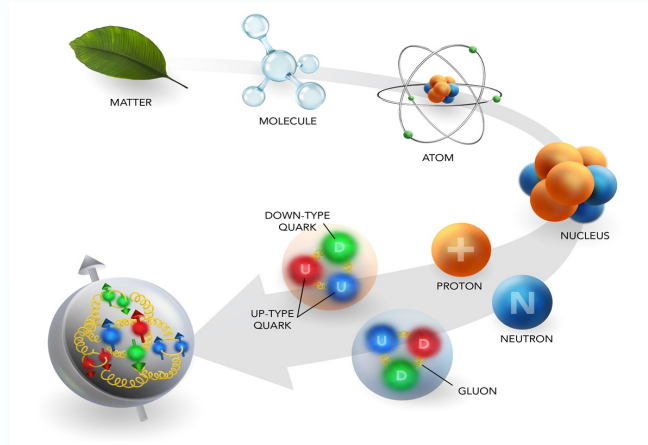


Lecture 2: TMD theory

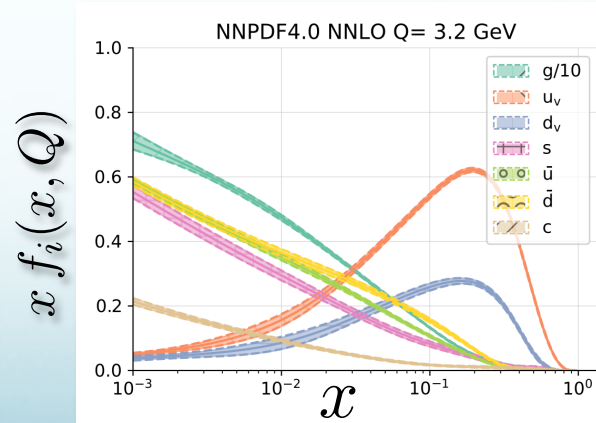
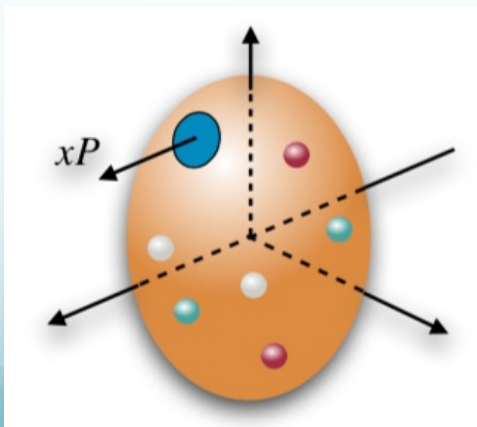
- TMD basic, parton model, symmetry

EIC era: Imaging of proton and nucleus

- Unraveling the mysteries of relativistic hadronic bound states

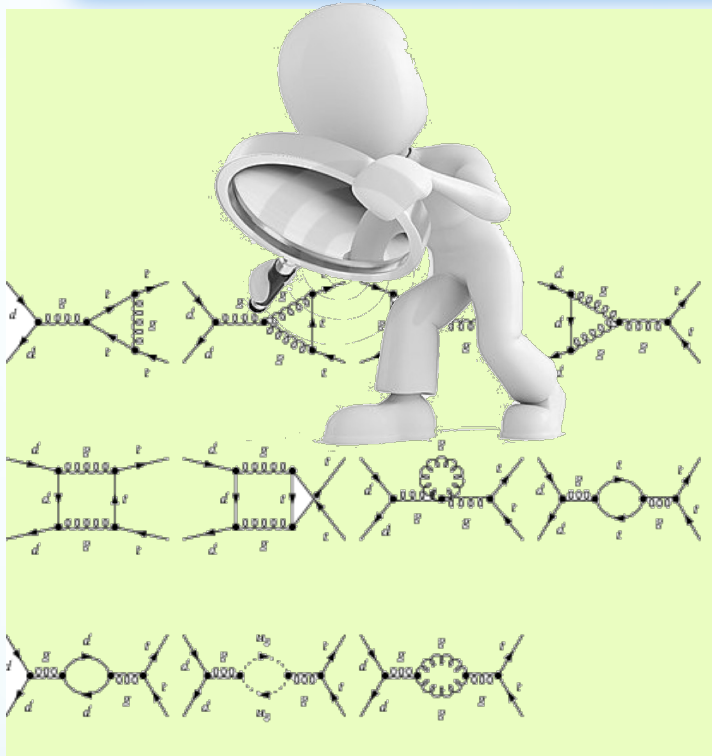


- Beyond 1D: collinear PDFs provide 1D structure – longitudinal motion



[Proton Movie](#)

Direction 1: deep in perturbative direction

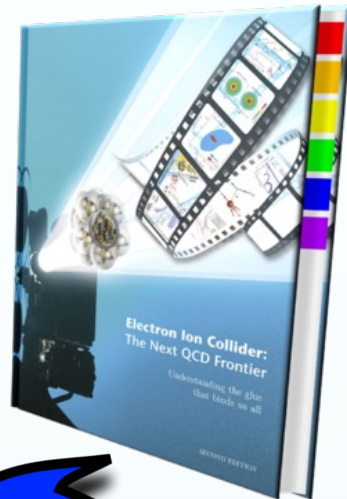


Dave Soper

Jets, Jet substructure, W/Z/H+Jet, ...
NLO, NNLO, Resummation, multi-
loops/legs, ...

$$\sigma(Q) = f_{q,g}(x) \otimes \hat{\sigma}_{q,g}(Q)$$

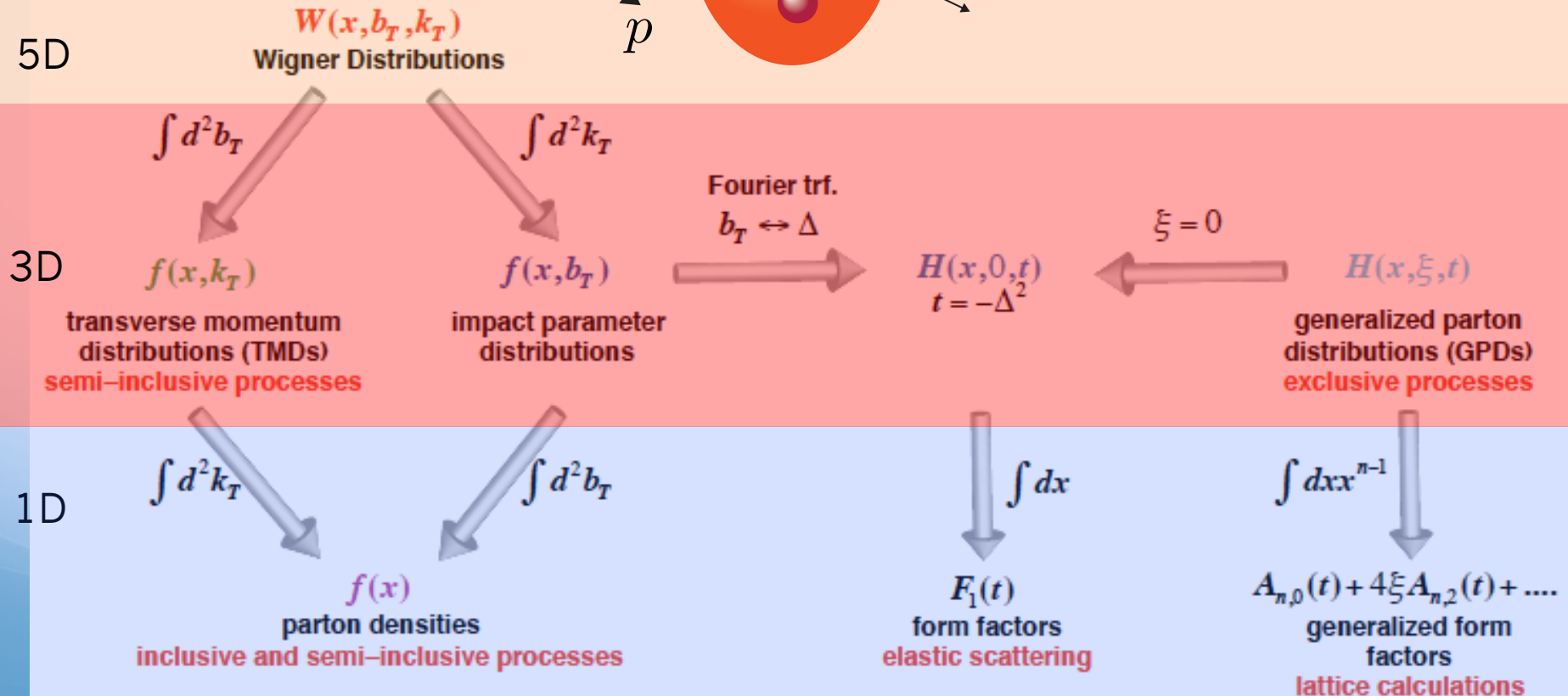
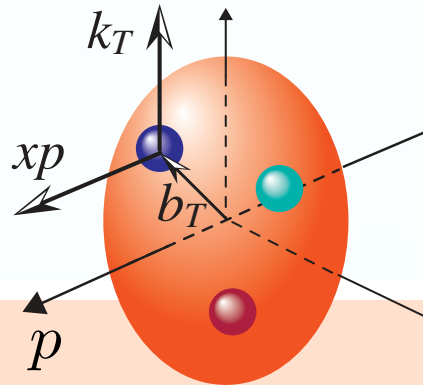
Direction 2: quantum tomography of nucleons



$$\sigma(Q) = f_{q,g}(x) \otimes \hat{\sigma}_{q,g}(Q)$$

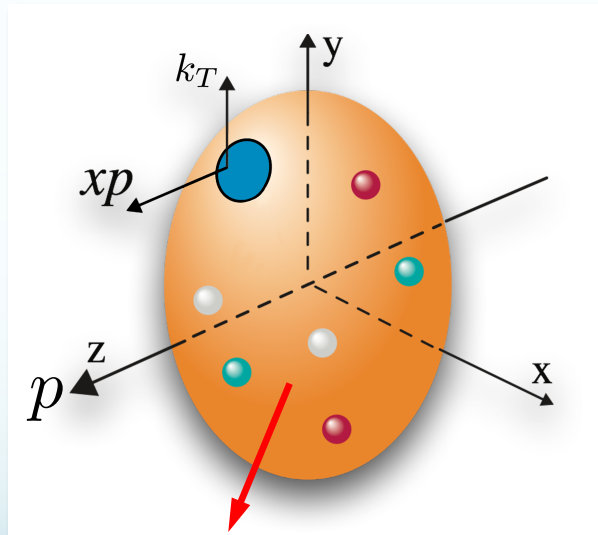
Unified view: internal landscape

- Wigner distributions: a quantum version of phase-space distribution

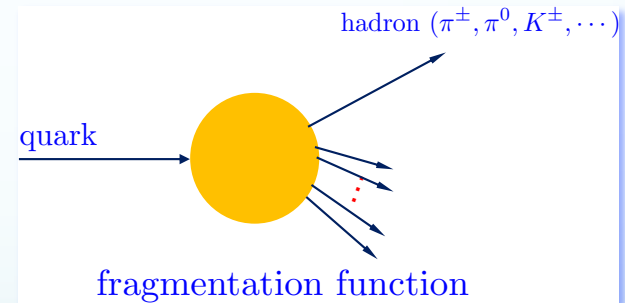


Transverse Momentum Dependent distributions (TMDs)

- 3D imaging in momentum space
 - Both longitudinal and transverse motion
 - What are the quantum correlations between the motion of the quarks/gluons, their spin and the spin of the proton? (TMD PDFs)
 - Similarly precision information on hadronization (TMD FFs)



Proton spin



TMDs with polarization

Leading Twist TMDs



Nucleon Spin



Quark Spin

TMD PDFs

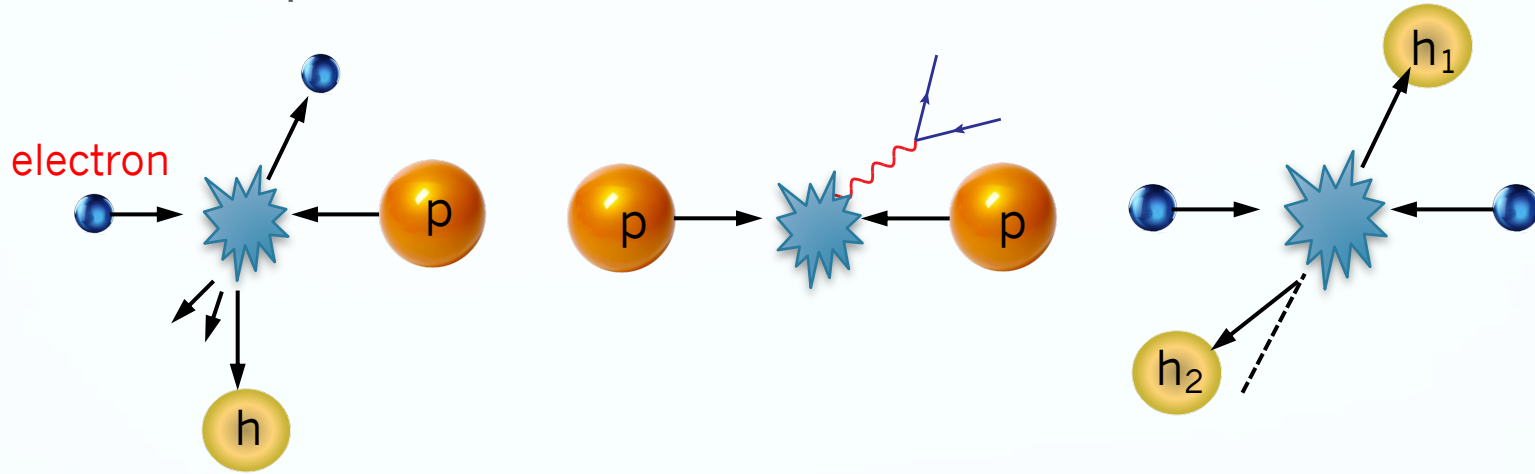
		Quark Polarization		
		Un-Polarized (U)	Longitudinally Polarized (L)	Transversely Polarized (T)
Nucleon Polarization	U	$f_1 =$		$h_1^\perp =$ - Boer-Mulders
	L		$g_{1L} =$ - Helicity	$h_{1L}^\perp =$ -
	T	$f_{1T}^\perp =$ - Sivers	$g_{1T} =$ - Transversal Helicity	$h_1 =$ - Transversity $h_{1T}^\perp =$ -

TMD FFs

		Quark Polarization		
		U	L	T
Pion	D_1			H_1^\perp Collins

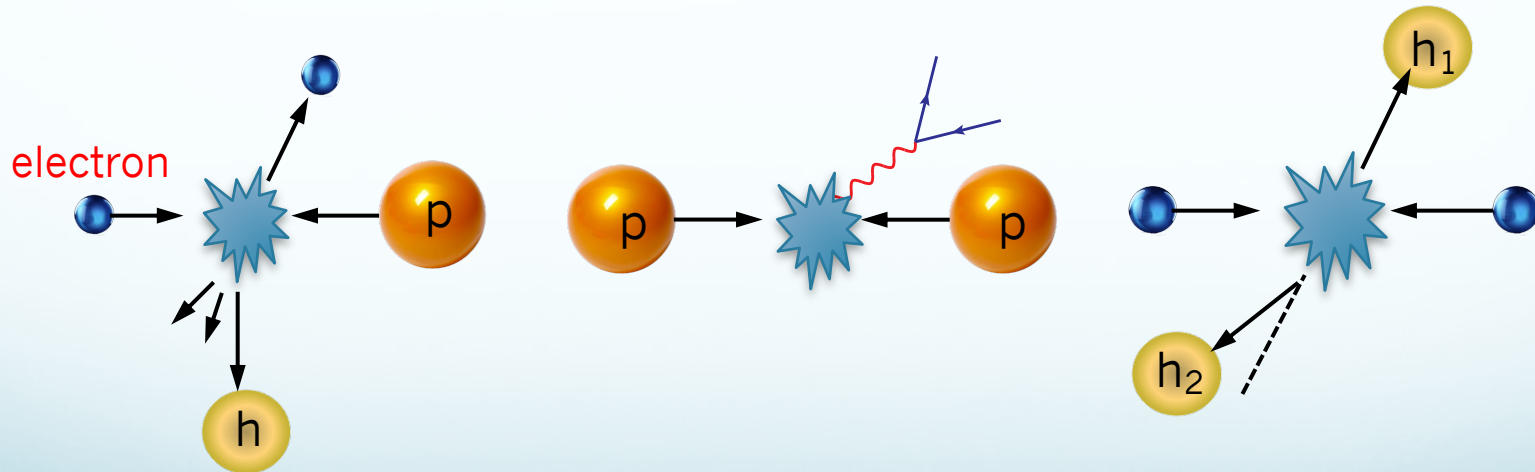
Processes to extract TMDs

- Standard processes: SIDIS, Drell-Yan, dihadron in e^+e^-



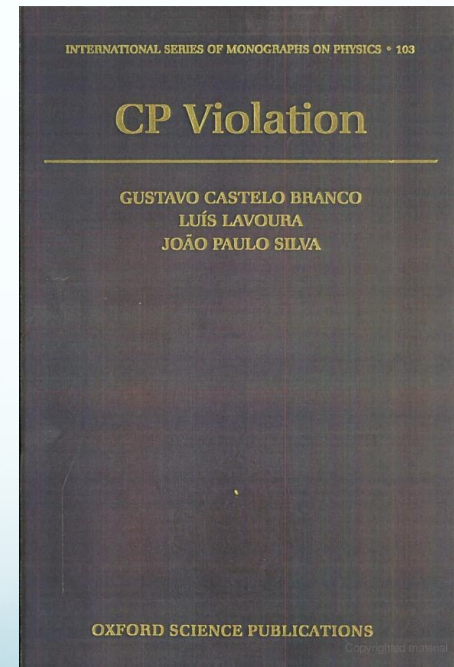
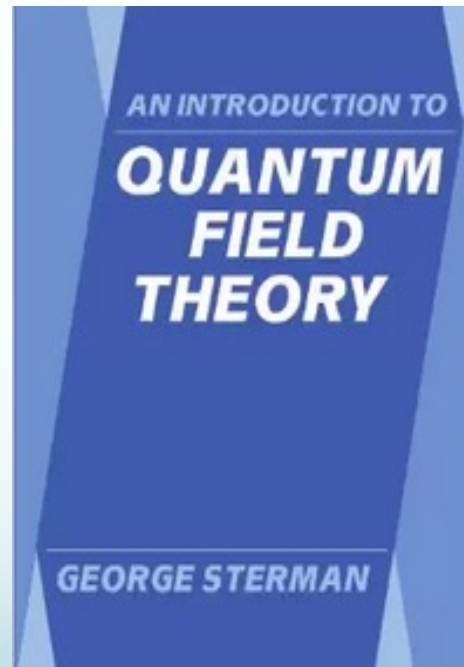
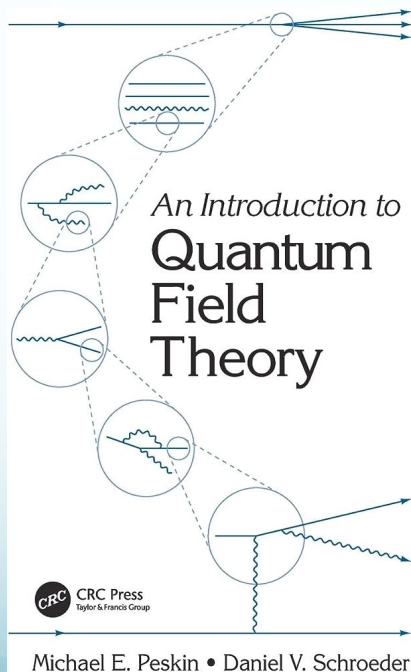
TMD parton model

- Two major questions we have to answer
 - How many TMD distributions are needed in order to fully characterize the proton structure?
 - How are the cross section (observables) connected to the TMDs?



Good textbooks

- Understand C, P, T discrete symmetry properties of the correlation function
 - Most textbooks on quantum field theory will give discussion on this topic, such as Peskin, Sterman: appendix of Sterman's book
 - If you want extensive discussion, see this book



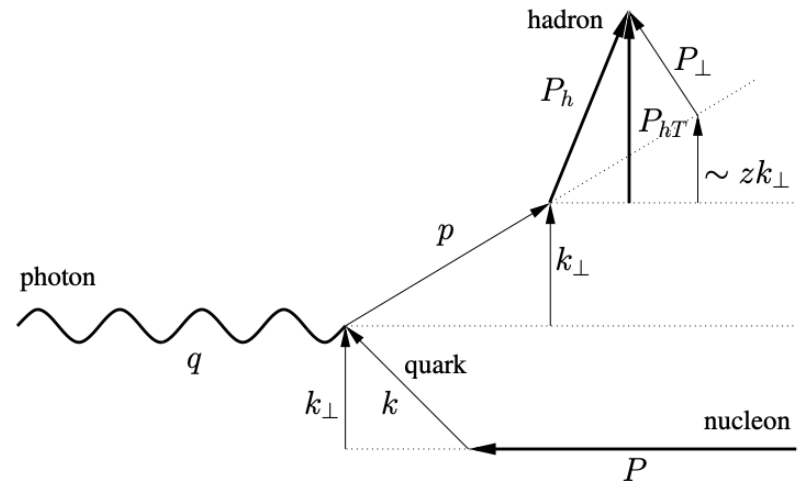
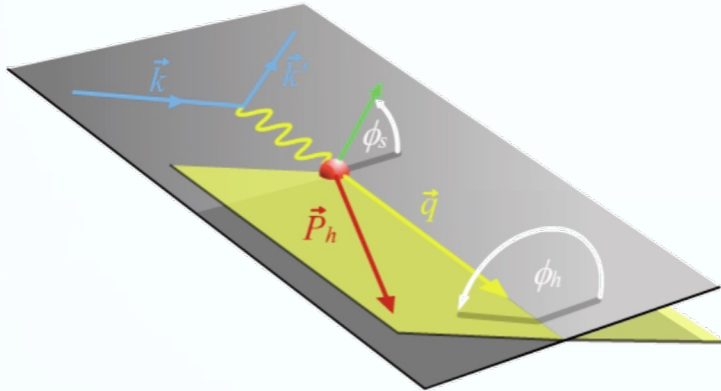
Operator analysis

- Operator analysis to figure out how many distributions are needed to characterize the nucleon structure
- For details, see
 - Mulders, Tangerman hep-ph/9403227
 - Mulders, Tangerman hep-ph/9510301
 - Mulders, <http://www.nat.vu.nl/~mulders/correlations-Onew.pdf>
- From now on, we will look at [this hand-writing note](#)

TMD parton model

- Here I provide [a detailed note](#) for deriving SIDIS cross section

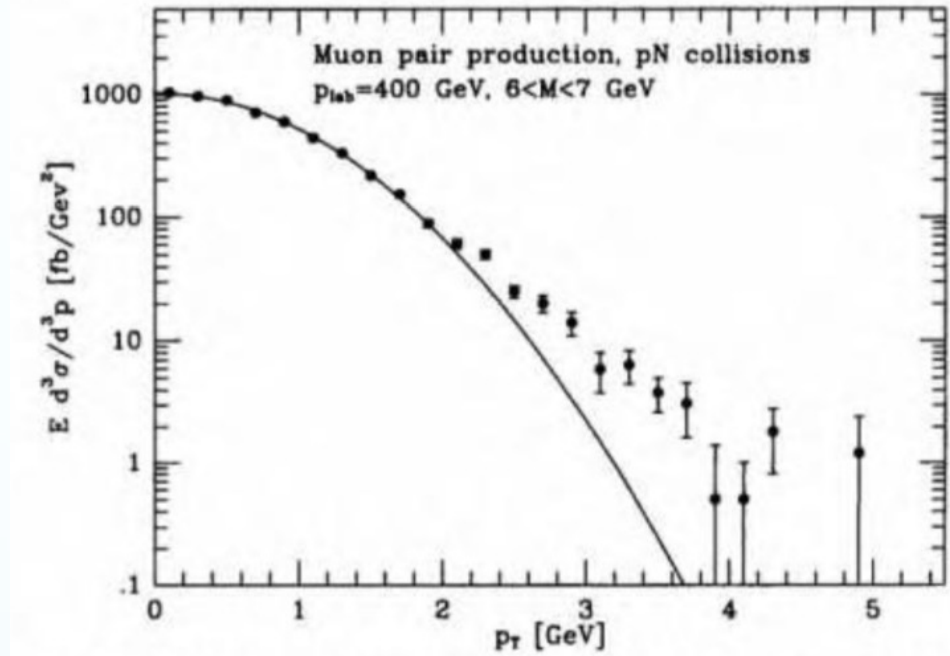
$$\ell + p^\uparrow \rightarrow \ell' + \pi(p_T) + X$$



$$\frac{d\sigma}{dx_B dy dz_h d^2 P_{hT}} = \sigma_0 \sum_q e_q^2 \int d^2 k_T d^2 p_T f_{q/p}(x_B, k_T^2) D_{h/q}(z_h, p_T^2) \delta^2(z_h \vec{k}_T + \vec{p}_T - \vec{P}_{hT})$$

Parton model for TMDs

- At low p_T , the distribution is described by a Gaussian form



$$f_q(x, k_T^2) = f_q(x)g(k_T)$$

$$g(k_T) = \frac{1}{\pi \langle k_T^2 \rangle} e^{-k_T^2 / \langle k_T^2 \rangle}$$

$$\text{since } \int d^2 k_T f_q(x, k_T^2) = f_q(x)$$

$$\rightarrow \int d^2 k_T g(k_T) = 1$$

TMD parton model phenomenology

- In the early days of TMD studies, one typically uses a Gaussian model to describe transverse momentum dependence

$$f_{q/p}(x, \mathbf{k}_\perp) = f_{q/p}(x) \frac{e^{-\mathbf{k}_\perp^2 / \langle \mathbf{k}_\perp^2 \rangle}}{\pi \langle \mathbf{k}_\perp^2 \rangle}$$

$$D_{h/q}(z, \mathbf{p}_\perp) = D_{h/q}(z) \frac{e^{-\mathbf{p}_\perp^2 / \langle \mathbf{p}_\perp^2 \rangle}}{\pi \langle \mathbf{p}_\perp^2 \rangle}$$

$$\begin{aligned} \frac{d\sigma^{\ell+p \rightarrow \ell' h X}}{dx_B dQ^2 dz_h dP_T^2} &= \frac{2\pi^2\alpha^2}{(x_B s)^2} \frac{[1 + (1-y)^2]}{y^2} \\ &\times \sum_q e_q^2 \int d^2\mathbf{k}_\perp d^2\mathbf{p}_\perp \delta^{(2)}(\mathbf{P}_T - z_h \mathbf{k}_\perp - \mathbf{p}_\perp) f_{q/p}(x, \mathbf{k}_\perp) D_{h/q}(z, \mathbf{p}_\perp) \\ &\equiv \frac{2\pi^2\alpha^2}{(x_B s)^2} \frac{[1 + (1-y)^2]}{y^2} F_{UU}. \end{aligned}$$

$$F_{UU} = \sum_q e_q^2 f_{q/p}(x_B) D_{h/q}(z_h) \frac{e^{-P_T^2 / \langle P_T^2 \rangle}}{\pi \langle P_T^2 \rangle}$$

$$\langle P_T^2 \rangle = \langle \mathbf{p}_\perp^2 \rangle + z_h^2 \langle \mathbf{k}_\perp^2 \rangle.$$

TMD parton model works well

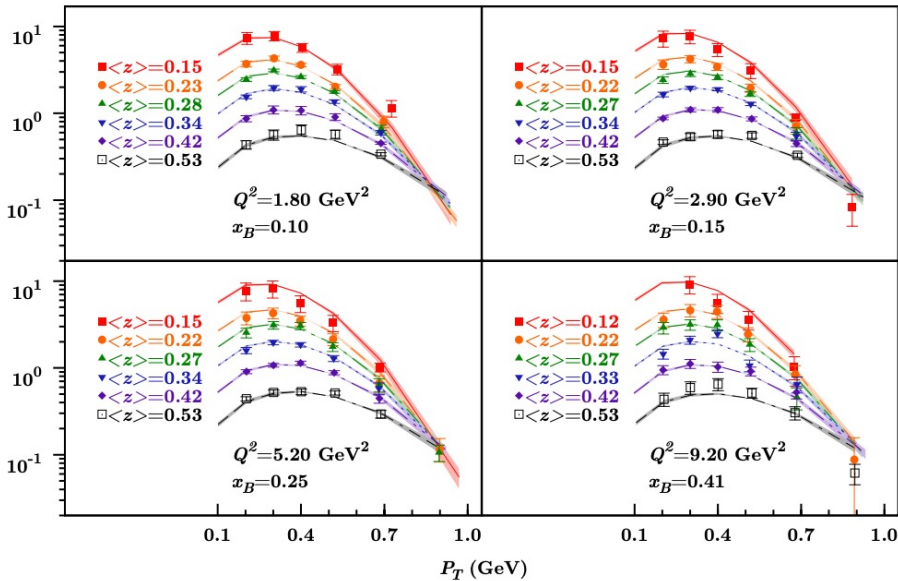
- When the energy (Q) range spans relatively small, TMD parton model does a decent job [fitted parameters]

$$\langle k_{\perp}^2 \rangle = 0.57 \pm 0.08 \text{ GeV}^2$$

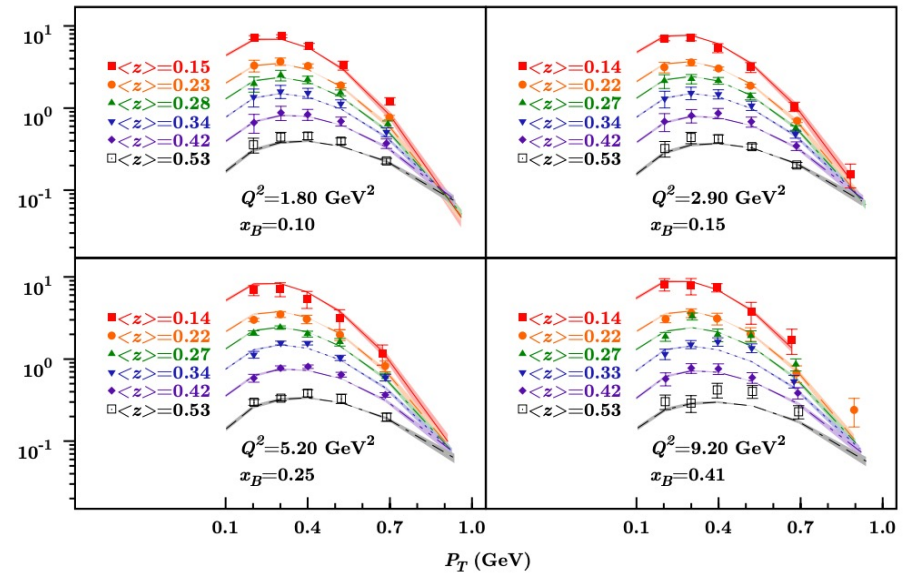
$$\langle p_{\perp}^2 \rangle = 0.12 \pm 0.01 \text{ GeV}^2$$

Anselmino, et.al. 1312.6261

HERMES $M_D^{\pi^+}$

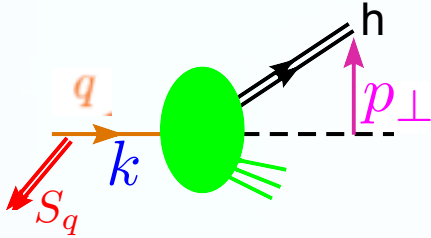


HERMES $M_D^{\pi^-}$



Polarized TMDs: Collins function

- Collins function: unpolarized hadron from a transversely polarized quark

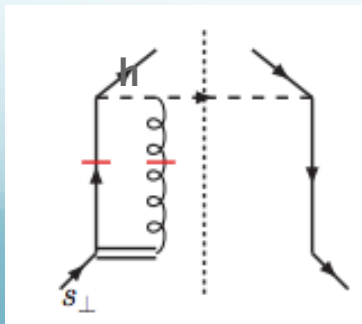


$$D_{h/q}(z, p_\perp) = D_1^q(z, p_\perp^2) + \frac{1}{zM_h} H_1^{\perp q}(z, p_\perp^2) \vec{S}_q \cdot (\hat{k} \times p_\perp)$$

Spin-independent

Spin-dependent

- ✓ 2002: A. Metz studied the universality property of Collins function in a model-dependent way – very subtle – finally found it is universal between SIDIS and e^+e^-
- ✓ 2004: Collins and Metz have general arguments
- ✓ 2008: Yuan generalizes to pp
- ✓ 2010: Boer, Kang, Vogelsang, and Yuan perform perturbative tail calculation, demonstrate the gauge link does not contribute

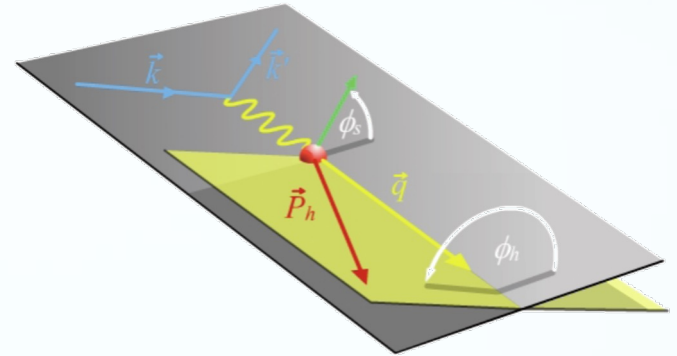


$$H_1^{\perp \text{SIDIS}}(z, p_\perp^2) = H_1^{\perp e^+e^-}(z, p_\perp^2) = H_1^{\perp \text{pp}}(z, p_\perp^2)$$

Metz 02, Collins, Metz 04, Yuan 08,
Boer, Kang, Vogelsang, Yuan, PRL 10, ...

Polarized TMDs: measurements

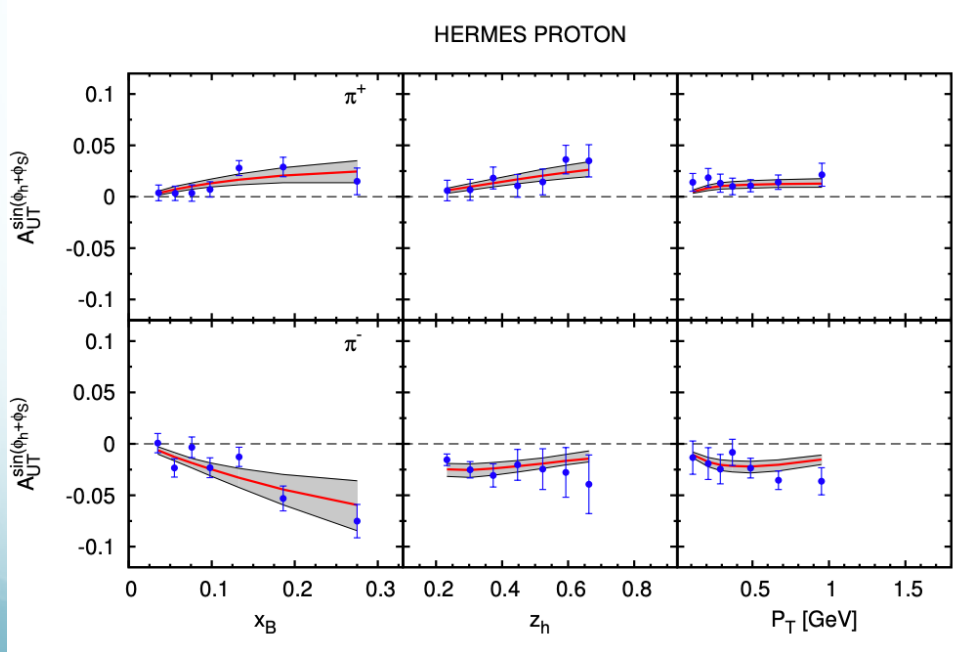
- Collins asymmetry in SIDIS



$$\frac{d^5\sigma(S_\perp)}{dx_B dy dz_h d^2P_{h\perp}} = \sigma_0(x_B, y, Q^2) \left[F_{UU} + \sin(\phi_h + \phi_s) \frac{2(1-y)}{1+(1-y)^2} F_{UT}^{\sin(\phi_h + \phi_s)} + \dots \right]$$

$$A_{UT}^{\sin(\phi_h + \phi_s)} \equiv 2\langle \sin(\phi_h + \phi_s) \rangle = \frac{\sigma_0(x_B, y, Q^2)}{\sigma_0(x_B, y, Q^2)} \frac{2(1-y)}{1+(1-y)^2} \frac{F_{UT}^{\sin(\phi_h + \phi_s)}}{F_{UU}}$$

$$F_{UT}^{\sin(\phi_h + \phi_s)} \propto \int d^2k_T d^2p_T h_1(x_B, k_T) H_1^\perp(z, p_T) \delta^2(z_h \vec{k}_T + \vec{p}_T - \vec{P}_{hT})$$



Anselmino, et.al. 1510.05389

Summary

- Operator analysis allows us to determine the independent TMD correlators
- To probe a small transverse momentum, one has to measure a small transverse momentum for the final state particles
- TMD parton model works well when the Q range spans relatively small