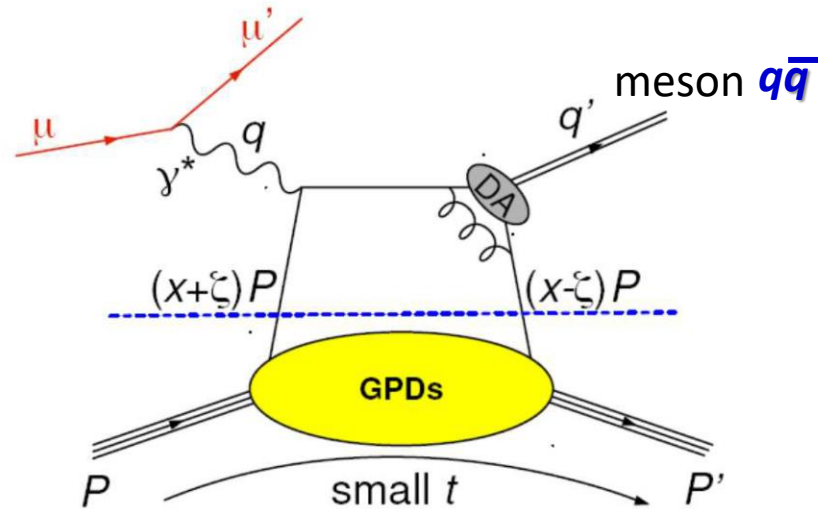
➤ The transverse-size evolution as a function of x_{Bj} → Expect at least 3 x_{Bj} bins from 2016-17 data



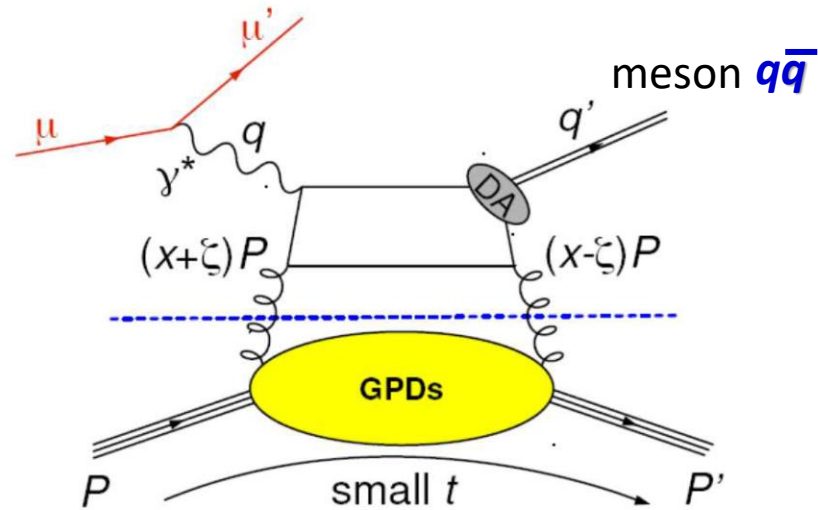
Hard Exclusive π^0 Production
@ Jlab Hall A

Deeply Virtual Meson Production (DVMP)

quark contribution



gluon contribution



4 chiral-even GPDs: helicity of parton unchanged

$$\begin{matrix} \mathbf{H}^q(x, \xi, t) & \mathbf{E}^q(x, \xi, t) \\ \tilde{\mathbf{H}}^q(x, \xi, t) & \tilde{\mathbf{E}}^q(x, \xi, t) \end{matrix}$$

+ 4 chiral-odd or transversity GPDs: helicity of parton changed

$$\begin{matrix} \mathbf{H}_T^q(x, \xi, t) & \mathbf{E}_T^q(x, \xi, t) \\ \tilde{\mathbf{H}}_T^q(x, \xi, t) & \tilde{\mathbf{E}}_T^q(x, \xi, t) \end{matrix} \quad \bar{\mathbf{E}}_T^q = 2 \tilde{\mathbf{H}}_T^q + \mathbf{E}_T^q$$

- Ability to probe the **chiral-odd GPDs**.
- Additional non-perturbative term from meson wave function → more difficult for GPD extraction
- In addition to nuclear structure, provide insights into reaction mechanism

Exclusive π^0 Production

$e p \rightarrow e \pi^0 p$

$$\frac{d^4\sigma}{dQ^2 dx_B dt d\phi} = \frac{1}{2\pi} \Gamma_\gamma(Q^2, x_B, E) \left[\frac{d\sigma_T}{dt} + \epsilon \frac{d\sigma_L}{dt} + \sqrt{2\epsilon(1+\epsilon)} \frac{d\sigma_{TL}}{dt} \cos(\phi) \right. \\ \left. + \epsilon \frac{d\sigma_{TT}}{dt} \cos(2\phi) + h \sqrt{2\epsilon(1-\epsilon)} \frac{d\sigma_{TL}}{dt} \sin(\phi) \right]$$

ϵ : degree of longitudinal polarization
 h : helicity of the initial lepton

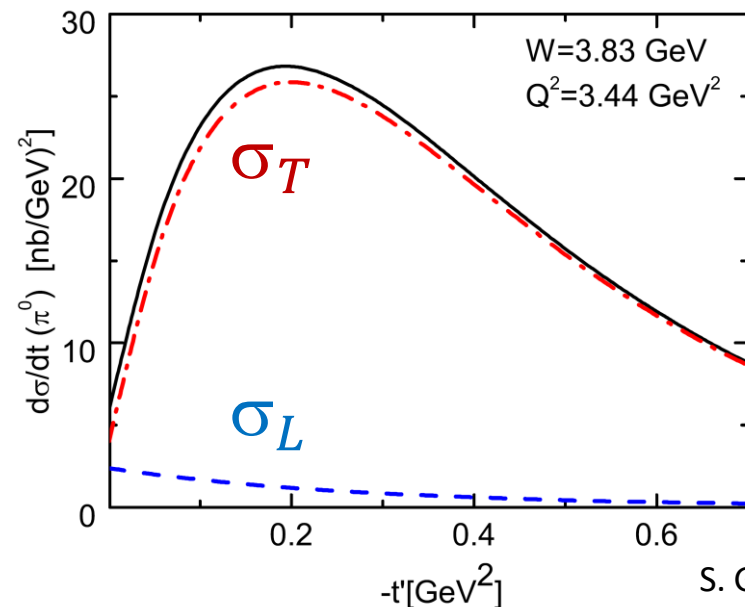
- Factorization proven only for σ_L , which depends on chiral-even GPDs only
- At sufficiently high Q^2 , expect $\sigma_L \propto Q^{-6}$ while σ_T asymptotically suppressed and $\propto Q^{-8}$
→ σ_L dominance
- Previous experiments with limited reach in Q^2 suggest the dominance of σ_T

Exclusive π^0 Production

$e p \rightarrow e \pi^0 p$

$$\frac{d^4\sigma}{dQ^2 dx_B dt d\phi} = \frac{1}{2\pi} \Gamma_\gamma(Q^2, x_B, E) \left[\frac{d\sigma_T}{dt} + \epsilon \frac{d\sigma_L}{dt} + \sqrt{2\epsilon(1+\epsilon)} \frac{d\sigma_{TL}}{dt} \cos(\phi) + \epsilon \frac{d\sigma_{TT}}{dt} \cos(2\phi) + h \sqrt{2\epsilon(1-\epsilon)} \frac{d\sigma_{TL}}{dt} \sin(\phi) \right]$$

ϵ : degree of longitudinal polarization
 h : helicity of the initial lepton

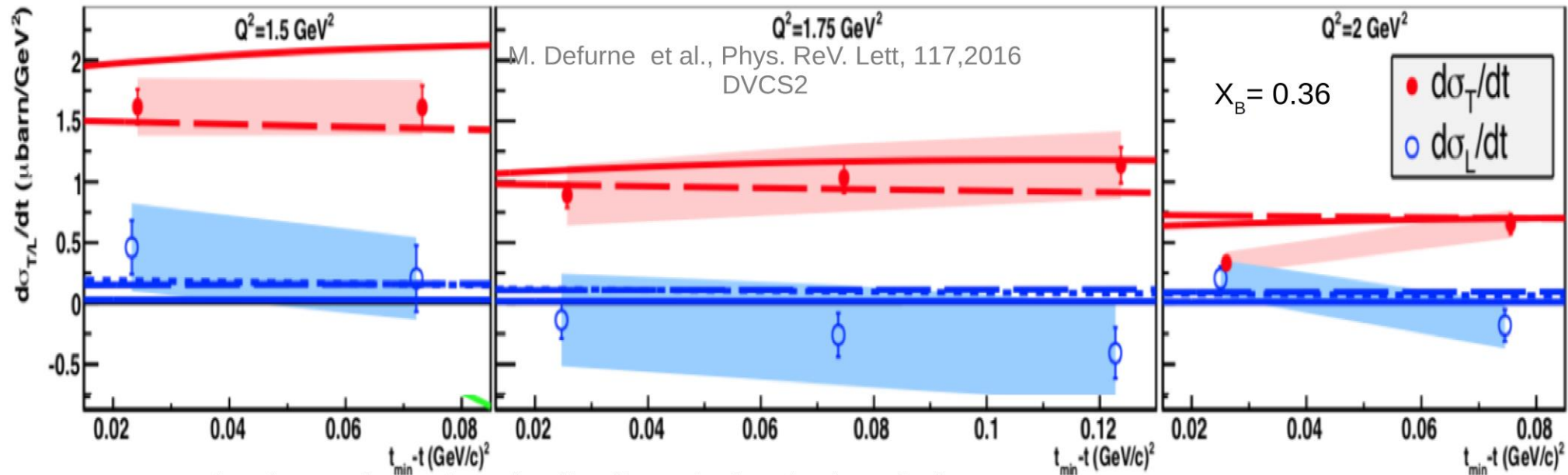


➤ Modeling of $\sigma_T \rightarrow$ coupling between transversity GPDs and twist-3 pion amplitude

Exclusive π^0 Production

$e p \rightarrow e \pi^0 p$

$$\frac{d^4\sigma}{dQ^2 dx_B dt d\phi} = \frac{1}{2\pi} \Gamma_\gamma(Q^2, x_B, E) \left[\frac{d\sigma_T}{dt} + \epsilon \frac{d\sigma_L}{dt} + \sqrt{2\epsilon(1+\epsilon)} \frac{d\sigma_{TL}}{dt} \cos(\phi) + \epsilon \frac{d\sigma_{TT}}{dt} \cos(2\phi) + h \sqrt{2\epsilon(1-\epsilon)} \frac{d\sigma_{TL}}{dt} \sin(\phi) \right]$$



Exclusive π^0 Production

$e p \rightarrow e \pi^0 p$

$$\frac{d^4\sigma}{dQ^2 dx_B dt d\phi} = \frac{1}{2\pi} \Gamma_\gamma(Q^2, x_B, E) \left[\frac{d\sigma_T}{dt} + \epsilon \frac{d\sigma_L}{dt} + \sqrt{2\epsilon(1+\epsilon)} \frac{d\sigma_{TL}}{dt} \cos(\phi) + \epsilon \frac{d\sigma_{TT}}{dt} \cos(2\phi) + h \sqrt{2\epsilon(1-\epsilon)} \frac{d\sigma_{TL}}{dt} \sin(\phi) \right]$$

ϵ : degree of longitudinal polarization
 h : helicity of the initial lepton

- $\frac{d\sigma_L}{dt} = \frac{4\pi\alpha}{k'} \frac{1}{Q^6} \left\{ (1-\xi^2) |\langle \tilde{H} \rangle|^2 - 2\xi^2 \text{Re} [\langle \tilde{H} \rangle^* \langle \tilde{E} \rangle] - \frac{t'}{4m^2} \xi^2 |\langle \tilde{E} \rangle|^2 \right\}$
- $\frac{d\sigma_T}{dt} = \frac{4\pi\alpha}{2k'} \frac{\mu_\pi^2}{Q^8} \left[(1-\xi^2) |\langle H_T \rangle|^2 - \frac{t'}{8m^2} |\langle \bar{E}_T \rangle|^2 \right]$
- $\frac{\sigma_{LT}}{dt} = \frac{4\pi\alpha}{\sqrt{2}k'} \frac{\mu_\pi}{Q^7} \xi \sqrt{1-\xi^2} \frac{\sqrt{-t'}}{2m} \text{Re} [\langle H_T \rangle^* \langle \tilde{E} \rangle]$
- $\frac{\sigma_{TT}}{dt} = \frac{4\pi\alpha}{k'} \frac{\mu_\pi^2}{Q^8} \frac{t'}{16m^2} |\langle \bar{E}_T \rangle|^2$

$$\bar{E}_T = 2\tilde{H}_T + E_T$$

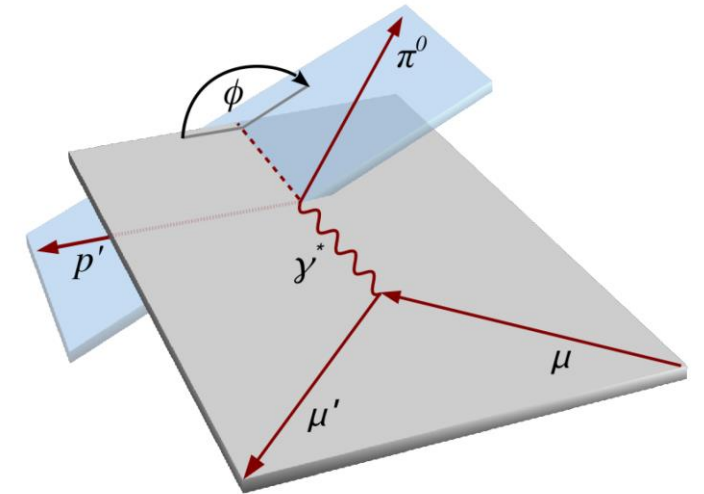
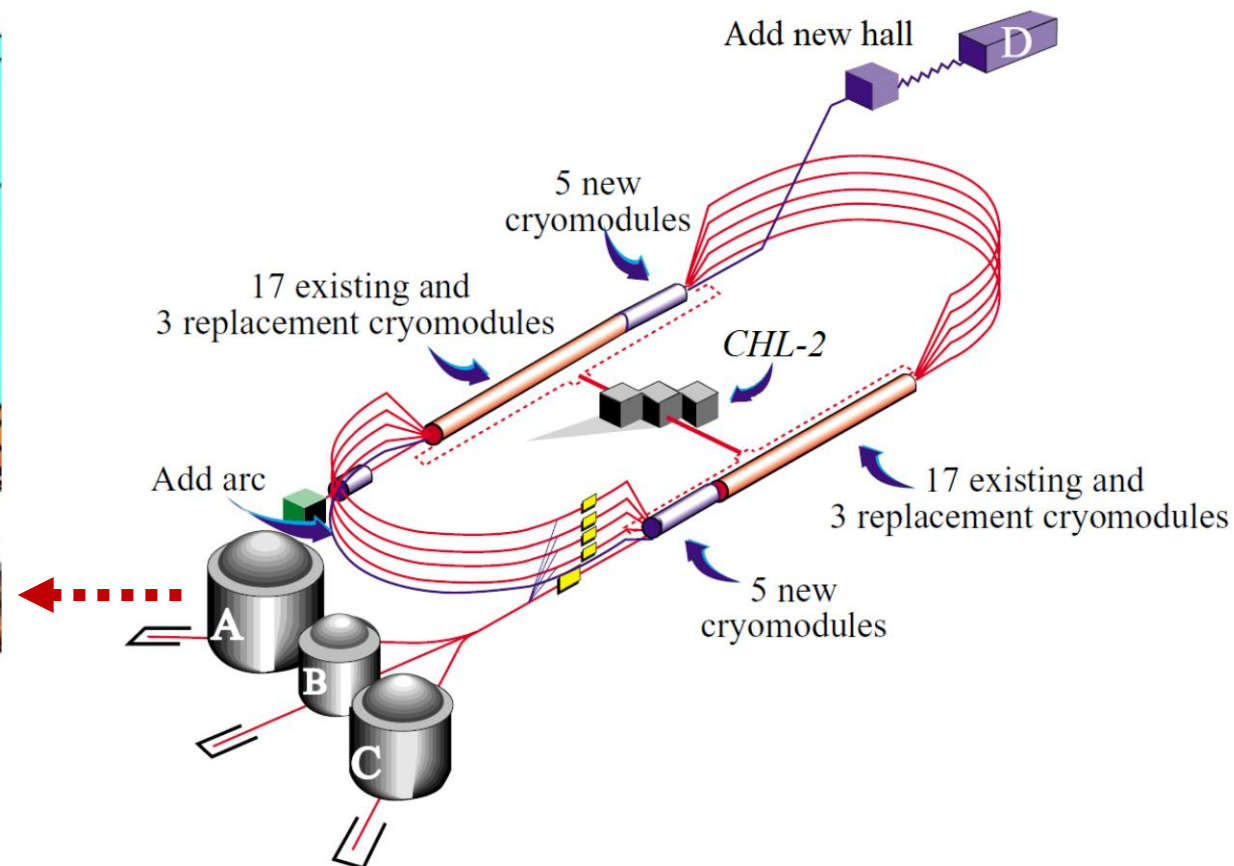
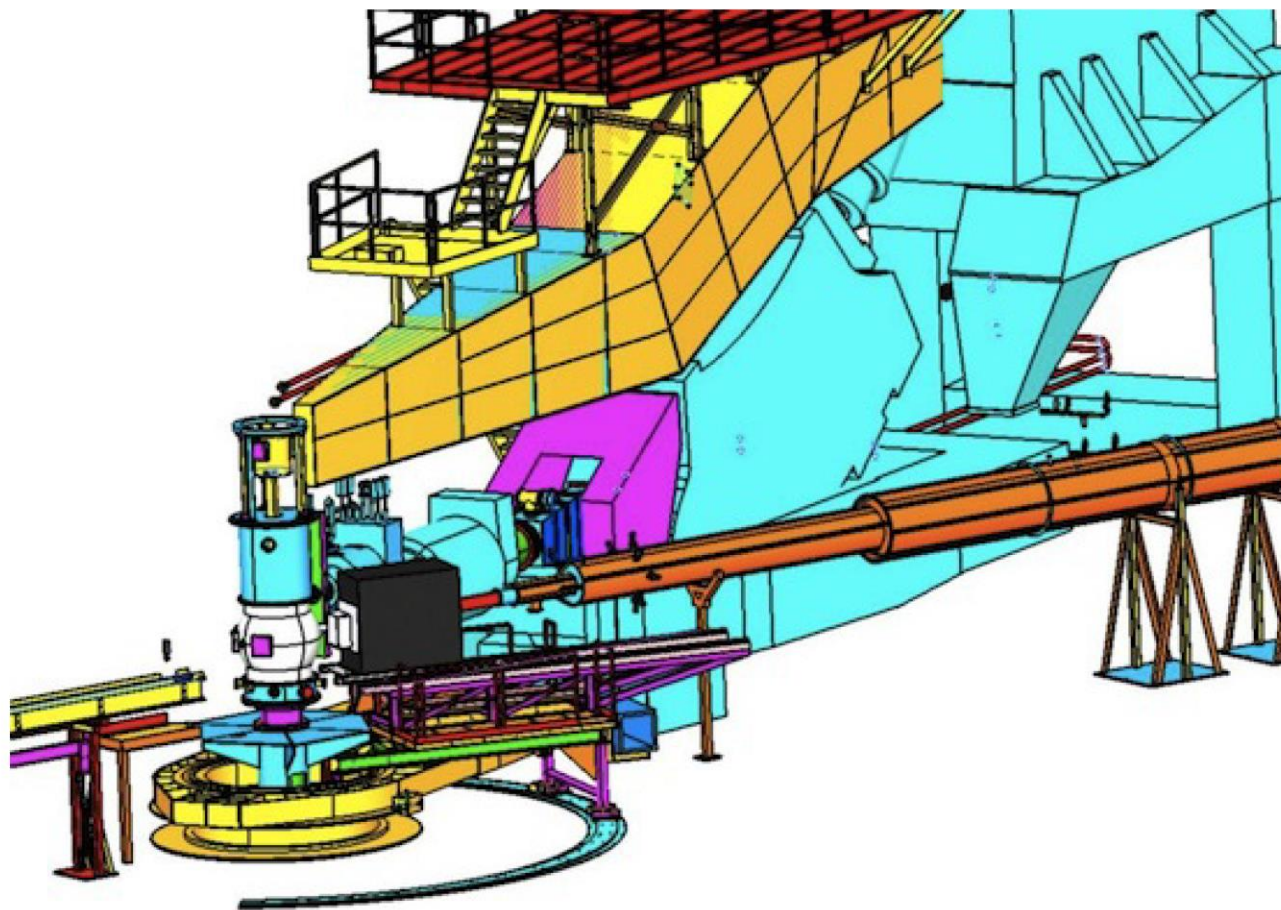


Fig: M.G. Alexeev et al. *Phys.Lett.B* 805 (2020)

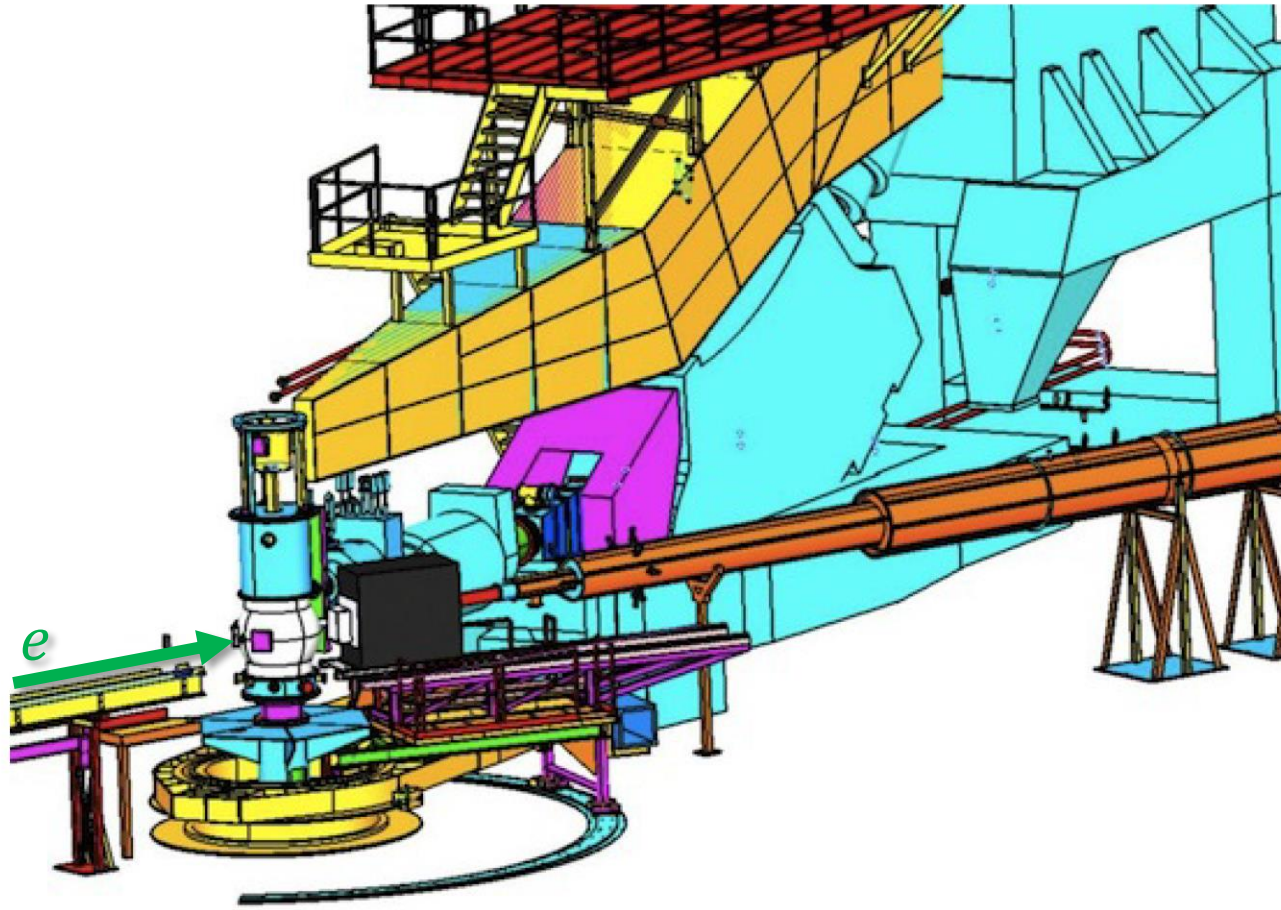
Jefferson Lab Hall A experiment E12-06-114

https://www.jlab.org/div_dept/physics_division/GeV/whitepaperv11/index.html

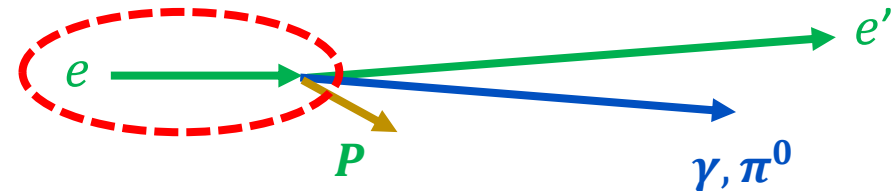


➤ 3rd Generation DVCS project @ Hall A → CEBAF12 grants the ability to explore high x_B with extended Q^2 .

E12-06-114 Experimental Setup

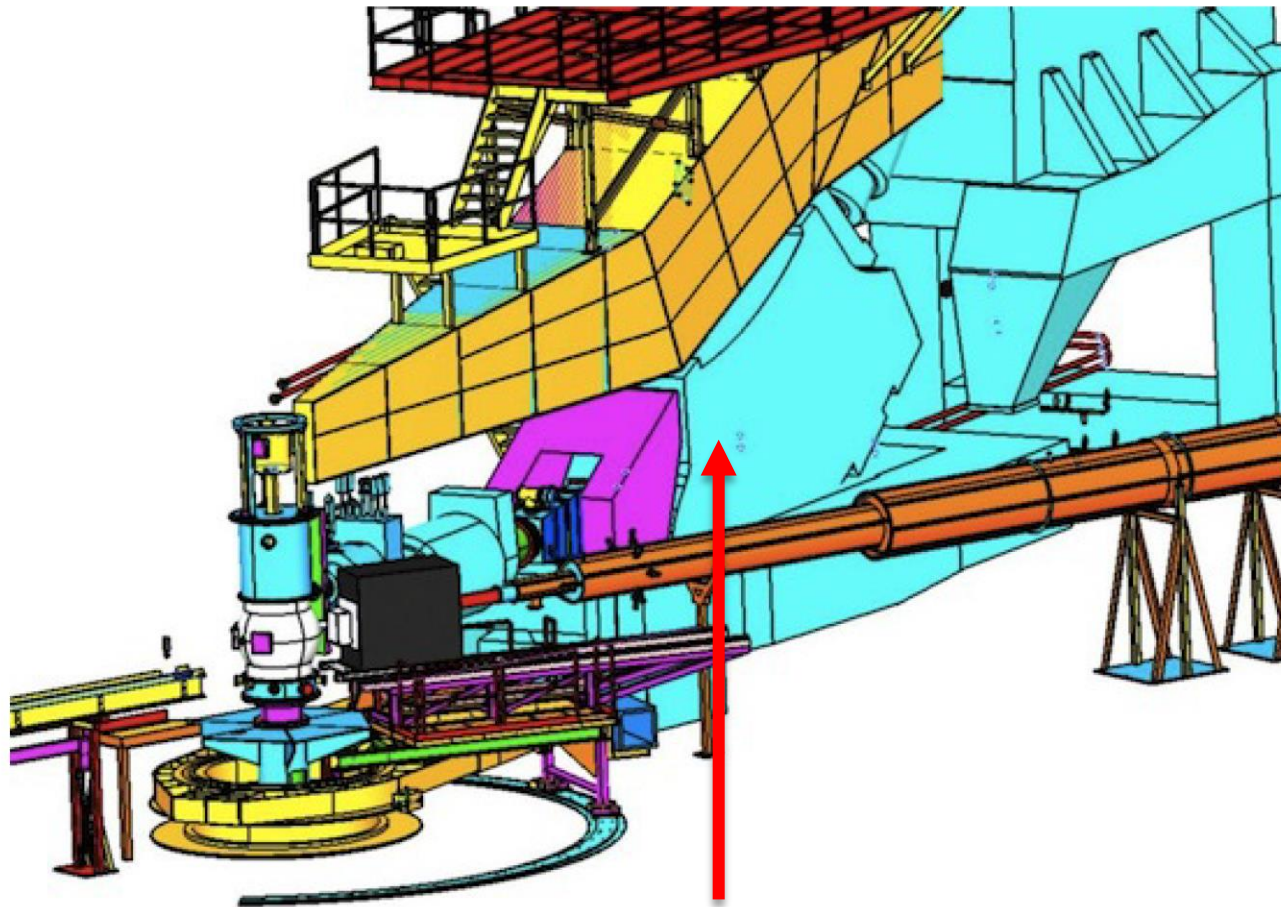


DVCS & Exclusive π^0 Production



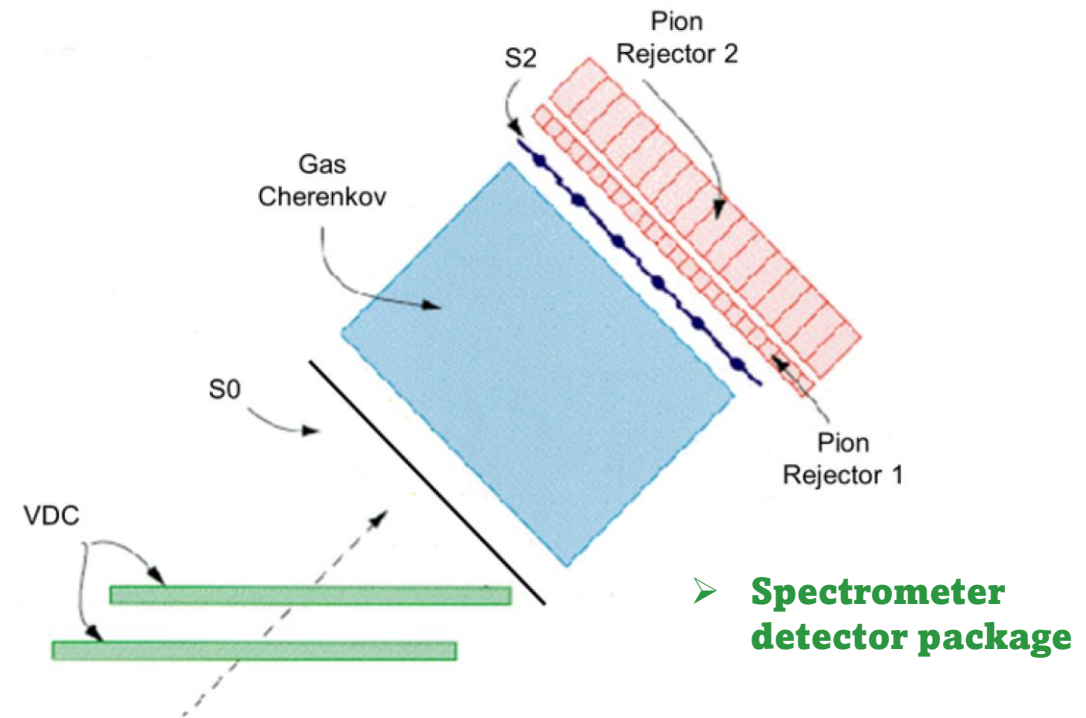
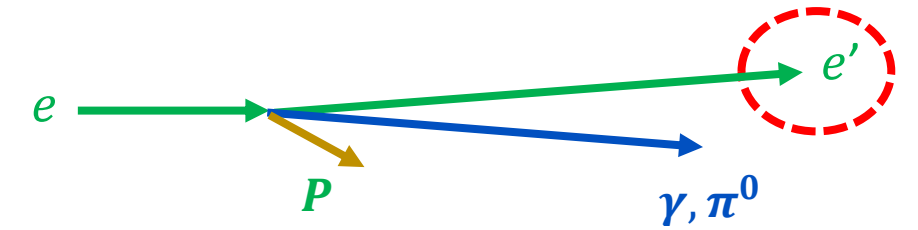
- **Electron beam**
 - polarisation $\sim 85\%$
 - helicity flipped at 30 Hz
 - luminosity: $\sim 10^{38}$ Hz/cm²
- **LH₂ target**
 - 6.35 cm diameter, 15 cm long

E12-06-114 Experimental Setup



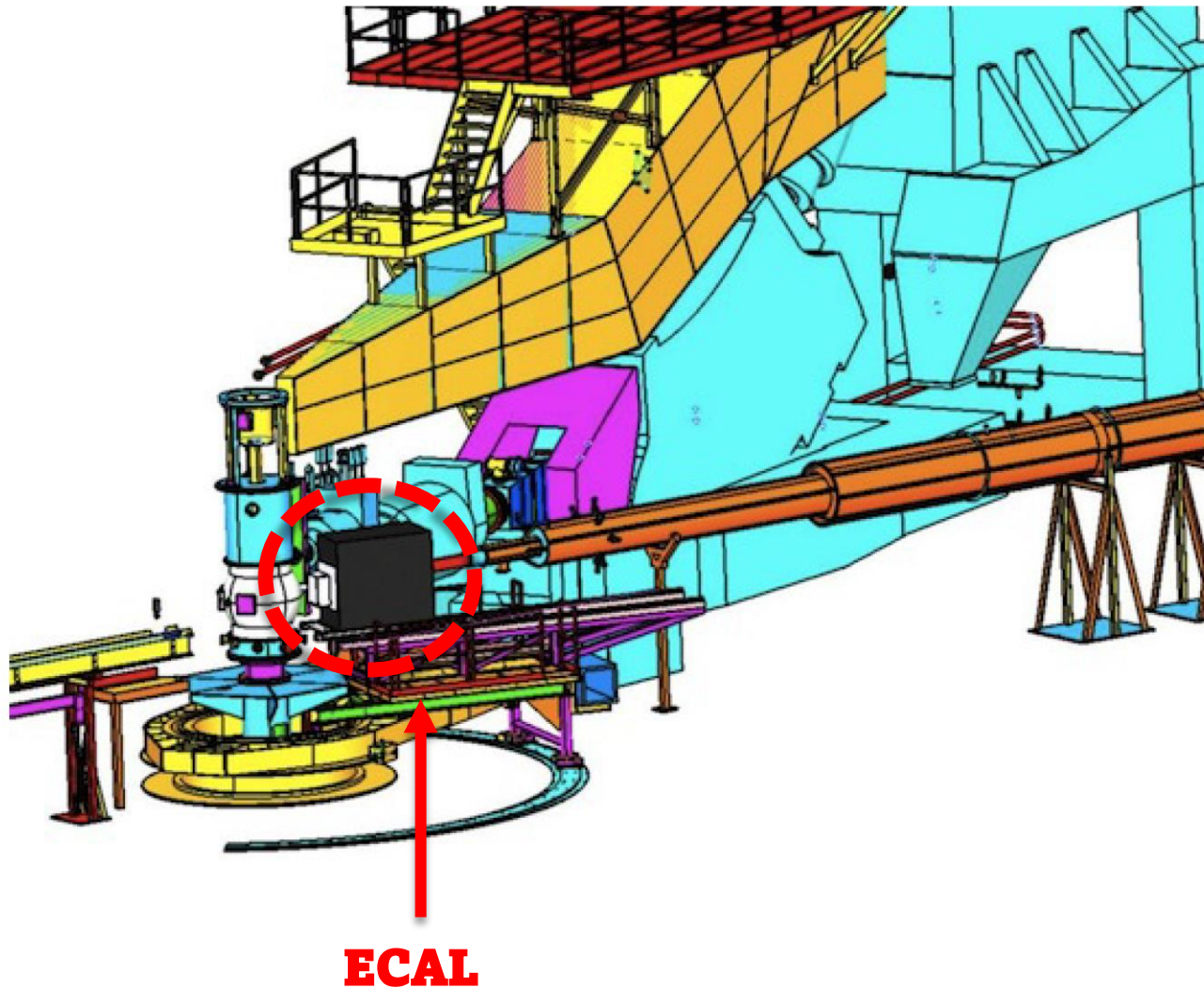
Left High Resolution Spectrometer (LHRS)

DVCS & Exclusive π^0 Production

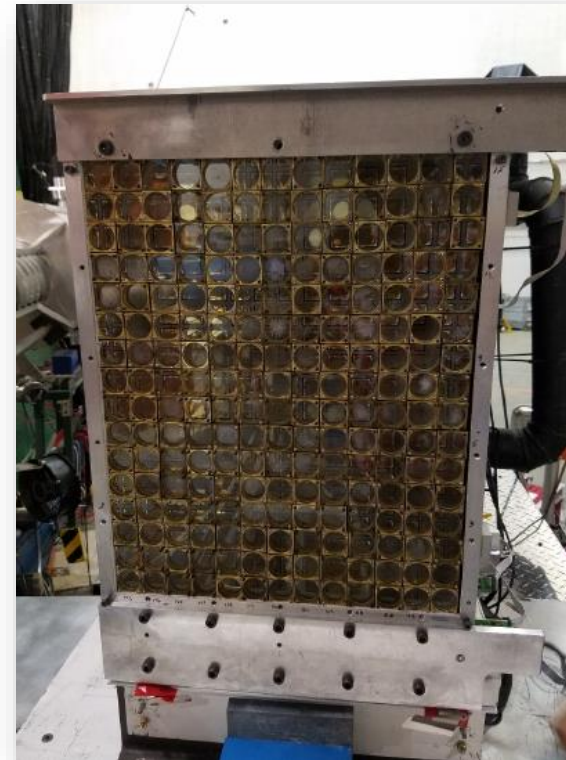
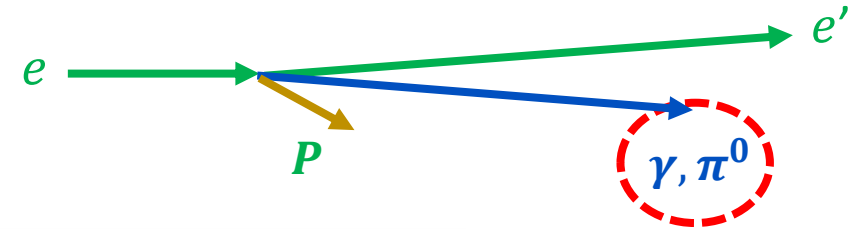


➤ $\delta P/P$ resolution $\sim 10^{-4}$ @ 4.3 GeV

E12-06-114 Experimental Setup



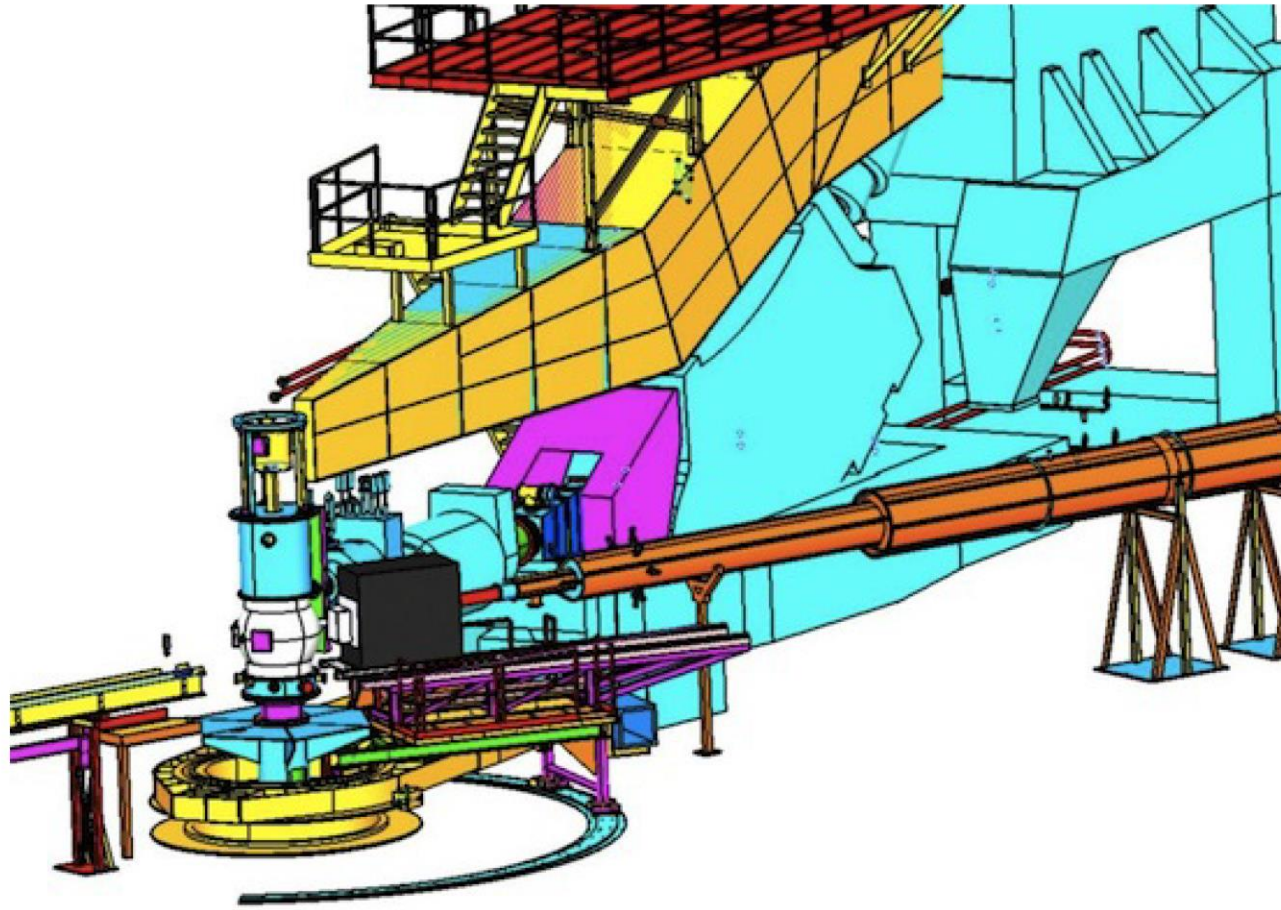
DVCS & Exclusive π^0 Production



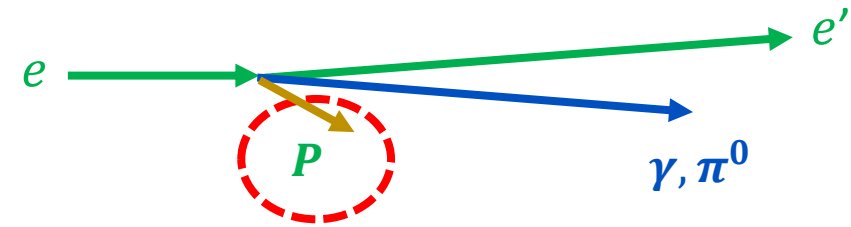
➤ ECAL

- 208 blocks of PF_2
- $3 \times 3 \times 18.6 \text{ cm}^3$ each
- ~ 20 radiation lengths
- Moliere radius 2.2cm
- E resolution $\sim 3\%$
@ 4.2 GeV

E12-06-114 Experimental Setup



DVCS & Exclusive π^0 Production

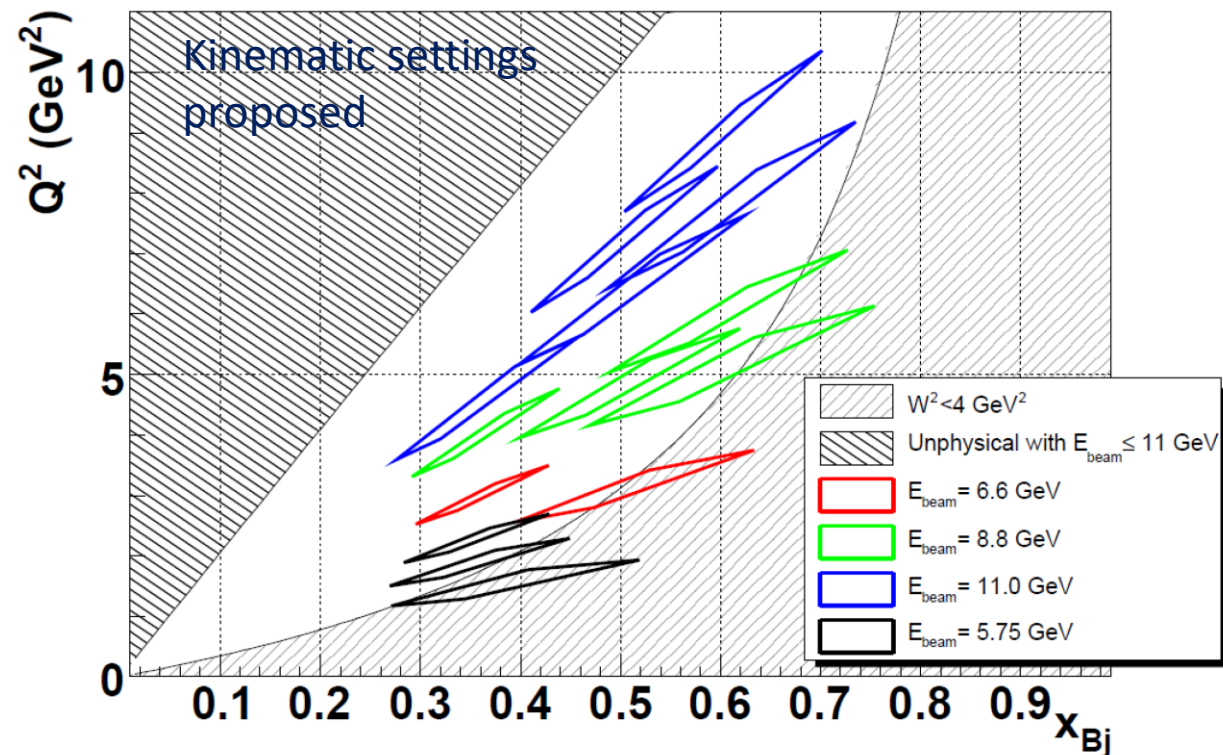


Recoil Proton

- Not detected
- Exclusivity of events ensured using missing mass, M_X^2

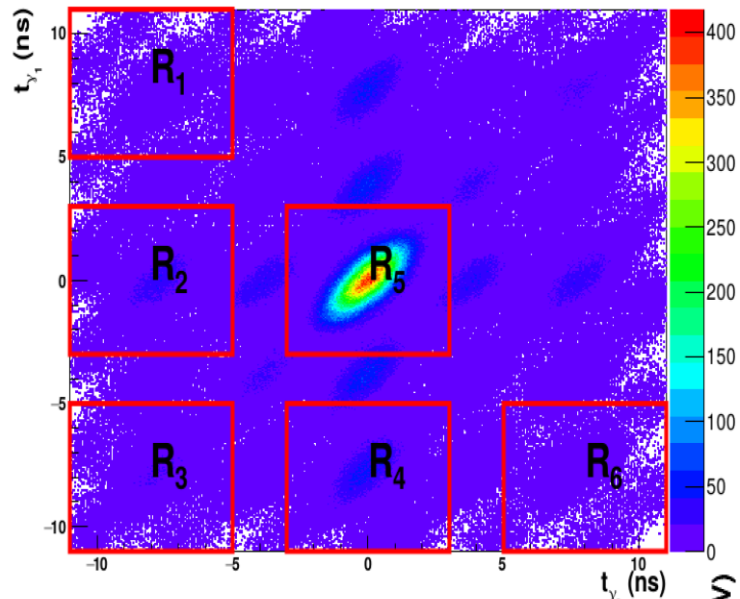
E12-06-114 Kienmatic Settings

x_B label	0.36			0.48				0.60	
$\langle x_B \rangle$	0.36	0.36	0.36	0.48	0.45	0.46	0.46	0.59	0.60
E (GeV)	7.38	8.52	10.59	4.49	8.85	8.85	10.99	8.52	10.59
Q^2 (GeV ²)	3.11	3.57	4.44	2.67	4.06	5.16	6.56	5.49	8.31
W^2 (GeV ²)	6.51	7.29	8.79	3.81	5.62	6.67	8.32	4.58	6.46
$-t_{\min}$ (GeV ²)	0.16	0.17	0.17	0.33	0.35	0.35	0.36	0.67	0.71
ϵ	0.61	0.62	0.63	0.51	0.71	0.55	0.52	0.66	0.50

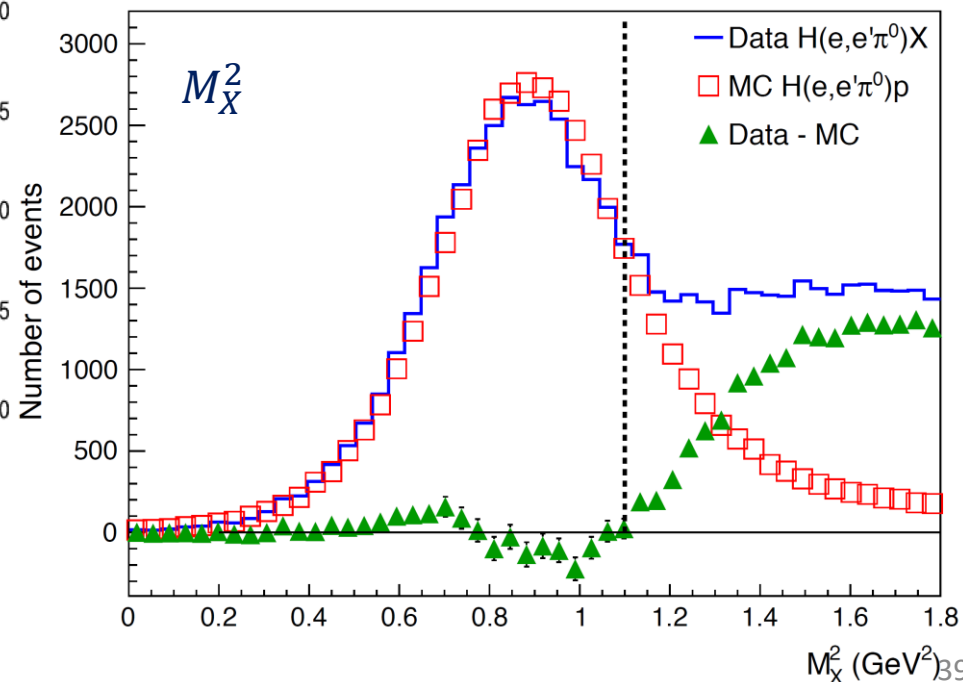
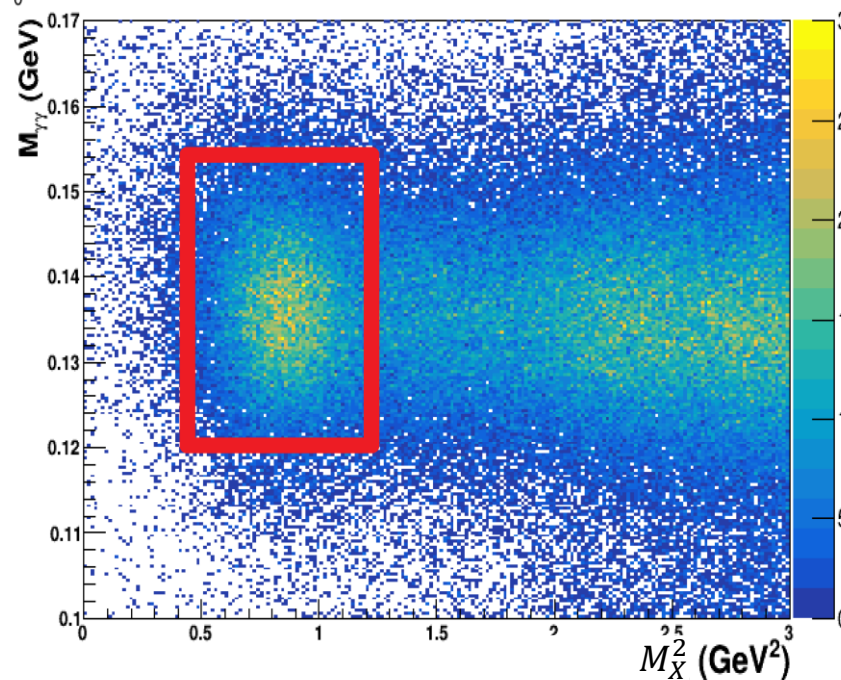
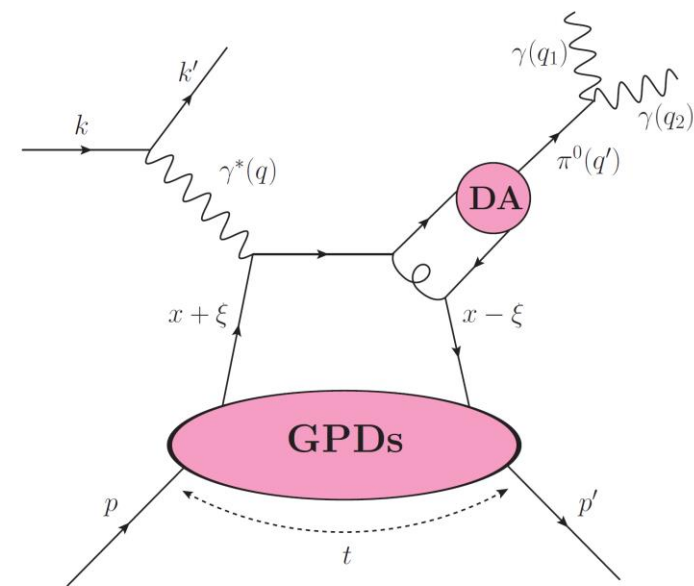


- Ran in 2014 & 2016
- 9 settings with x_B of 0.36, 0.48, and 0.6 and Q^2 ranging from about 3 to 8 GeV²
- About 50% of allocated 100 PAC days
- Missing PAC days reallocated to the future experiment @ Hall C

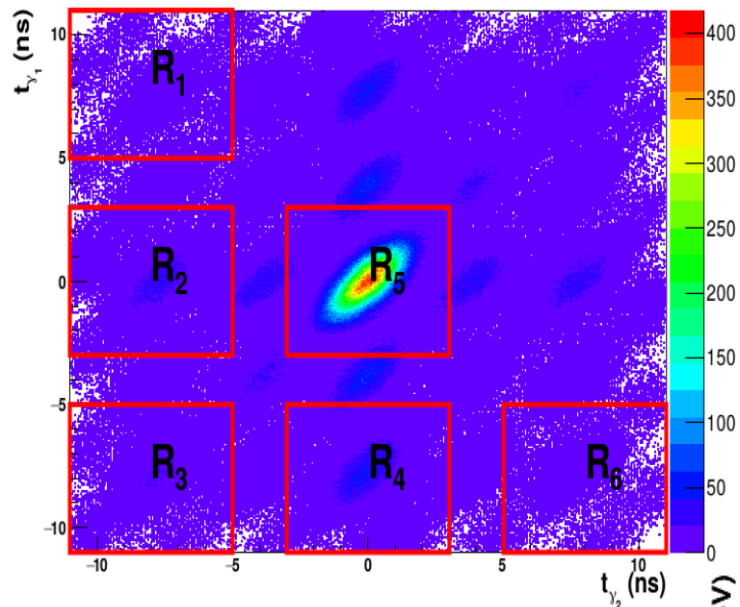
Exclusive π^0 Event Selection



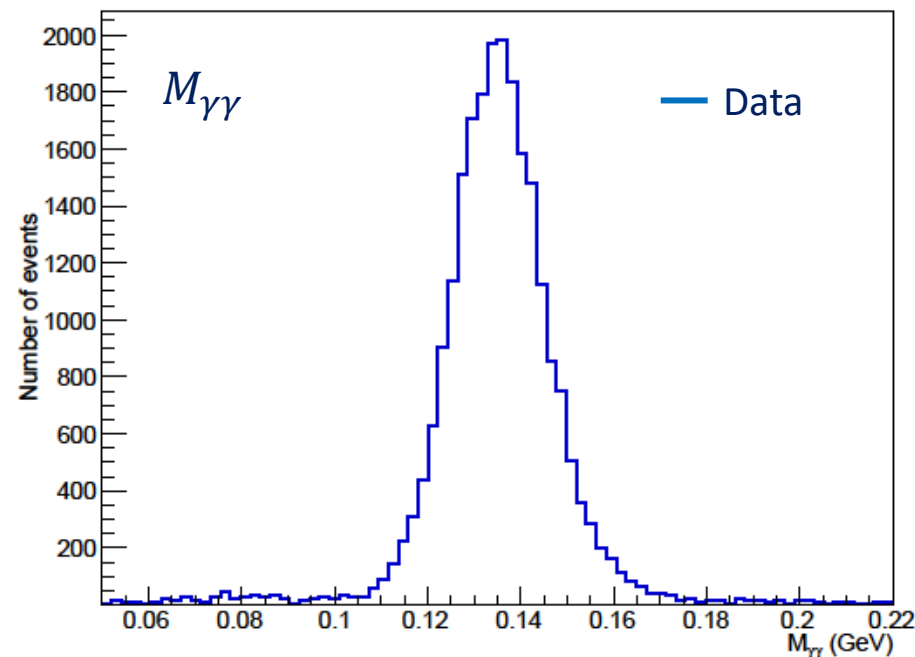
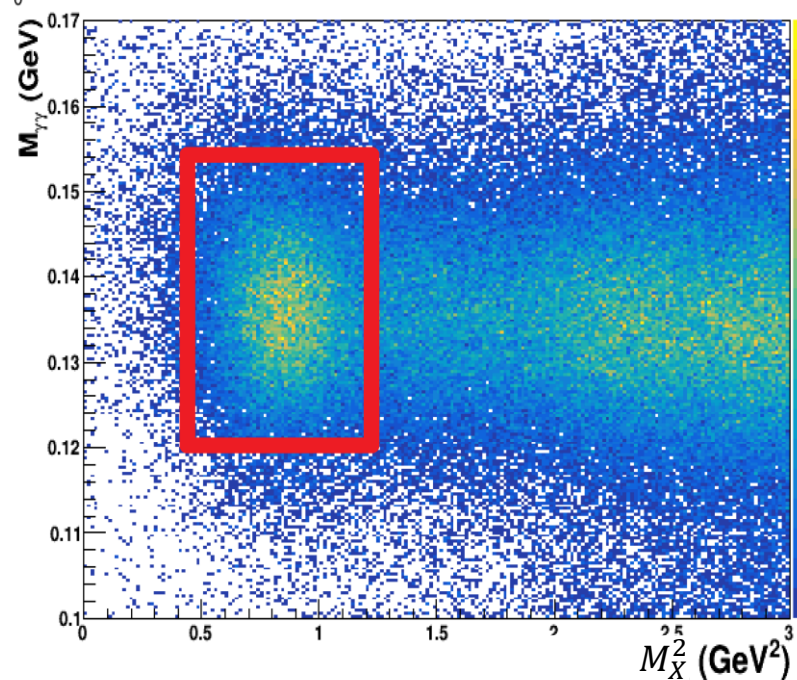
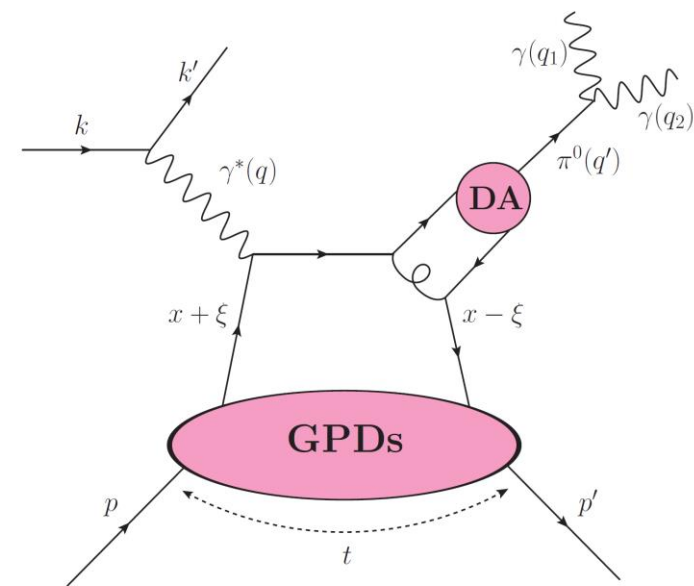
- Main background: accidentals. The background in the signal coincidence window, $[-3,3]$ ns, is estimated via other time windows.
- Exclusivity \rightarrow remove the $M_X^2 = (k + P - k' - q_1 - q_2)^2$ contribution from inclusive channels, threshold $\approx 1.15 \text{ GeV}^2$
- π^0 events \rightarrow select events with invariant mass $M_{\gamma\gamma} = \sqrt{(q_1 + q_2)^2}$ around the π^0 mass



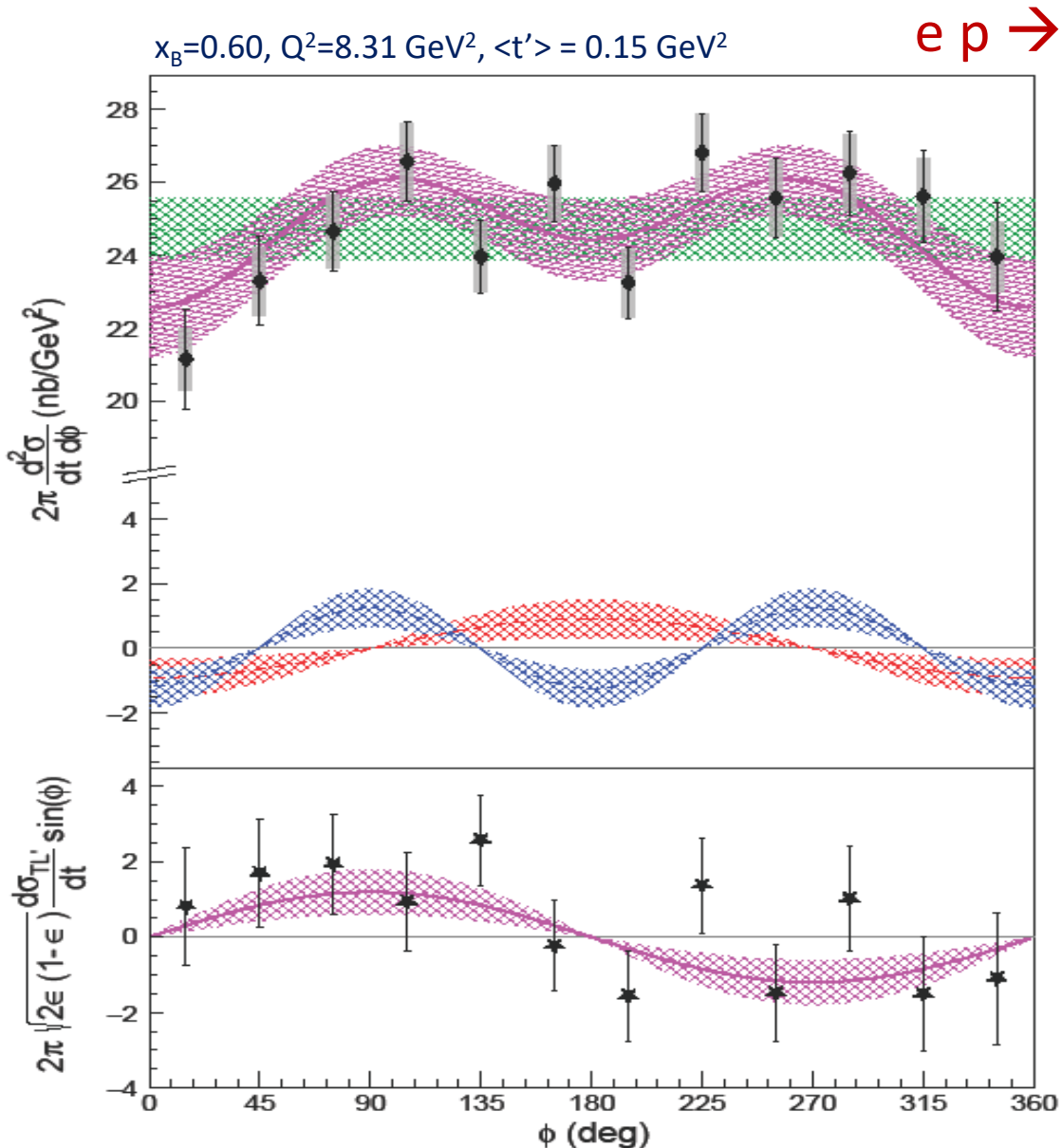
Exclusive π^0 Event Selection



- Main background: accidentals. The background in the signal coincidence window, $[-3,3]$ ns, is estimated via other time windows.
- Exclusivity \rightarrow remove the $M_X^2 = (k + P - k' - q_1 - q_2)^2$ contribution from inclusive channels, threshold $\approx 1.15 \text{ GeV}^2$
- π^0 events \rightarrow select events with invariant mass $M_{\gamma\gamma} = \sqrt{(q_1 + q_2)^2}$ around the π^0 mass



Structure-function Extraction



$$\frac{d^4\sigma}{dQ^2 dx_B dt d\phi} = \frac{1}{2\pi} \frac{d^2 \Gamma_\gamma}{dQ^2 dx_B}(Q^2, x_B, E)$$

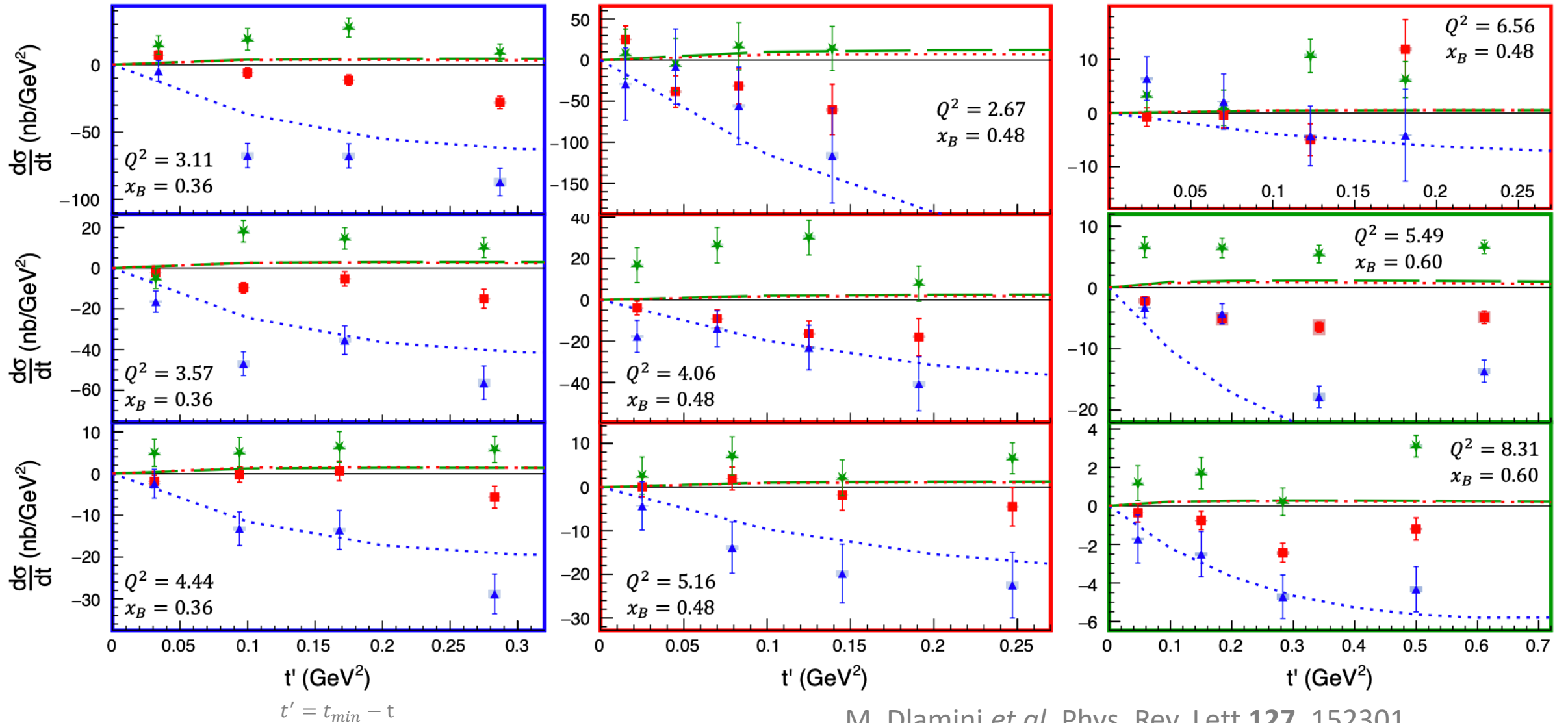
$$\left[\frac{d\sigma_T}{dt} + \epsilon \frac{d\sigma_L}{dt} \right] + \sqrt{2\epsilon(1+\epsilon)} \left[\frac{d\sigma_{LT}}{dt} \cos(\phi) + \epsilon \frac{d\sigma_{TT}}{dt} \cos(2\phi) \right. \\ \left. + h \sqrt{2\epsilon(1-\epsilon)} \left[\frac{d\sigma_{LT'}}{dt} \sin(\phi) \right] \right]$$

- Structure functions extracted for all 9 kinematic settings
- Extract different terms via their corresponding ϕ dependence
- $d\sigma_T$ and $d\sigma_L$ can't be separated, extracted as $d\sigma_U = d\sigma_T + \epsilon d\sigma_L$
- Main systematic errors come from deviation observed in DIS events and the exclusivity cuts

Structure Functions

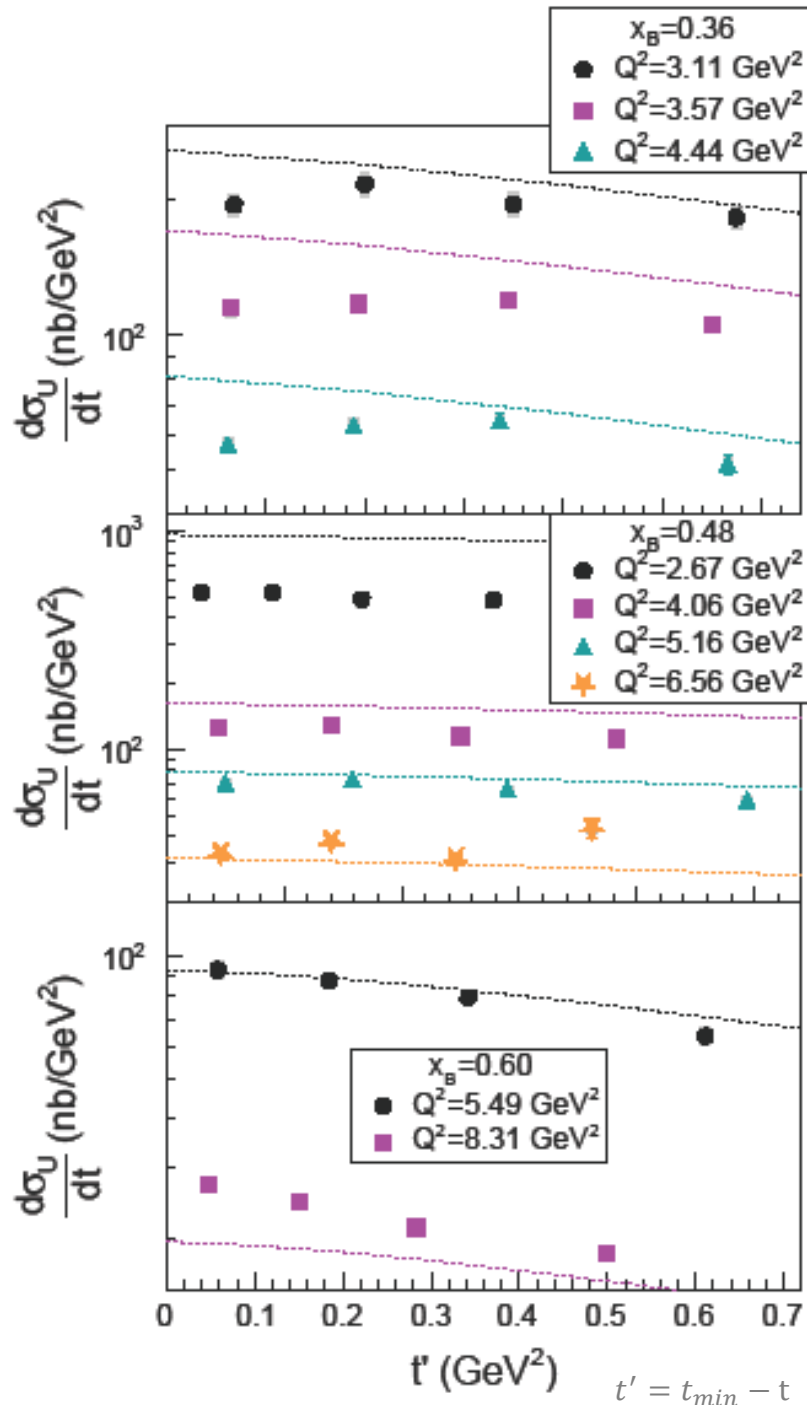
- Dotted (dashed) curves: P. Kroll, private communications

▲ $d\sigma_{TT}$ ■ $d\sigma_{LT}$ ★ $d\sigma_{LT'}$

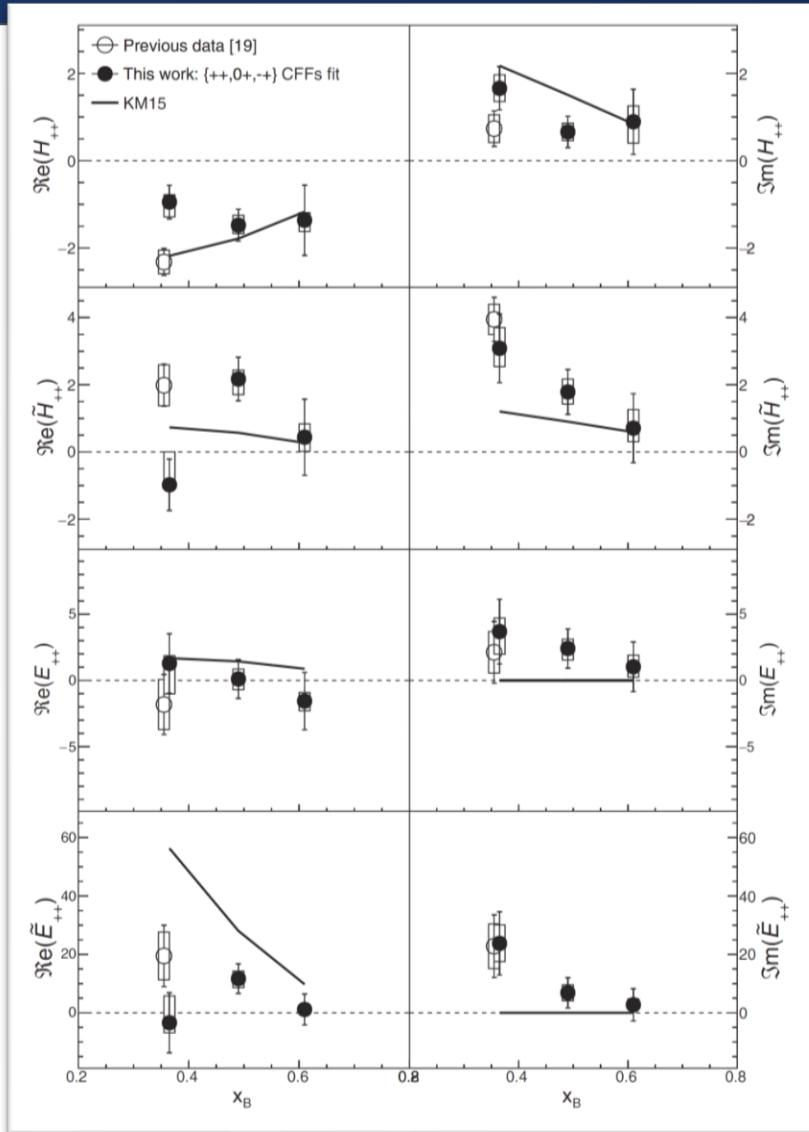


Structure Functions

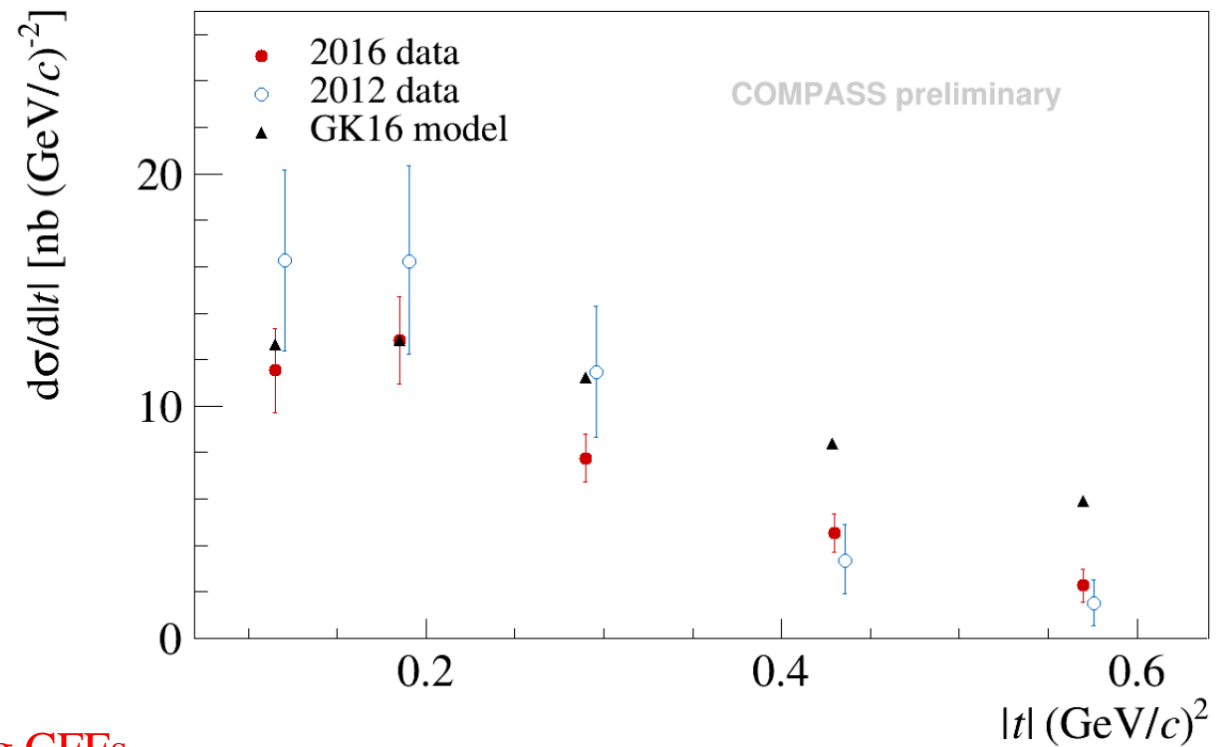
- Solid Markers: Measured $d\sigma_U = d\sigma_T + \epsilon d\sigma_L$
 - Dotted curves: P. Kroll, private communications
- Reasonable agreement in $d\sigma_U$ and $d\sigma_{TT}$
 - $d\sigma_{TT}$ larger than $d\sigma_{LT}$ & $d\sigma_{LT'}$ in general
 - Hint the dominance of $\sigma_T \rightarrow$ as suggested by the GK model
 - GK underestimates both σ_{LT} & $\sigma_{LT'}$
 - Suggest a larger contribution of the longitudinal amplitude than the one expected by GK.
 - Sign difference in σ_{LT}
 - Different from Hall B or COMPASS results
 - Provide useful input for understanding the GPDs involved in the valence domain



DVCS @ Hall A & Exclusive π^0 @ COMPASS



- Preliminary results from 2016 data at low ξ ($\langle x_B \rangle = 0.0096$), with statistics about 2.3 times larger than the published 2012 pilot run.
- **New inputs for phenomenological models.**

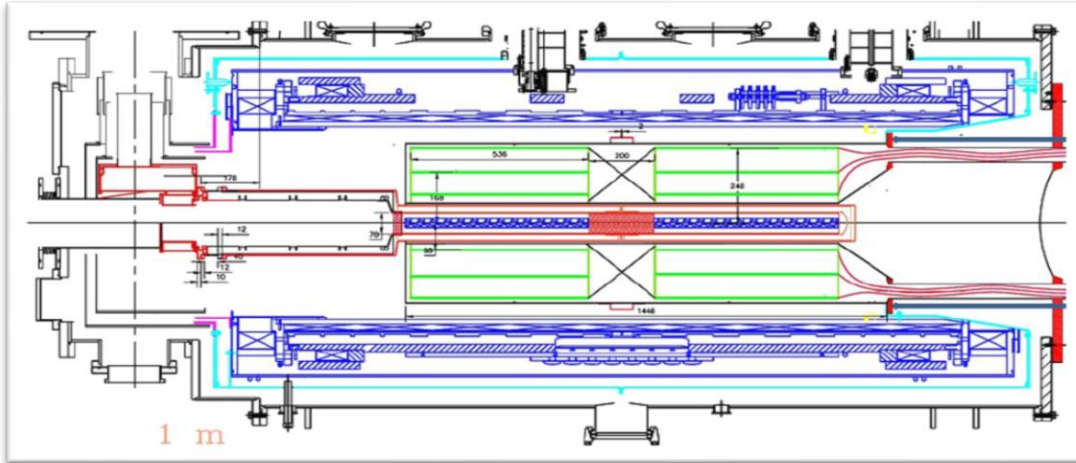
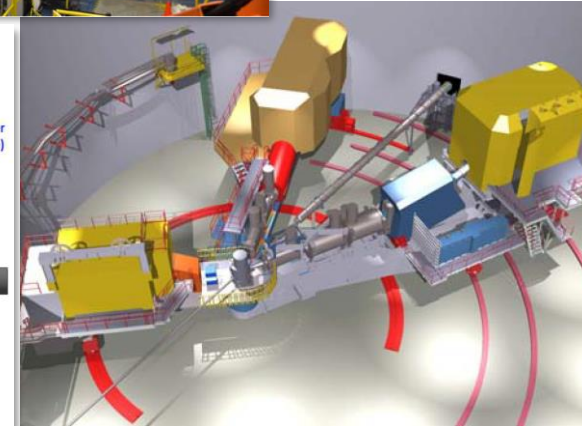
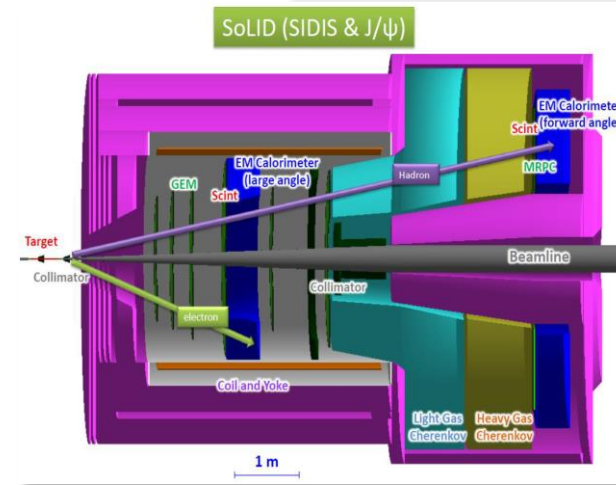


- **First experimental extraction of all four helicity-conserving CFFs**

Near Future

➤ Expect fruitful measurements coming from JLab-12

- Continuation of DVCS & π^0 at Hall C
- DVCS, DVMP, TCS, even DDVCS



Silicone proton recoil detector between target & polarizing magnet

➤ COMPASS/AMBER

- DVCS \rightarrow ReH with charge-spin asymmetry
- DVMP of $\pi^0, \omega, \rho, J/\psi$
- Transversely polarized target in AMBER?

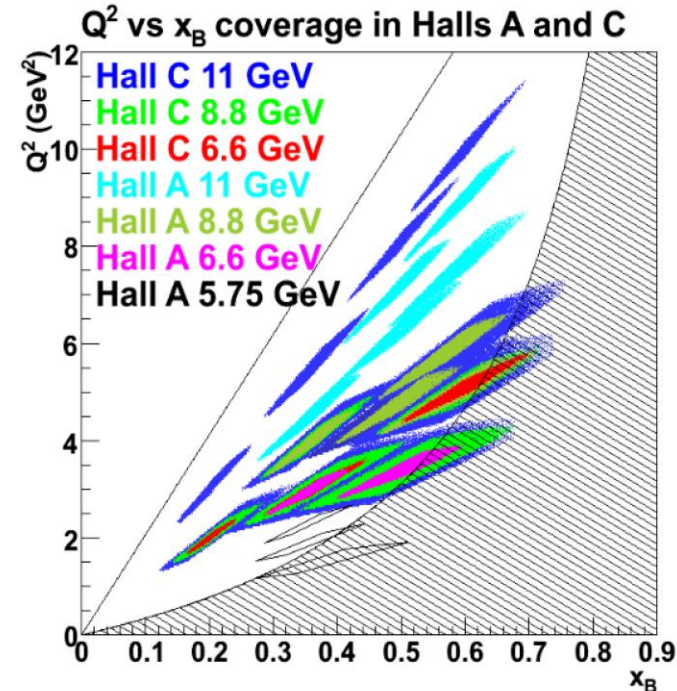
Summary & Outlook

COMPASS: DVCS x-sections with polarized μ^+ and μ^-

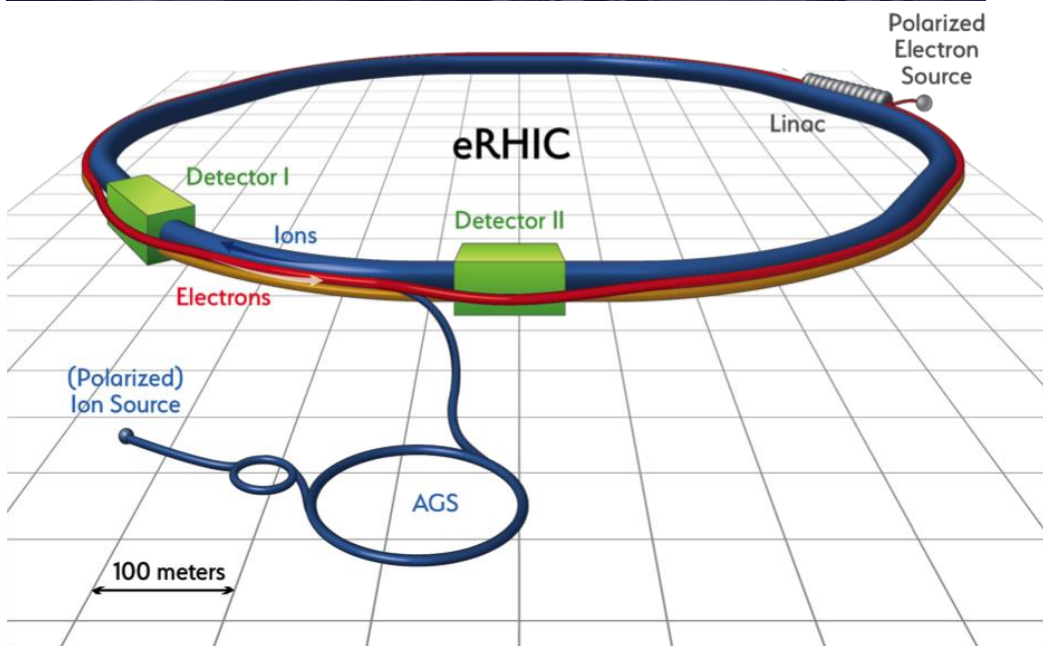
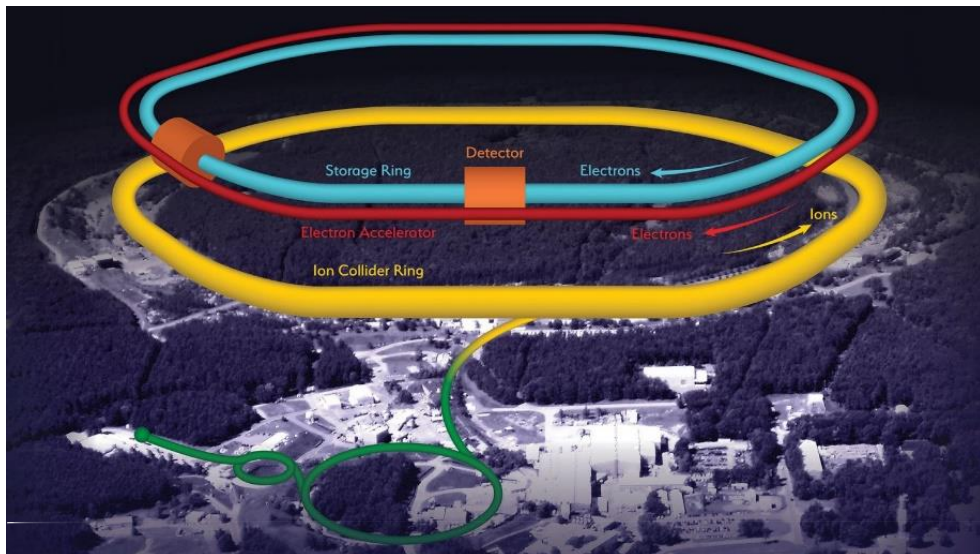
- Beam charge-spin sum $\rightarrow \text{Im}\mathcal{H}(\xi,t) \rightarrow$ Transverse extension of partons
- Beam charge-spin difference $\rightarrow \text{Re}\mathcal{H}(\xi,t) \rightarrow$ D-term, pressure distribution
- Meson Production \rightarrow Flavor decomposition, gluon GPD

Jlab Hall A E12-06-114: Exclusive π^0 Production

- Reasonable description of results by GK model, improvements required
 - Provide inputs for transversity GPD parameterization
 - Extension to higher Q^2 and lower x_B
 - σ_T and σ_L separation of π^0 production
- } Hall C E12-13-010



Looking forward to GPD Measurements at EIC



hadronic calorimeters

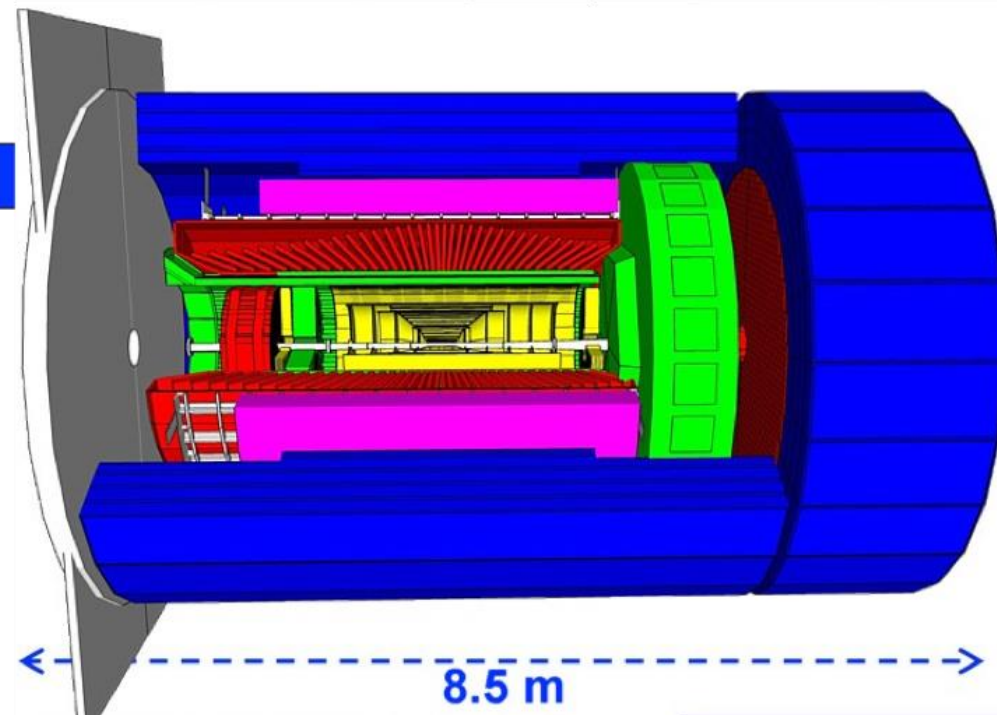
solenoid coils

e/m calorimeters

ToF, DIRC,
RICH detectors

MPG trackers

MAPS tracker



Thank you



Backup Slides

COMPASS⁺⁺ / AMBER



A new QCD facility
at the M2 beam line of the CERN SPS



- Unique beam line with polarised μ^\pm and high-intensity **Pion** beam
- Possible high-intensity **antiproton** and **Kaon** beams, provided by RF-separation technique
- With upgraded apparatus

Proposed physics goals

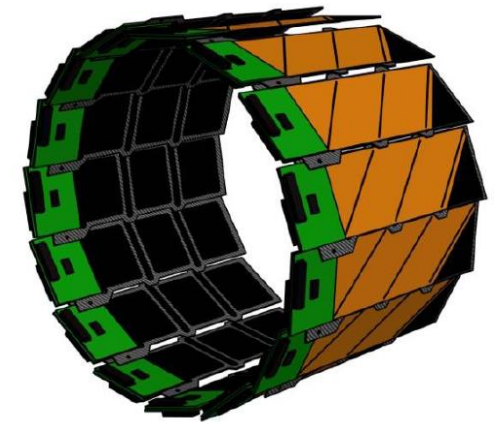
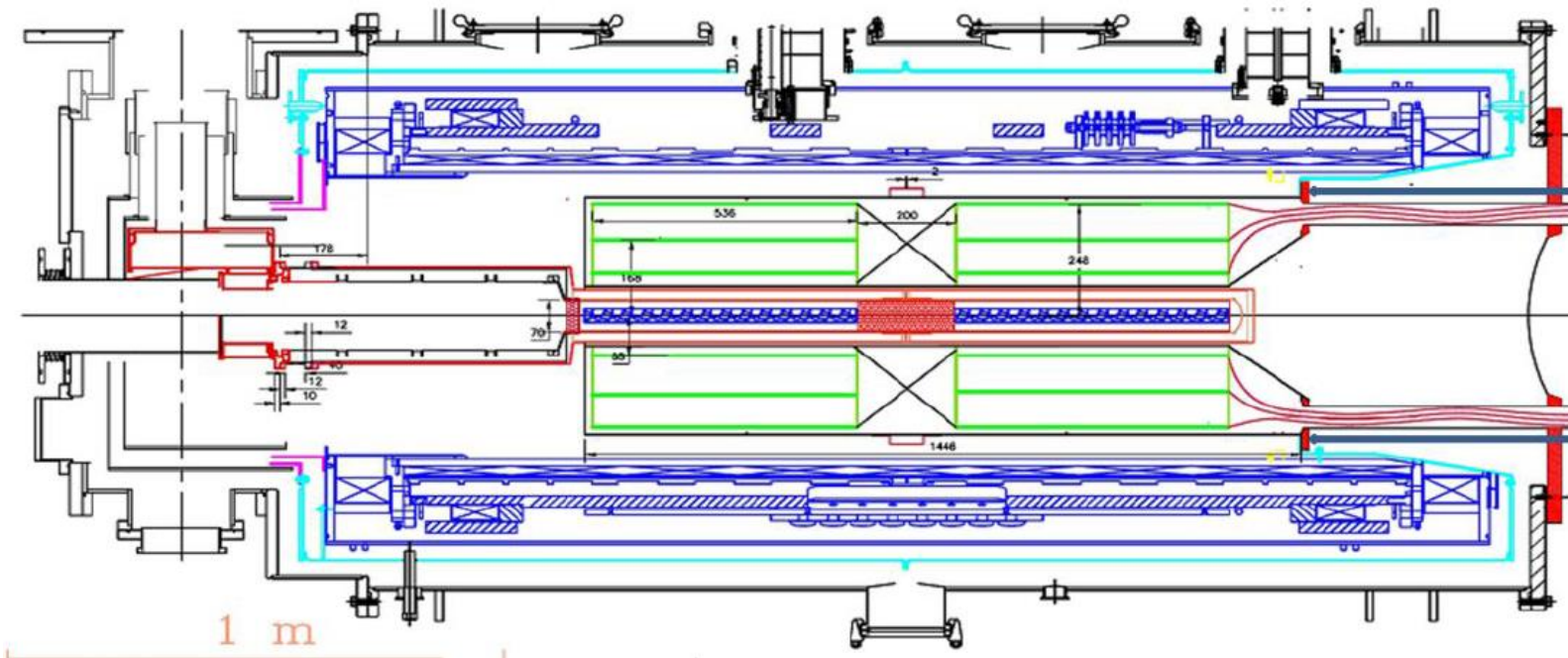
Proton Radius
Meson PDF – gluon PDF
Proton spin structure
3D imaging (TMDs and GPDs)
Hadron spectroscopy
Anti-matter cross section

Program	Physics Goals	Beam Energy [GeV]	Beam Intensity [s ⁻¹]	Trigger Rate [kHz]	Beam Type	Target	Earliest start time, duration	Hardware Additions
μp elastic scattering	Precision proton-radius measurement	100	$4 \cdot 10^6$	100	μ^\pm	high-pressure H2	2022 1 year	active TPC, SciFi trigger, silicon veto,
Hard exclusive reactions	GPD E	160	$2 \cdot 10^7$	10	μ^\pm	NH ₃ [†]	2022 2 years	recoil silicon, modified PT magnet
Input for Dark Matter Search	\bar{p} production cross section	20-280	$5 \cdot 10^5$	25	p	LH2, LHe	2022 1 month	LHe target
\bar{p} -induced Spectroscopy	Heavy quark exotics	12, 20	$5 \cdot 10^7$	25	\bar{p}	LH2	2022 2 years	target spectr.: tracking, calorimetry
Drell-Yan	Pion PDFs	190	$7 \cdot 10^7$	25	π^\pm	C/W	2022 1-2 years	
Drell-Yan (RF)	Kaon PDFs & Nucleon TMDs	~100	10^8	25-50	K^\pm, \bar{p}	NH ₃ [†] , C/W	2026 2-3 years	"active absorber", vertex det.
Primakoff (RF)	Kaon polarisability & pion life time	~100	$5 \cdot 10^6$	> 10	K^-	Ni	non-exclusive 2026 1 year	
Prompt Photons (RF)	Meson gluon PDFs	≥ 100	$5 \cdot 10^6$	10-100	K^\pm π^\pm	LH2, Ni	non-exclusive 2026 1-2 years	hodoscope
K -induced Spectroscopy (RF)	High-precision strange-meson spectrum	50-100	$5 \cdot 10^6$	25	K^-	LH2	2026 1 year	recoil TOF, forward PID
Vector mesons (RF)	Spin Density Matrix Elements	50-100	$5 \cdot 10^6$	10-100	K^\pm, π^\pm	from H to Pb	2026 1 year	

Possible RPD for COMPASS⁺⁺ / AMBER



A recoil proton detector (RPD) is mandatory to ensure the exclusivity. A Silicon detector is included *between* the target surrounded by the modified MW cavity *and* the polarizing magnet



A technology developed at JINR for NICA for the BM@N experiment

No possibility for ToF → PID of p/π with dE/dx
Momentum and trajectory measurements
 $|t|_{\min} \sim 0.1 \text{ GeV}$

Beam Charge-spin Difference



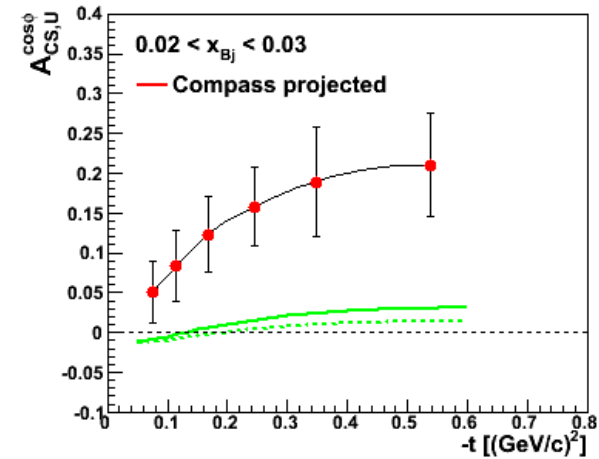
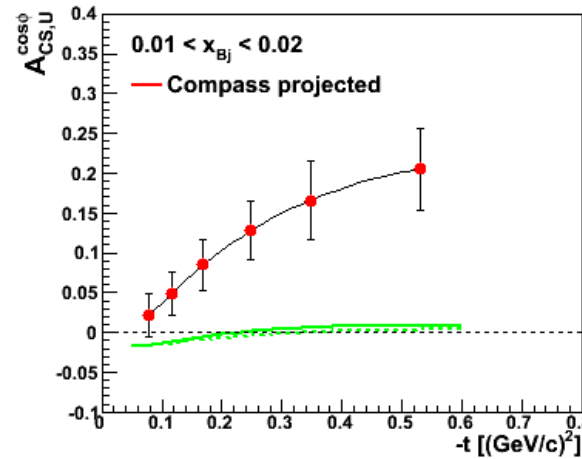
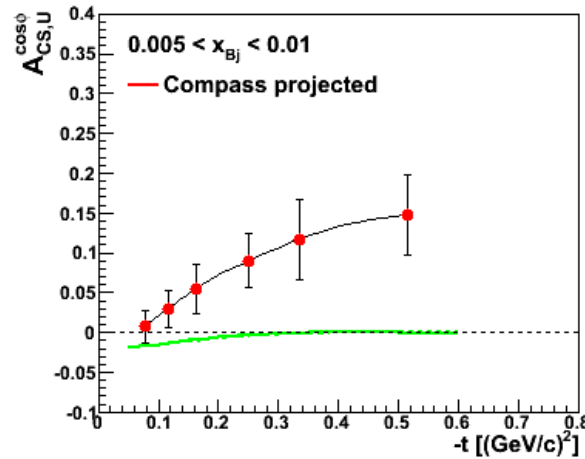
$$\mathcal{D}_{CS,U}(\phi) \equiv d\sigma(\mu^{+\leftarrow}) - d\sigma(\mu^{-\rightarrow}) \rightarrow c_0^I + c_1^I \cos \phi$$

$$BCSA = \mathcal{D}_{CS,U} / S_{CS,U} = A_0 + A_{CS,U}^{\cos\phi} \cos\phi + A_2 \cos 2\phi$$

$$c_1^I = \text{Re } F_1 \mathcal{H}$$

— VGG

— KM10 – fit to data

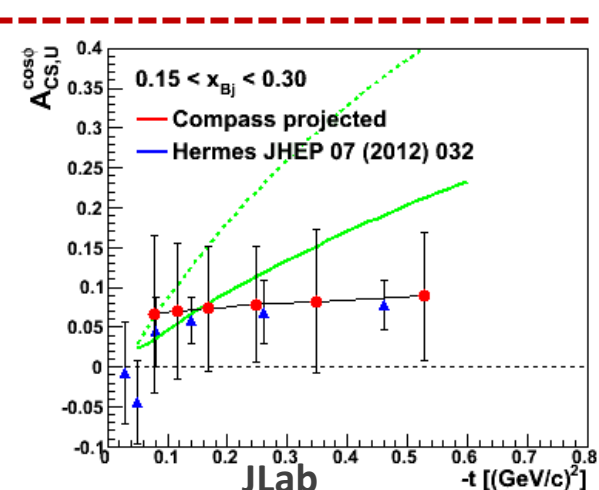
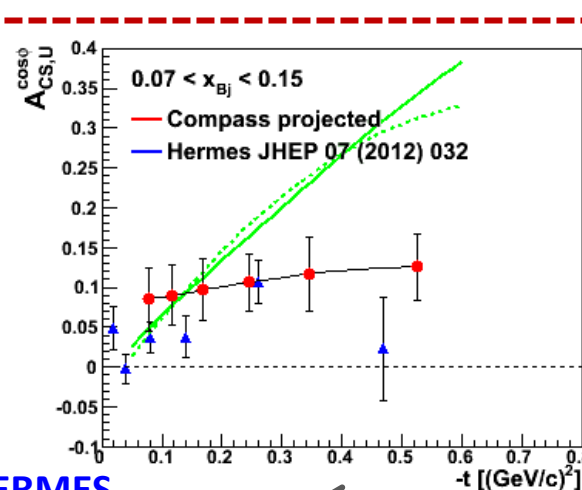
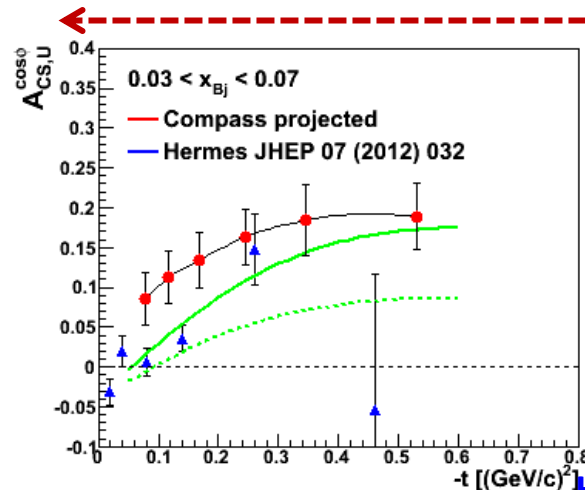


- With $\text{Re } F_1 \mathcal{H}$ and $\text{Im } F_1 \mathcal{H}$
→ Extraction of **D-term**

$\text{Re } \mathcal{H} > 0$ at H1

< 0 at HERMES

Value of x_{Bj} for the node?



← COMPASS 2 years of data E_μ= 160 GeV 1 < Q² < 8 GeV²

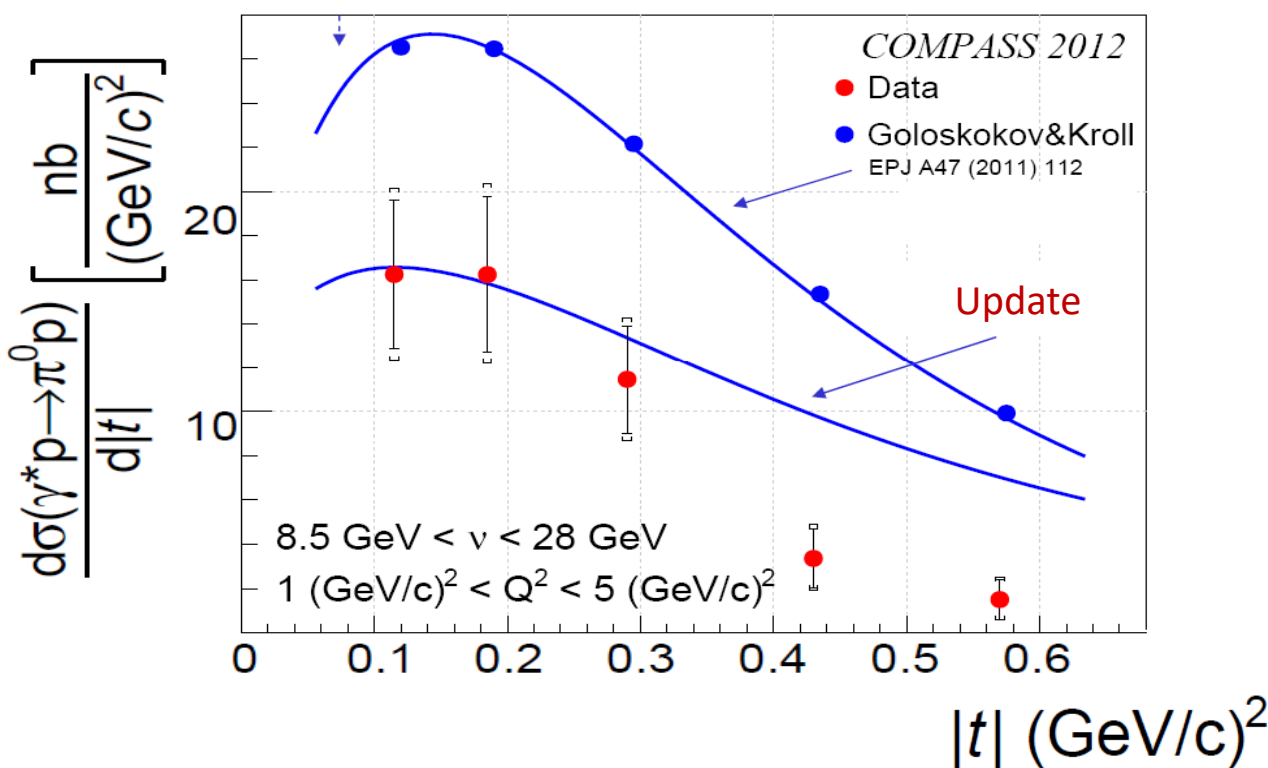
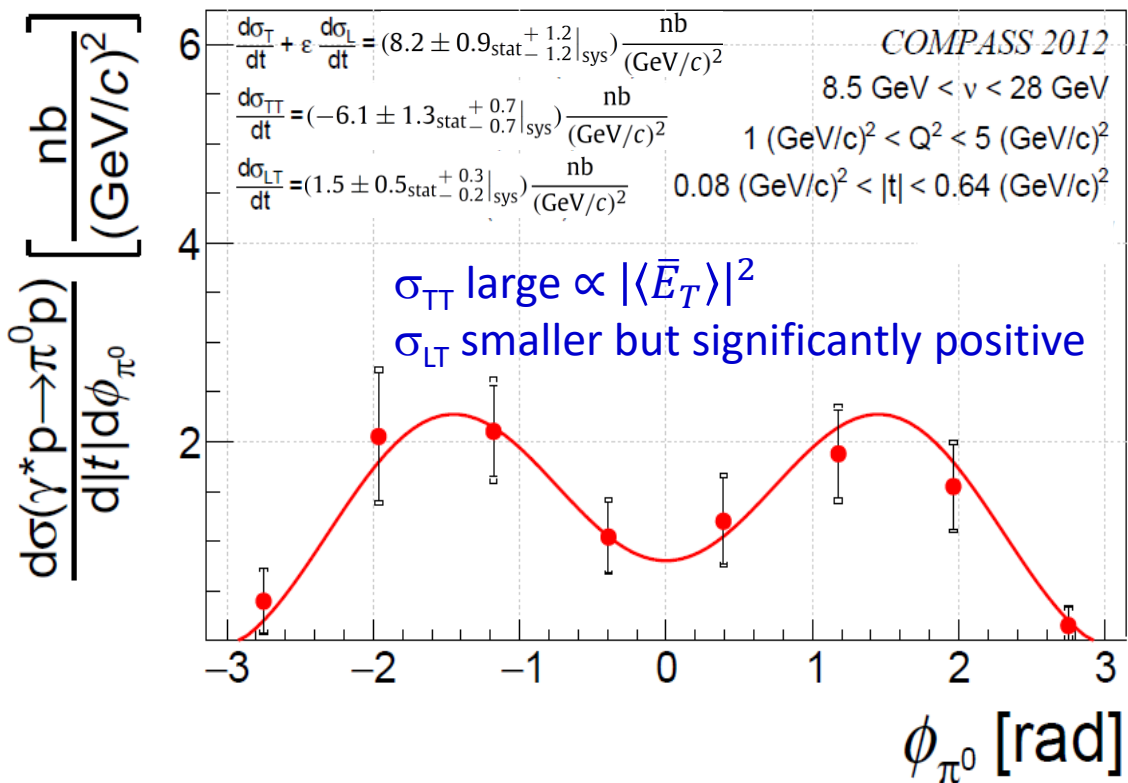
2012 Exclusive π^0 Prod. on Unpolarized Proton



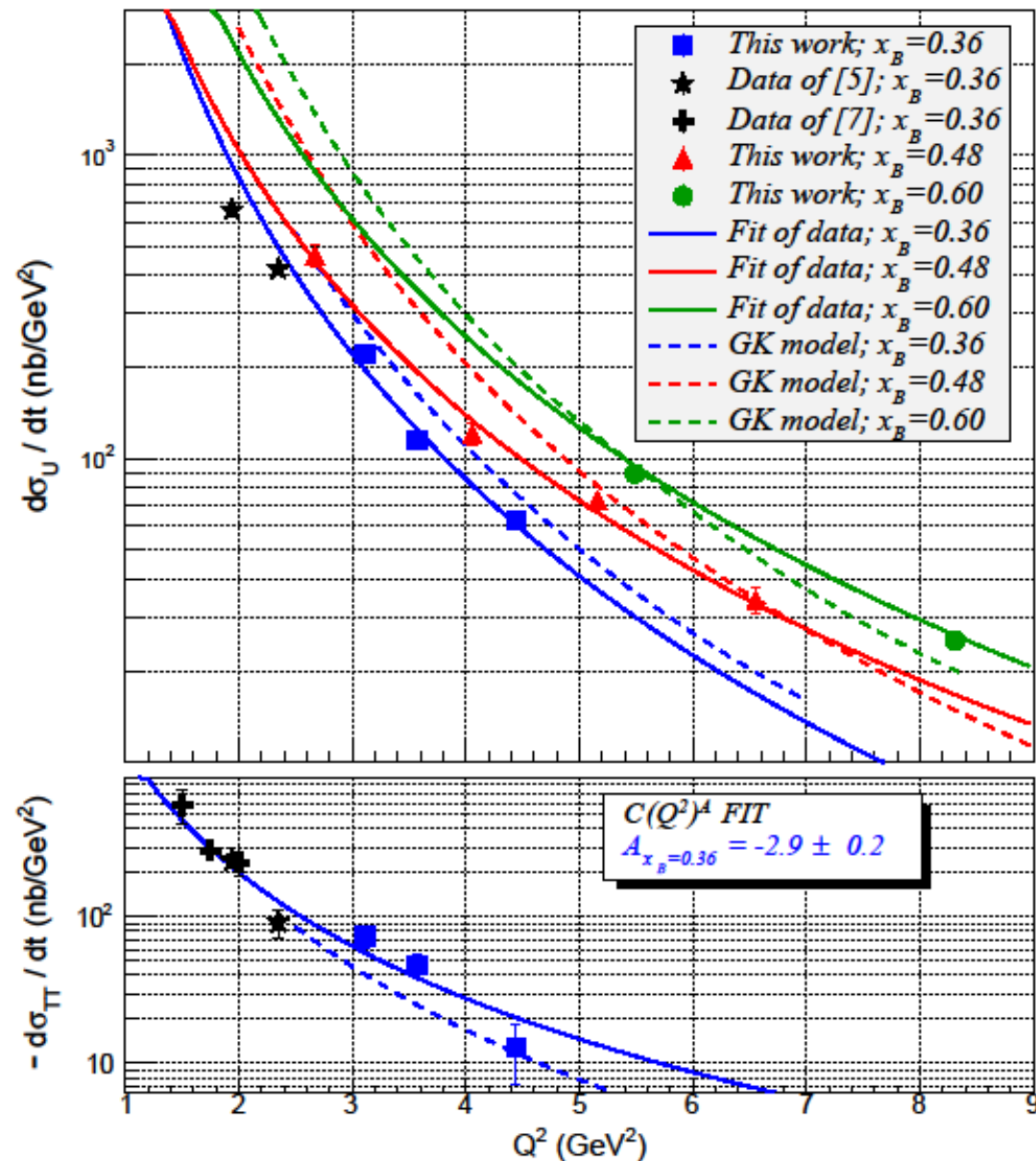
$$\frac{d^2\sigma}{dt d\phi_\pi} = \frac{1}{2\pi} \left[\left(\frac{d\sigma_T}{dt} + \epsilon \frac{d\sigma_L}{dt} \right) + \epsilon \cos 2\phi_\pi \frac{d\sigma_{TT}}{dt} + \sqrt{2\epsilon(1+\epsilon)} \cos \phi_\pi \frac{d\sigma_{LT}}{dt} \right]$$



A dip at small t would indicate the significance of $\overline{E_T}$



Q² Dependence



- Dashed curves: P. Kroll, private communications
- Solid Markers: Experimental measurements $\langle t' \rangle = 0.1 \text{ GeV}^2$

■ This work, $x_B = 0.36$

▲ This work, $x_B = 0.48$

● This work, $x_B = 0.60$

★ E. Fuchey *et al*, Phys. Rev. C 83, 025201 (2011)

■ M. Defurne *et al*, Phys. Rev. Lett. 117, 262001 (2016)

➤ $C(Q^2)^A \exp(-Bt')$ fit to experimental results of $d\sigma_U$ in different $x_B \rightarrow$ solid curves

$x_B = 0.36 \rightarrow A = -3.3 \pm 0.1$

$x_B = 0.48 \rightarrow A = -2.9 \pm 0.1$

$x_B = 0.60 \rightarrow A = -3.1 \pm 0.1$

➤ Q^2 dependence closer to Q^{-6} , rather than Q^{-8} as expected for σ_T at high Q^2

Hall A DVCS Cross Sections

$x_B=0.36, Q^2=3.7\text{GeV}^2, t=-0.33\text{GeV}^2;$

$x_B=0.48, Q^2=5.4\text{ GeV}^2, t=-0.33\text{ GeV}^2;$

$x_B=0.6, Q^2=8.4\text{ GeV}^2, t=-0.91\text{ GeV}^2$

